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Cross-Spectral Analysis of Global Oscillation Time Series

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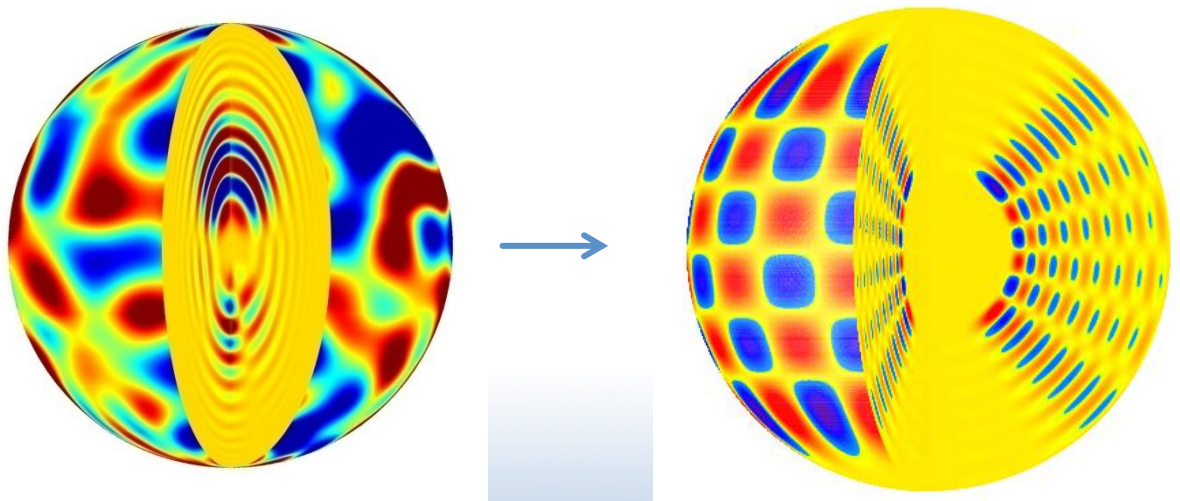
Kiepenheuer-Institut für Sonnenphysik, Freiburg

Tenerife, March 12, 2014

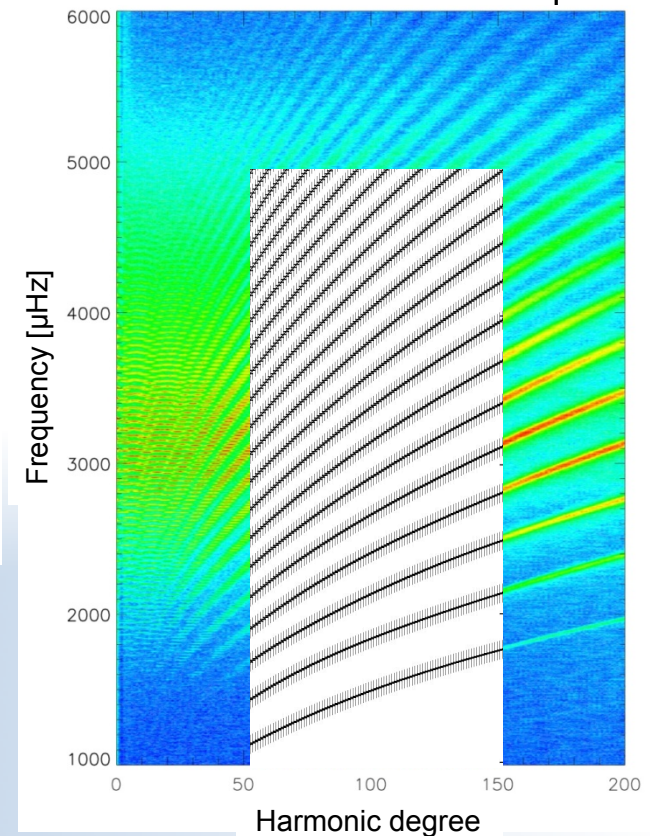
Many discoveries by using the following concept:

- Decomposing normal standing modes on the basis of spherical harmonics

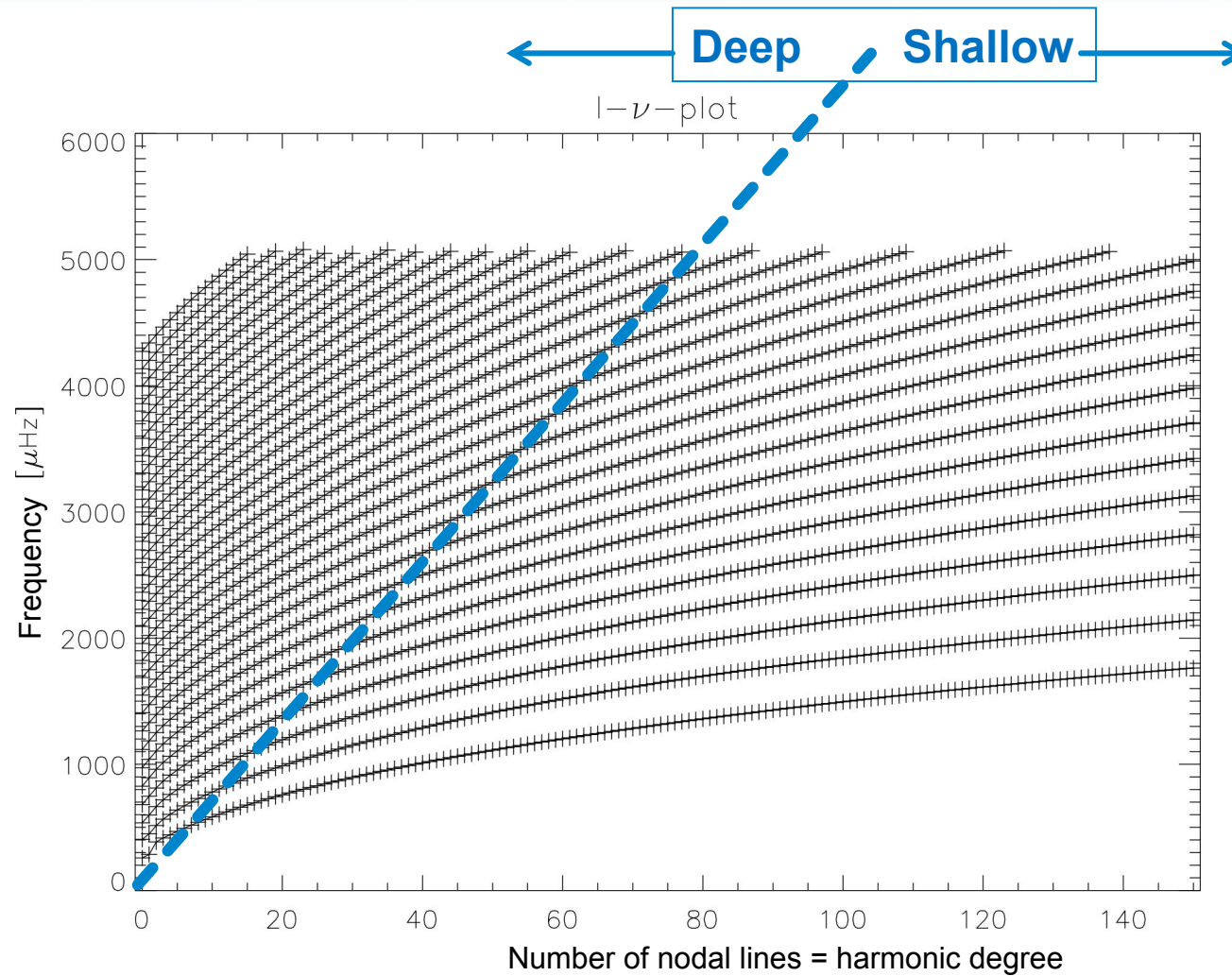
$$\xi(R_{\odot}, \theta, \phi, t) = \text{Re} \left[\sum_{n,l,m} A_{nlm}(t) \xi_{nlm}(R_{\odot}) Y_{lm}(\theta, \phi) \exp(i\omega_{nlm}t + \varphi_{nlm}) \right]$$



- Results are averages over the entire Sun
 - No longitudinal information
 - Symmetric in latitude across the equator



Penetration Depth of Modes



Perturbing the equilibrium model with a slow flow (small perturbation)

$$\begin{aligned}
 -\omega_k^2 \rho_0 \xi_k &= -\nabla p_1 + \rho_0 \mathbf{g}_1 + \rho_1 \mathbf{g}_0 \\
 -\omega_k^2 \rho_0 \xi_k - 2i\omega_k \rho_0 (\mathbf{v} \cdot \nabla) \xi_k &= -\nabla p_1 + \rho_0 \mathbf{g}_1 + \rho_1 \mathbf{g}_0
 \end{aligned}$$

$$-\rho_0 \omega_k^2 \xi_k = H_0(\xi_k) + \varepsilon H_1 \xi_k$$

Mode Coupling: Perturbation matrix elements for calculation of new eigenvalues

$$H_{k'k} = \langle \xi_k | 2i\omega_k \rho_0 (\mathbf{v} \cdot \nabla) | \xi_{k'} \rangle$$

Eigenvalues of perturbation matrix are frequency corrections: $\tilde{\omega}_k = \omega_k + \delta\omega_{1,k}$

Decomposition into a

- toroidal flow (includes differential rotation) and
- poloidal flow (includes meridional flow and giant cells)

$$v(r) = \sum_{s=0}^{\infty} \sum_{t=-s}^s T_s^t(r; \mu; \dot{A}) + P_s^t(r; \mu; \dot{A})$$

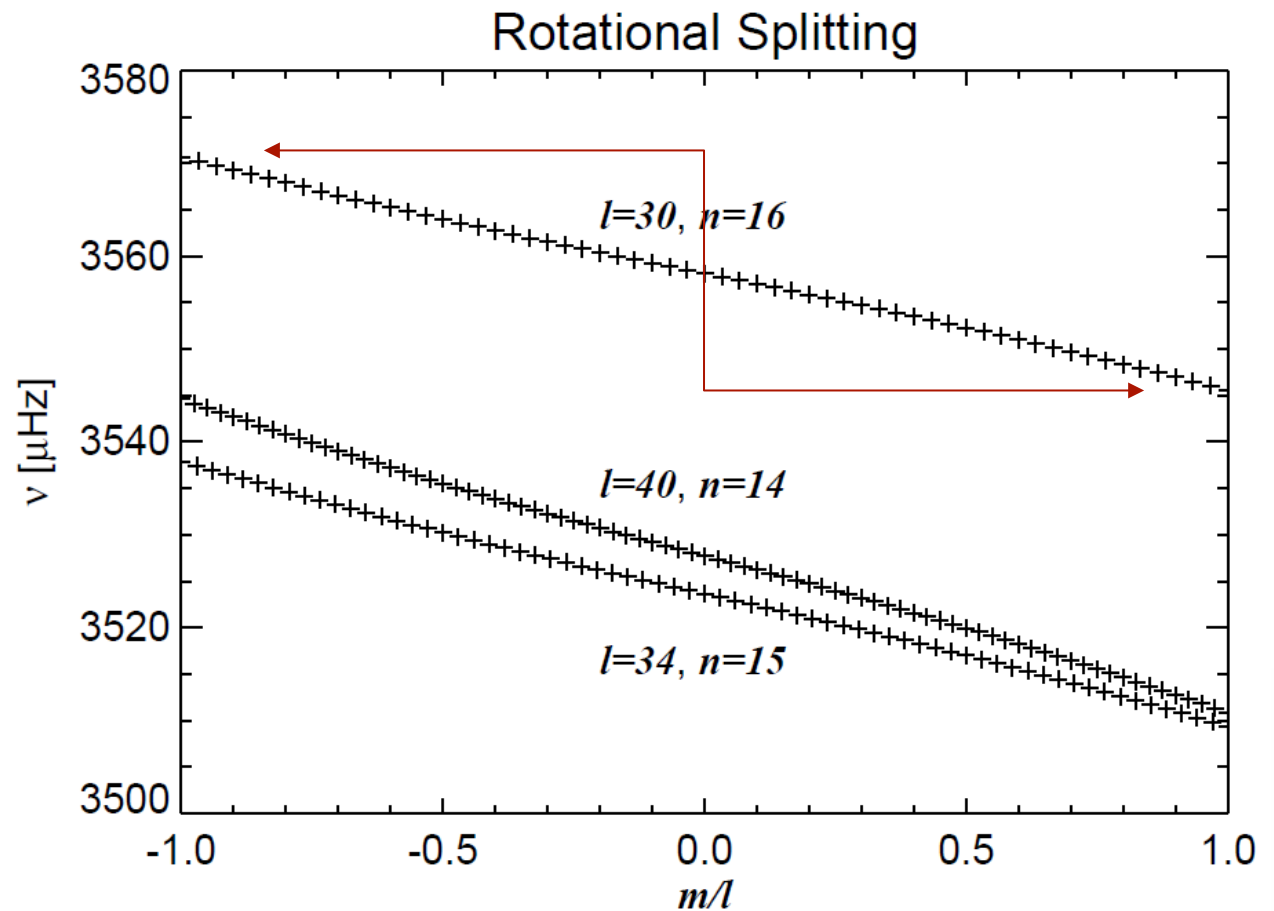
where components are expanded in terms of spherical harmonics

$$T_s^t(r; \mu; \dot{A}) = i w_s^t(r) e_r \otimes r_h Y_s^t(\mu; \dot{A})$$

$$P_s^t(r; \mu; \dot{A}) = u_s^t(r) Y_s^t(\mu; \dot{A}) e_r + v_s^t(r) r_h Y_s^t(\mu; \dot{A})$$

Lifting of degeneracies

! "Self-coupling" within
multiplets



Differential Rotation: Frequency Splitting

Toroidal Flow:
$$\mathbf{v}_0(\mathbf{r}) = \sum_s \sum_{t=-s}^s -w_s^t(r) \mathbf{e}_r \times \nabla_h Y_s^t(\theta, \phi)$$

Differential rotation (s odd, t=0) :
$$\omega_k(m) = \omega_k(m=0) + \delta\omega(m)$$

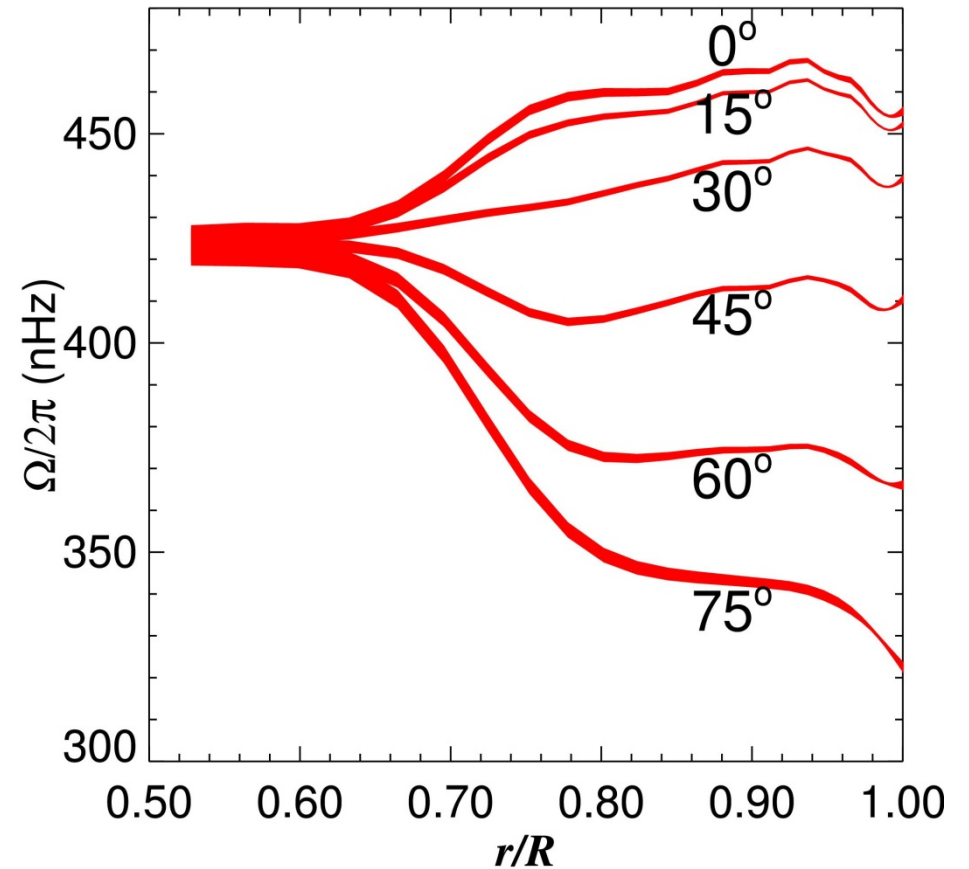
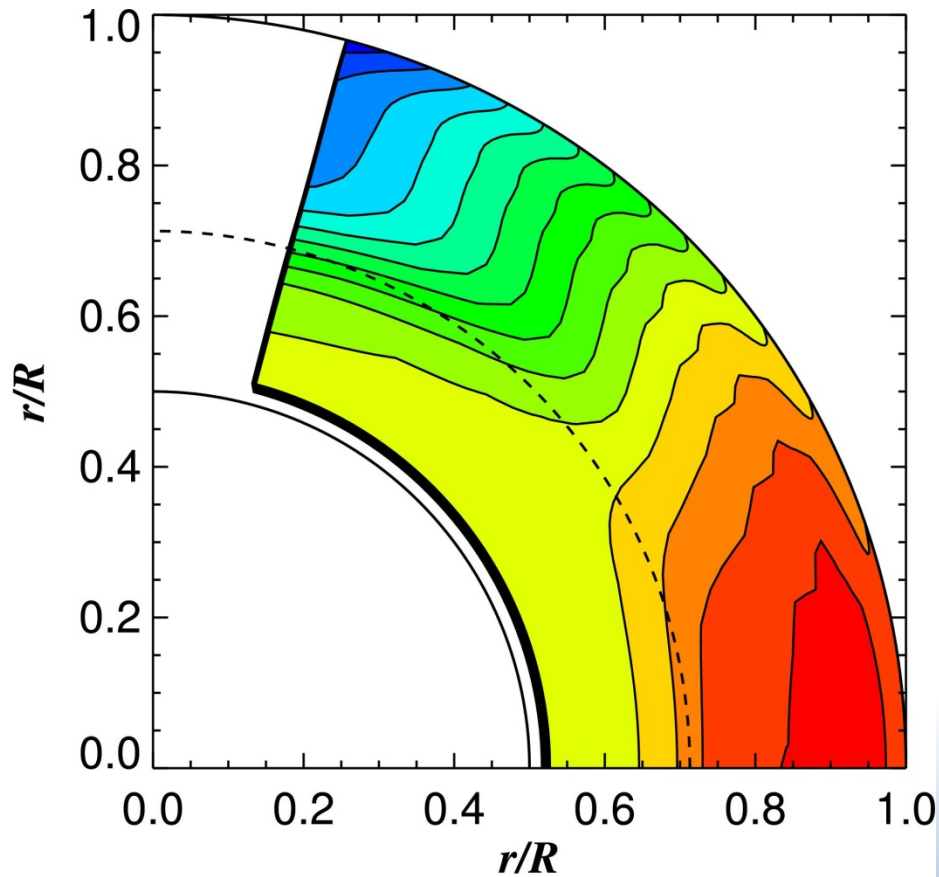
with
$$\delta\omega(m) = \sum_{s=1,3,5,\dots} c_{nl,s} \gamma_{nl,s}(m)$$

where
$$c_{nl,s} = s_0^R w_s(r) K_{nl,s}(r) r^2 dr$$

and $\gamma_{nl,s}$ orthogonal functions (Clebsch-Gordon coefficients)

! Inversion problem for $w_s(r)$

Inversion of Frequency Splittings



Poloidal Flows: Additional Frequency Shifts

Poloidal Flows:

$$\mathbf{v}_0(\mathbf{r}) = \sum_{s=0} \sum_{t=-s}^s u_s^t(r) Y_s^t(\theta, \phi) \mathbf{e}_r + v_s^t(r) \nabla_h Y_s^t(\theta, \phi)$$

$s \neq 0, t \neq 0$: giant cells

$s \neq 0, t = 0$: meridional flow

Frequency shifts:

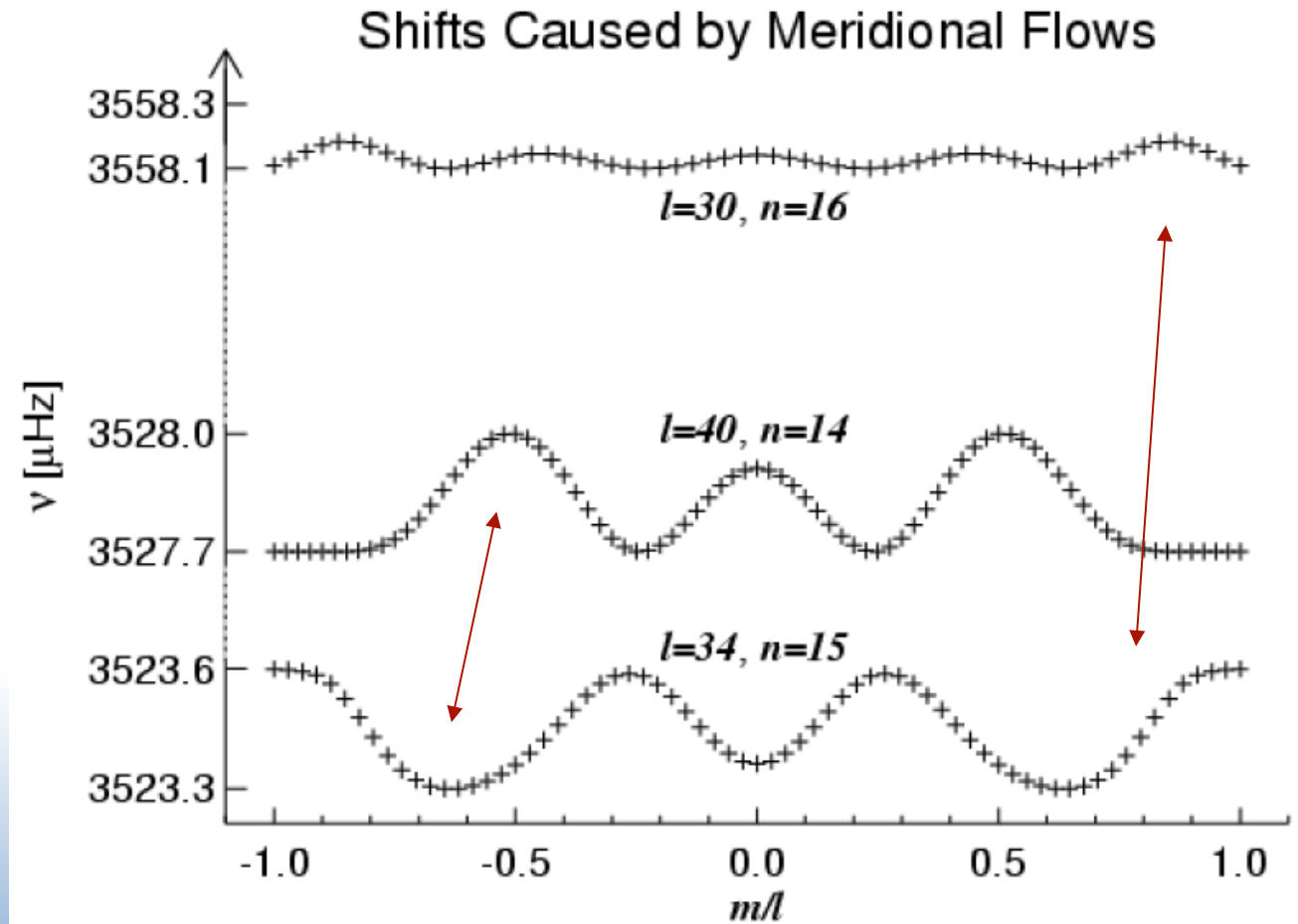
$$\delta\omega(m) = \sum_{s+l+l' \text{ even}} (c_{nn',ll',s} \gamma_{nn',ll',s}(m))^2$$

basis functions are squared Clebsch-Gordon coefficients ! **Orthogonality?**

No global helioseismology inversions with the existing frequency analysis tools

Frequency Shifts caused by Meridional Flow

- $V_{\max} = 100$ m/s
- $s=8, t=0$

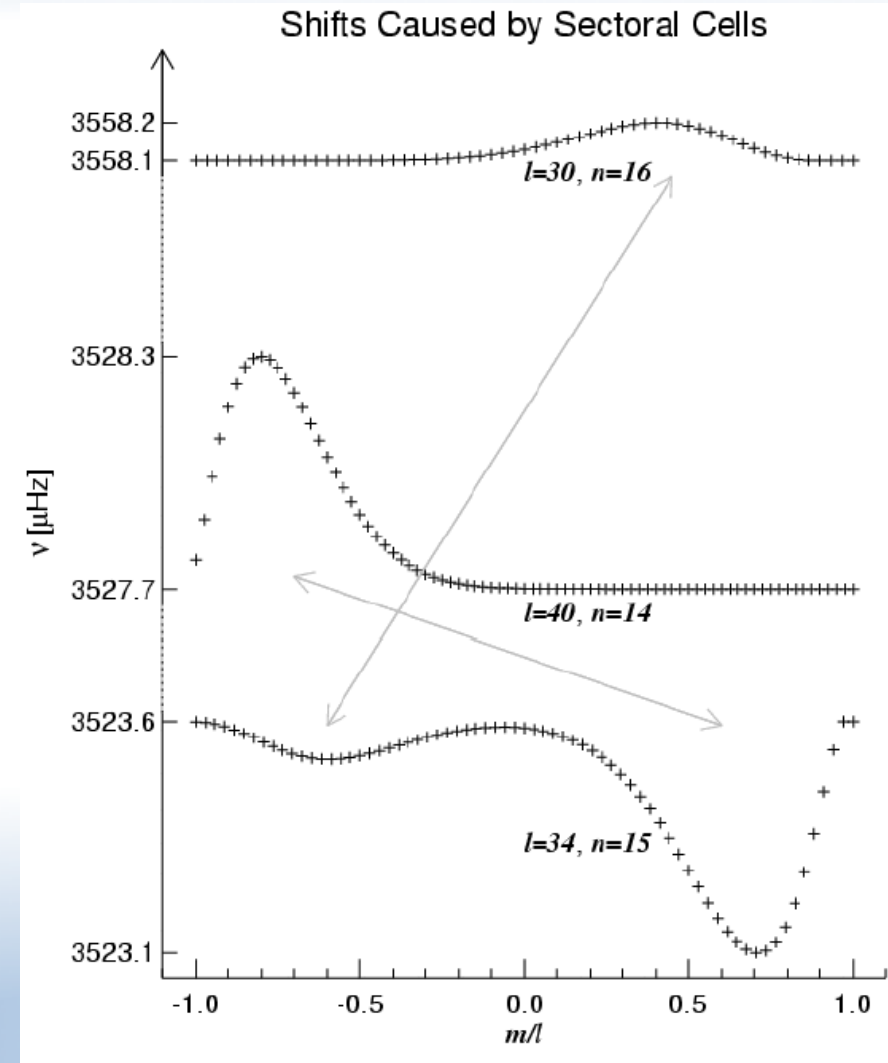


- $V_{\max} = 100$ m/s
- $s=8, t=8$

Additional frequency shifts

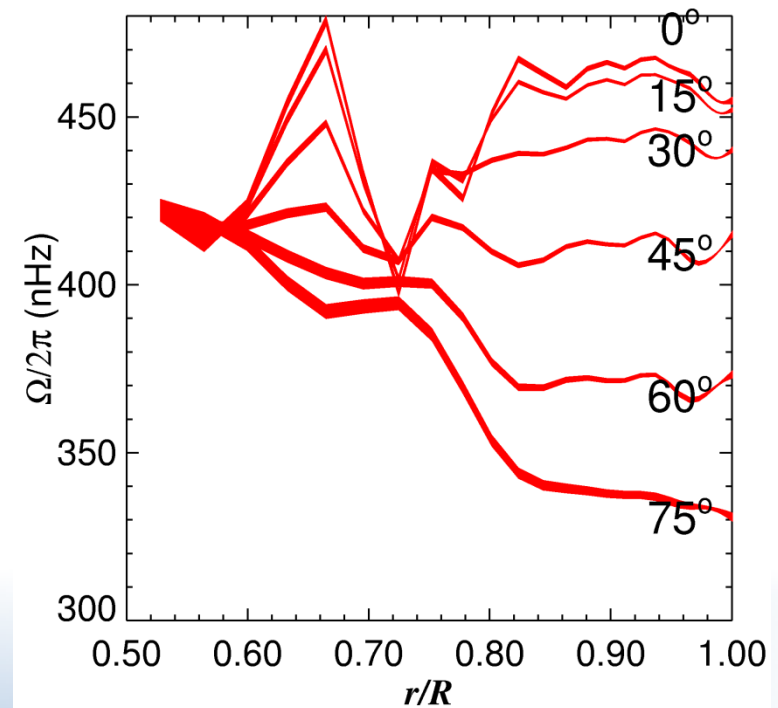
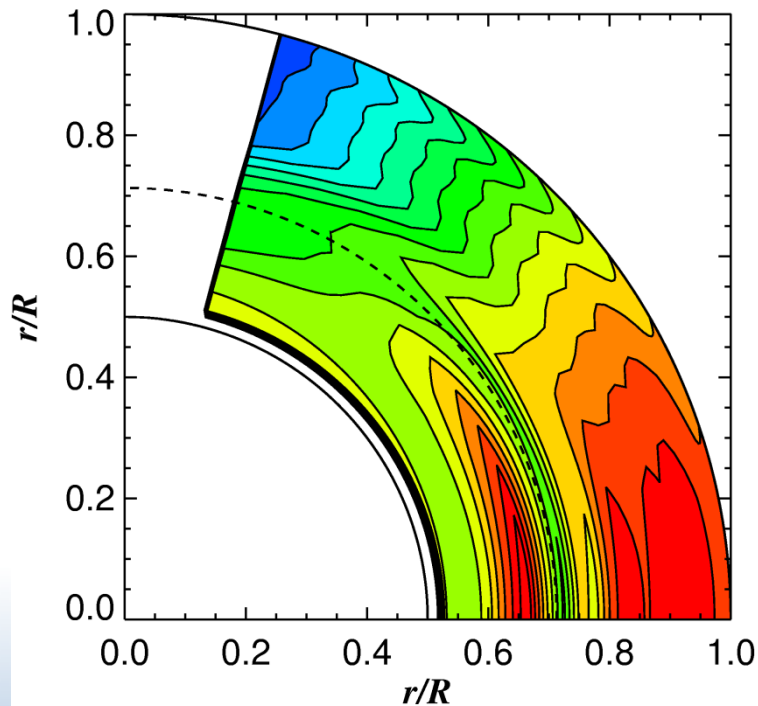
**Giant cells & meridional circulation
leave signature in global data**

Real effect is very small in comparison to
rotational splitting.



Effect of Giant Cells on Frequency Inversions for Rotation

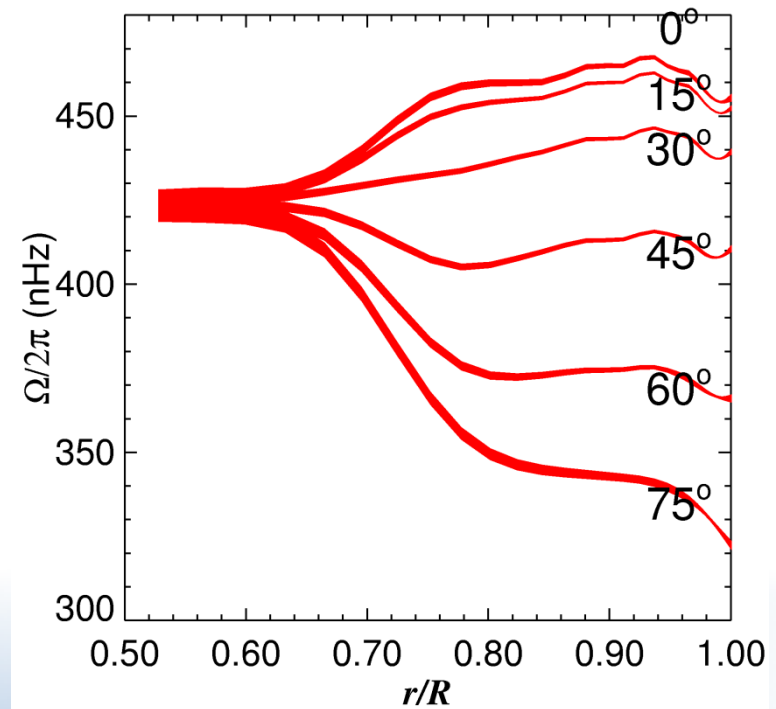
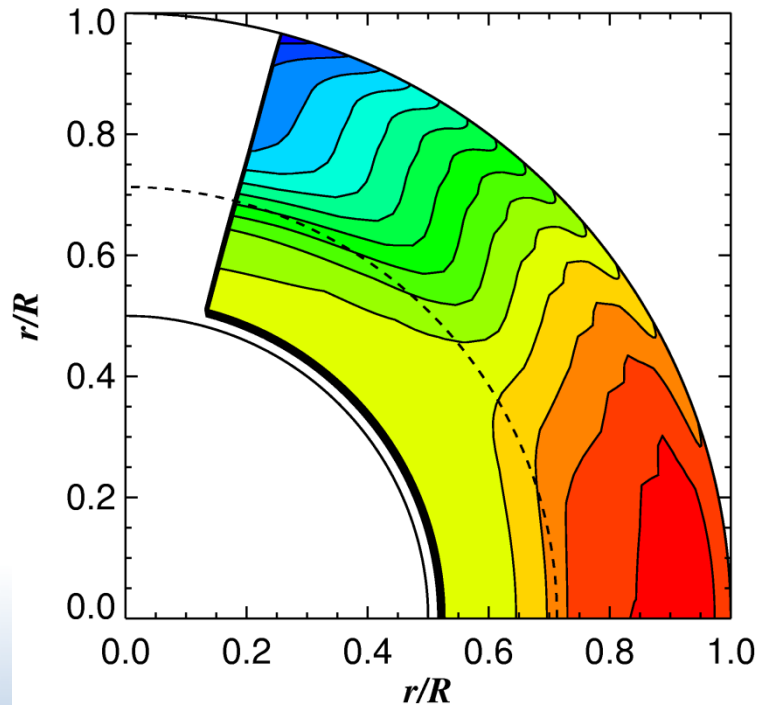
$V = 10 \text{ m/s}$



Given limits for giant cell velocities in the Sun do not cause a problem.
What about stars?

Effect of Meridional Flow on Frequency Inversions for Rotation

~~V = 000 m/s~~
V = 000 m/s



The effect of large-scale poloidal flows on solar oscillations

Perturbation theory:
**in first order no effect on the frequencies,
only in higher orders**

Resulting average effect:

Frequencies corrected down

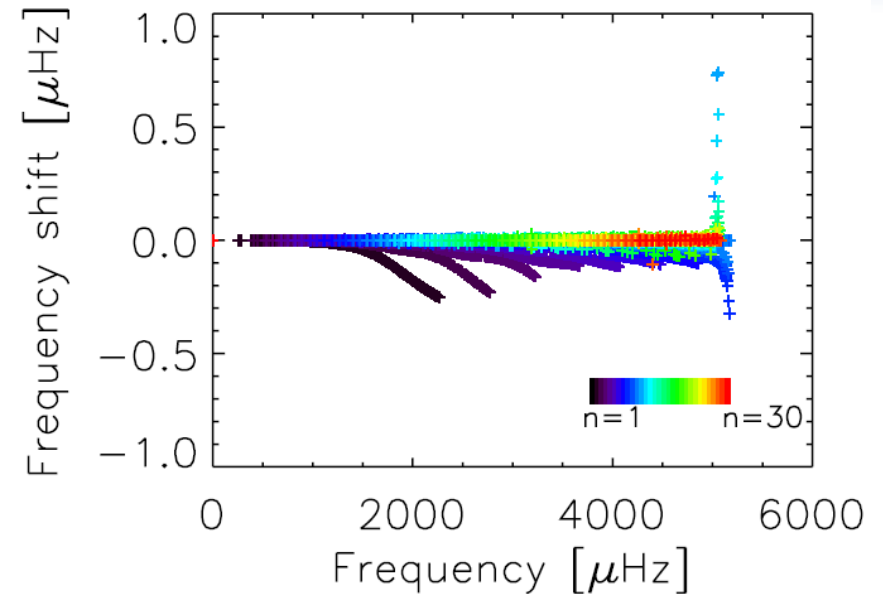
Stix & Zhugzhda 1994:

in the order to explain discrepancies between theoretical and observed frequencies

→ giant cells & meridional flow not measurable
with global helioseismology?

But: **first order perturbations in the eigenfunctions**

Observed oscillations are no pure eigenstates any more



Perturbation theory applied to the oscillation equations of p modes:

> coupling of modes in a “neighborhood” K_k of a target mode $k=(n,l,m)$

$$\mathbf{v}_k(r, \theta, \phi, t) = \underbrace{\alpha_k(t)}_{\text{osc. amplitude}} \underbrace{\boldsymbol{\xi}_k(r, \theta, \phi)}_{\text{perturbed eigenfunction}} = \alpha_k(t) \sum_{k' \in K_k} c_{kk'} \underbrace{\boldsymbol{\xi}_{k'}^0(r, \theta, \phi)}_{\text{unperturbed eigenfunction}}, \quad (\text{Lavelly \& Ritzwoller 1992})$$

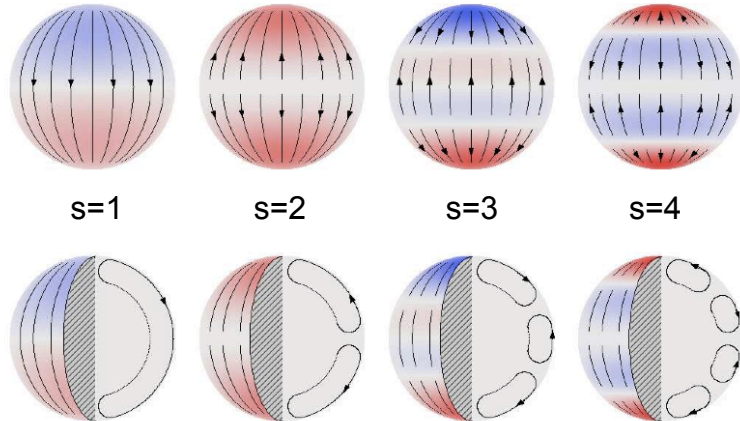
> $c_{kk'}$ – coupling coefficient, 1. order approximation

$$c_{kk'} \approx \begin{cases} 1 & \text{for } k' = k \\ \frac{H_{k'k}}{\omega_k^2 - \omega_{k'}^2} & \text{for } k' \in K \setminus \{k\} \end{cases} .$$

> $H_{k'k}$ – matrix element of mode coupling between mode k, k' (\approx coupling strength):

$$H_{k'k} = -2i\omega_{ref} \int \rho_0 \overline{\boldsymbol{\xi}_{k'}^0} \cdot \underbrace{(\mathbf{u} \cdot \nabla \boldsymbol{\xi}_k^0)}_{\text{advection of acoustic wave}} d^3\mathbf{r}$$

Global Helioseismology – Perturbation Theory



$$\mathbf{u}(\mathbf{r}) = \sum_{s=1}^{\infty} \left[\underbrace{u_s^0(r) Y_s^0(\theta, \phi) \mathbf{e}_r}_{\text{radial component}} + \underbrace{v_s^0(r) \partial_{\theta} Y_s^0(\theta, \phi) \mathbf{e}_{\theta}}_{\text{horizontal component}} \right]$$

Mass conservation: $\rho_0 r s(s+1) v_s^0 = \partial_r (r^2 \rho_0 u_s^0)$

(Figure source: D. Hathaway, NASA)

→ coupling matrix element $H_{k'k} = H_{n'l',nl}(m) = i\omega_{ref} \sum_s b_{n'l',nl}^s \mathcal{P}_{l'l}^s(m)$,

complete set of orthogonal polynomials: $\mathcal{P}_{l'l}^s(m) := (-1)^{-m} \begin{pmatrix} l' & s & l \\ -m & 0 & m \end{pmatrix}$,

b-coefficients: $b_{n'l',nl}^s = \int_0^R \rho_0(r) K_s^{n'l',nl}(r) u_s^0(r) r^2 dr$

→ knowing $b_{n'l',nl}^s$ one can infer the radial flow strength $u_s(r)$!

Observable Effect on Global Oscillations

SHT of Dopplergrams:
$$o_{l'm'}(t) = \int \overline{Y_{l'}^{m'}}(\theta, \phi) W(\theta, \phi) v_D(\theta, \phi, t) d\Omega,$$

FT – frequency domain \Downarrow
$$= \sum_k \alpha_k(t) \sum_{k'' \in K_k} c_{kk''} L_{k'k''}, \quad L_{k'k''} \text{ - leakage matrix}$$

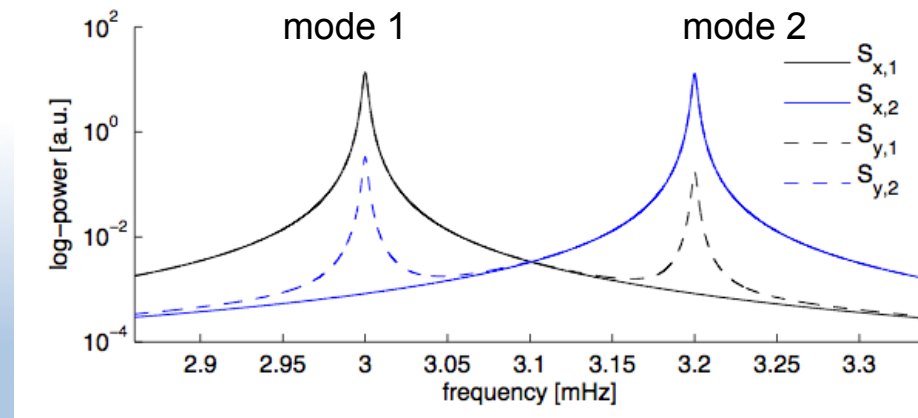
$\tilde{o}_{l'm'}(\omega)$

Observable: amplitude ratio between target mode $k=(n,l,m)$ and coupling modes k' in K_k :

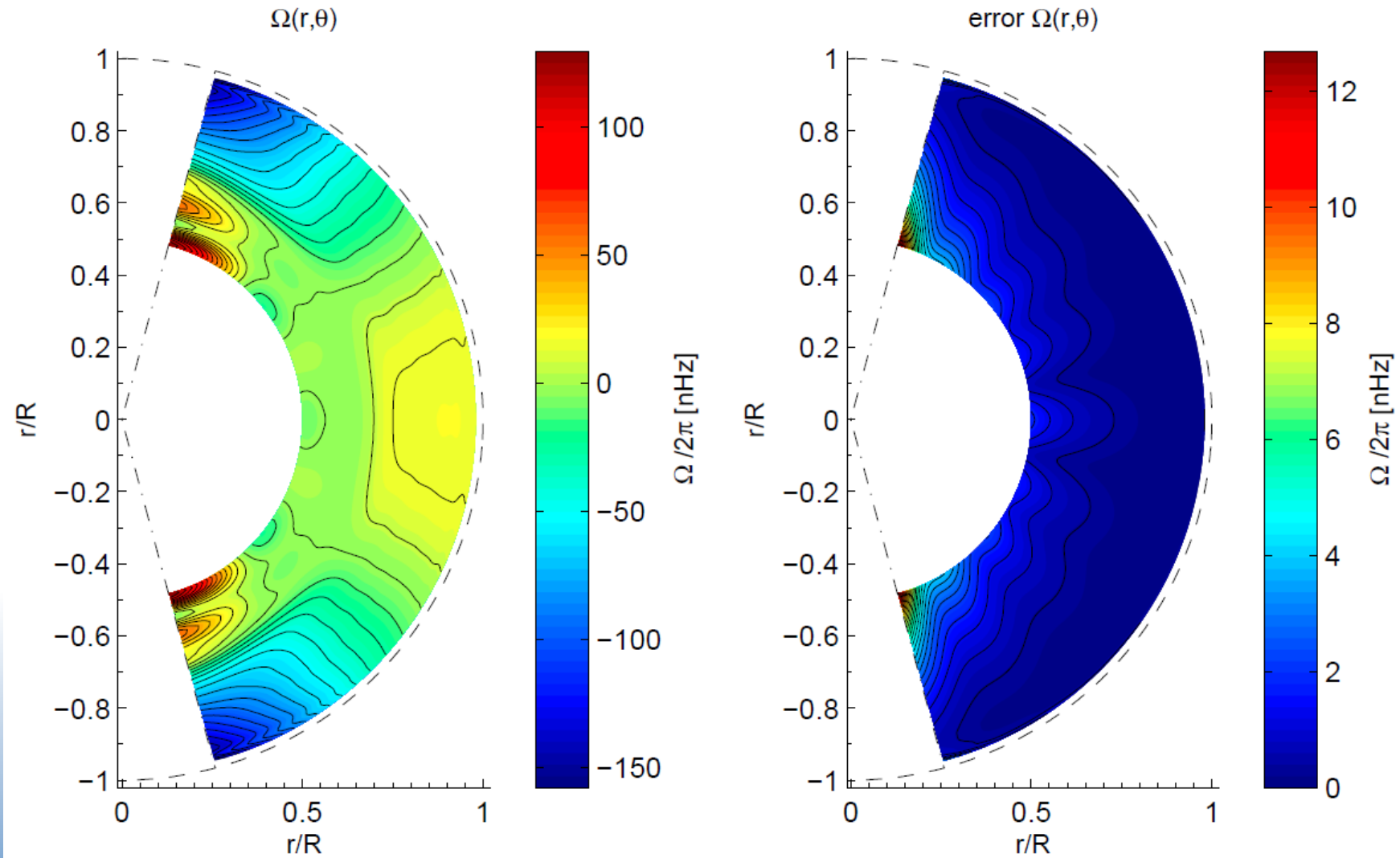
$$y_{lm'l'm}(\omega_{nlm}) := \frac{\tilde{o}_{l'm}(\omega_{nlm})}{\tilde{o}_{lm}(\omega_{nlm})} \approx \frac{\sum_{k'' \in K_k} c_{kk''} L_{k'k''}}{\sum_{k'' \in K_k} c_{kk''} L_{kk''}} \in \mathbb{C}$$

target mode

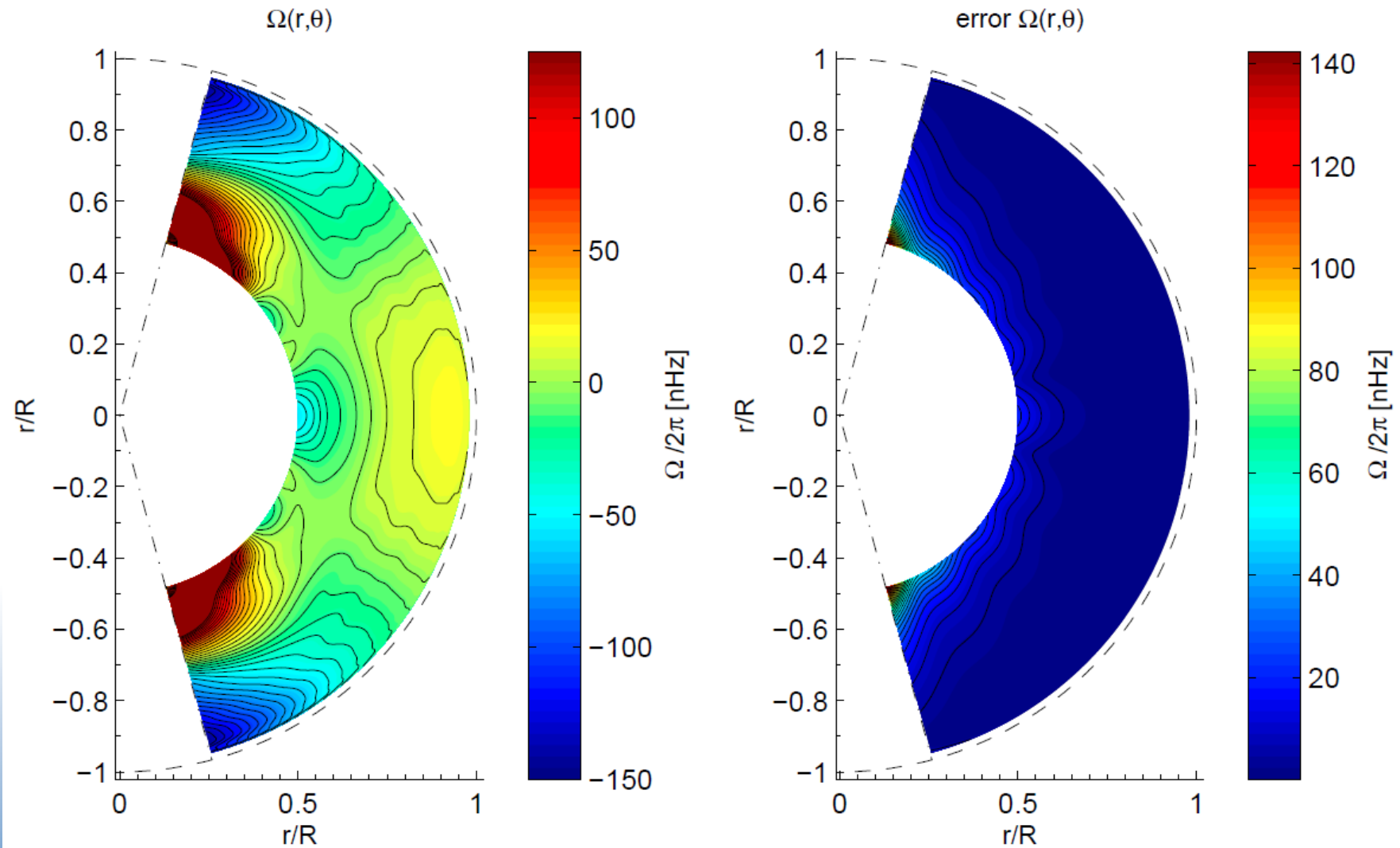
„ratio between global oscillations in the frequency domain that corresponds to coupling modes“



Inversion for Rotation from Frequency Splittings



Inversion for Rotation from Amplitude Ratios



Operation scheme:

- estimation of amplitude ratios $y_{kk'}$ for the multiplets > cross-spectral analysis / gain
- estimation of b-coefficients $b_{kk'}^s$ > least squares fitting routine
- inversion of radial flow strength $u_s(r)$ > e.g. SOLA inversion method (as used for diff. rotation)
- reconstruction of the horizontal flow strength $v_s(r)$ from $u_s(r)$ > e.g. polynomial fit

$$\rho_0 r s(s+1) v_s^0 = \partial_r (r^2 \rho_0 u_s^0)$$

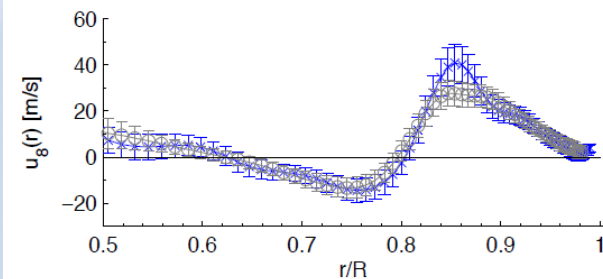
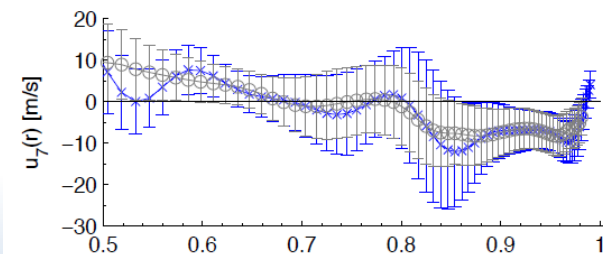
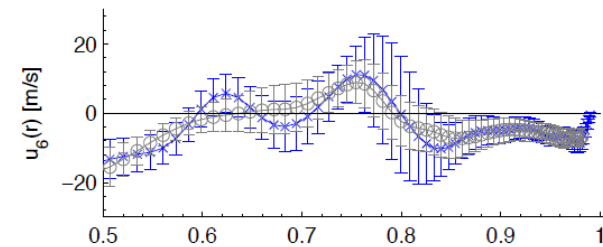
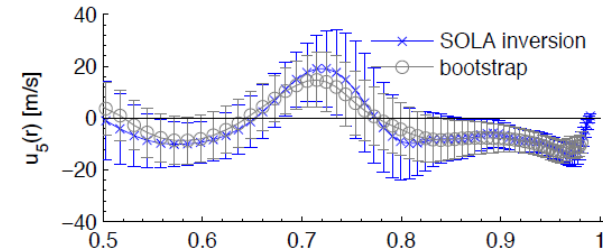
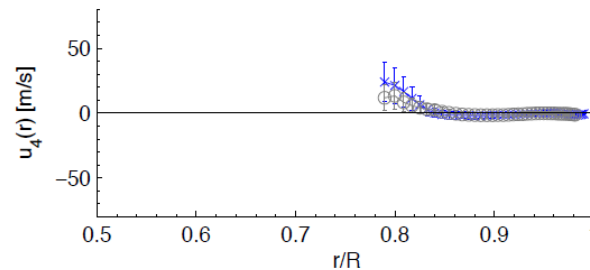
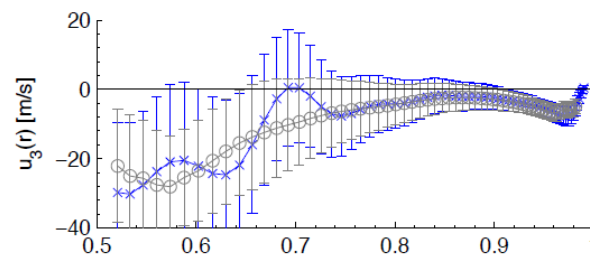
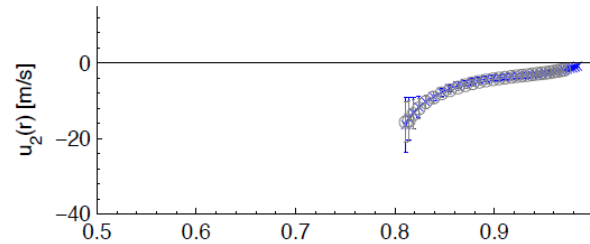
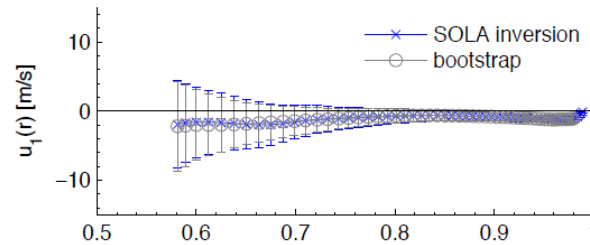
Application to MDI data

MDI data: 2004 – 2010
Harmonic degree:
 $1 \leq l \leq 200$

Radial flow for different
components

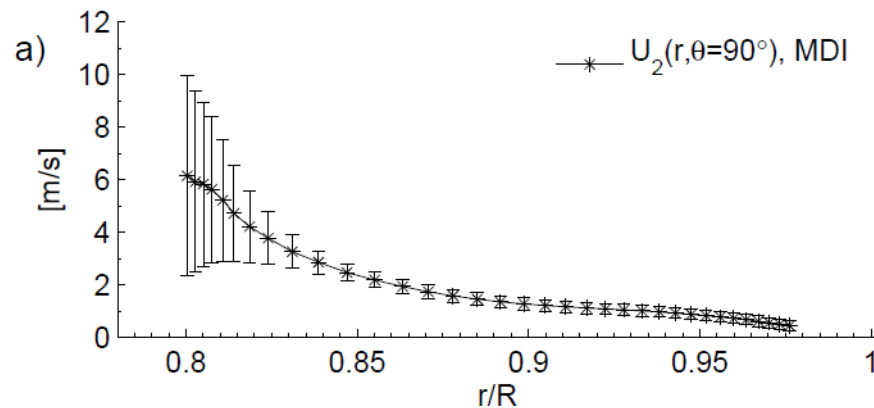
Deepest depth depends
on degree s
(between 0.5 and 0.8 R)

$s=2$ and $s=8$
to be studied in detail

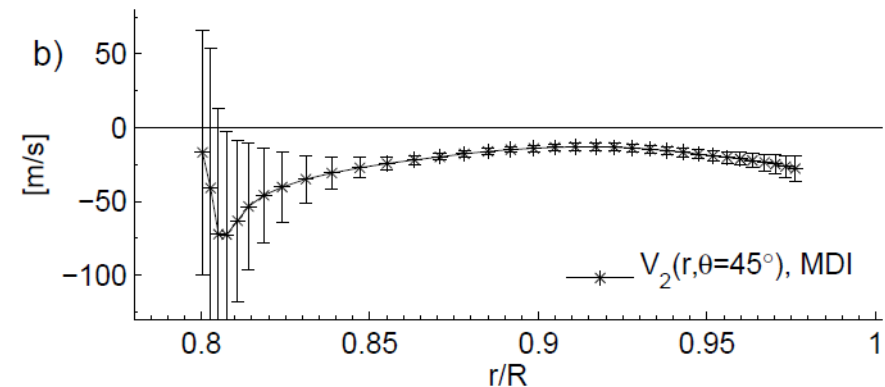


Application to MDI Data – Results for $s=2$

Radial flow strength U_2



Horizontal flow strength V_2

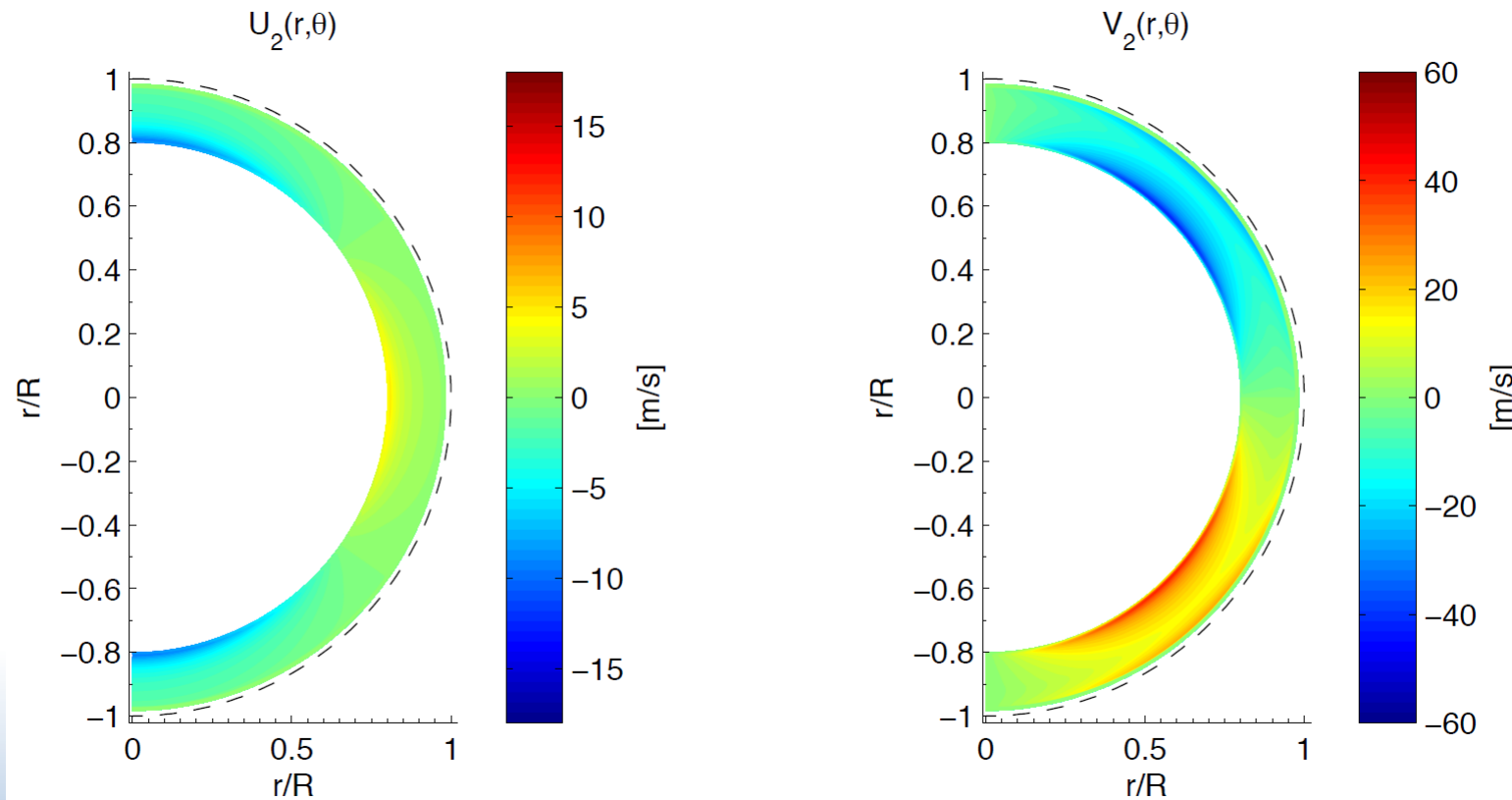


- > inversions for $0.82 \leq r/R \leq 0.97$ (depth 20.9 – 125.3 Mm)
- > u_2 grows \approx linear with depth
- > $v_2 \approx$ constant with r
- > 1σ -error u_2 : 0.6 – 12 m/s
- > 1σ -error v_2 : 2.1 – 100 m/s

> **no return flow within $0.82 < r/R < 0.97$!**

Application to MDI Data – Results for $s=2$

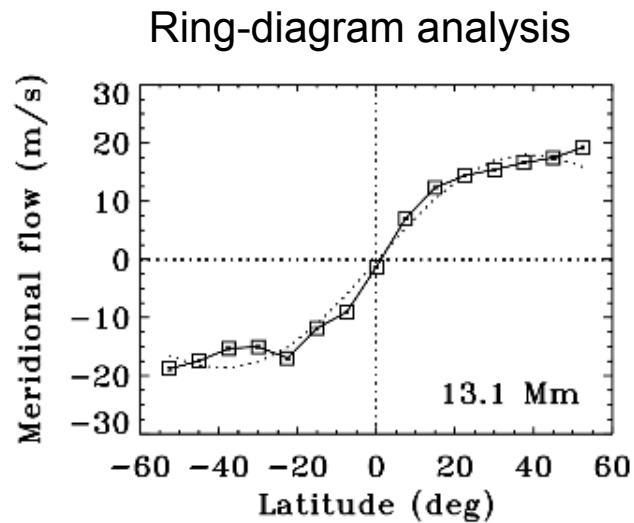
Cross-section through radial & horizontal flow profile with latitude ($s=2$):



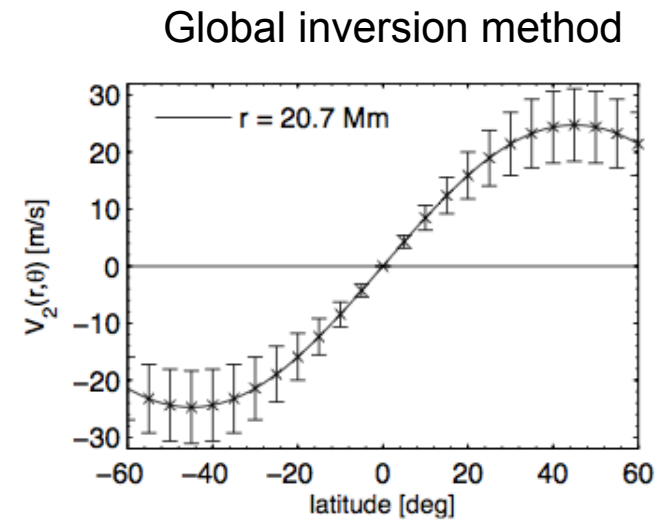
→ poleward directed horizontal flow

Comparison with Ring Diagram Analysis (for $s=2$)

Horizontal flow component of $s=2$ at 20.7 Mm depth:



(Komm et al., 2005, ApJ)

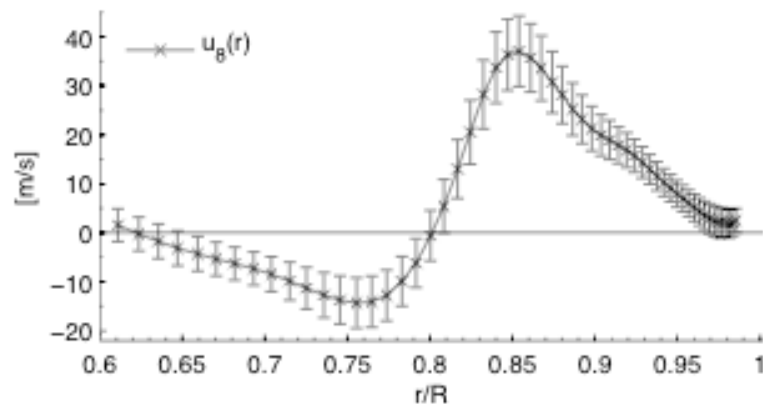


→ max. horizontal flow near surface at $\theta=45^\circ$: $V_2 \approx 28 \pm 9$ m/s

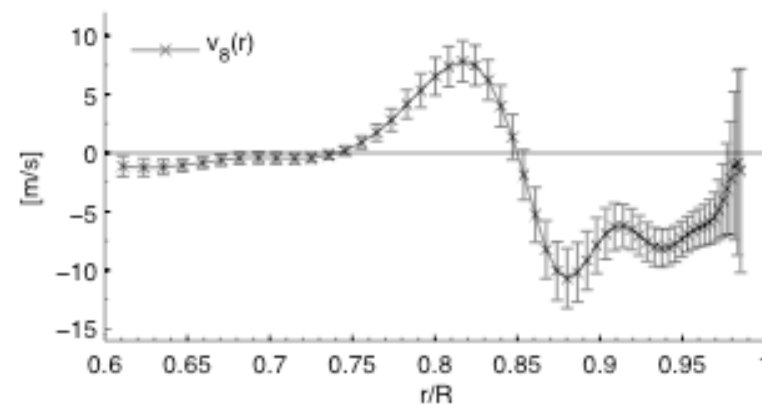
(Schad et al., 2013, ApJL)

Application to MDI Data – Results for $s=8$

radial flow strength u_8



horizontal flow strength v_8



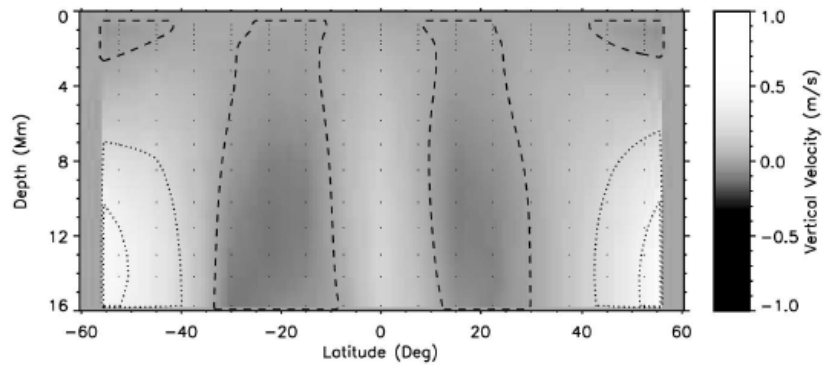
- inversions for $0.61 \leq r/R \leq 0.984$ (depth 11 – 271 Mm)
- two flow cells in depth

Comparison with Ring Diagram Analysis (for $s=8$)

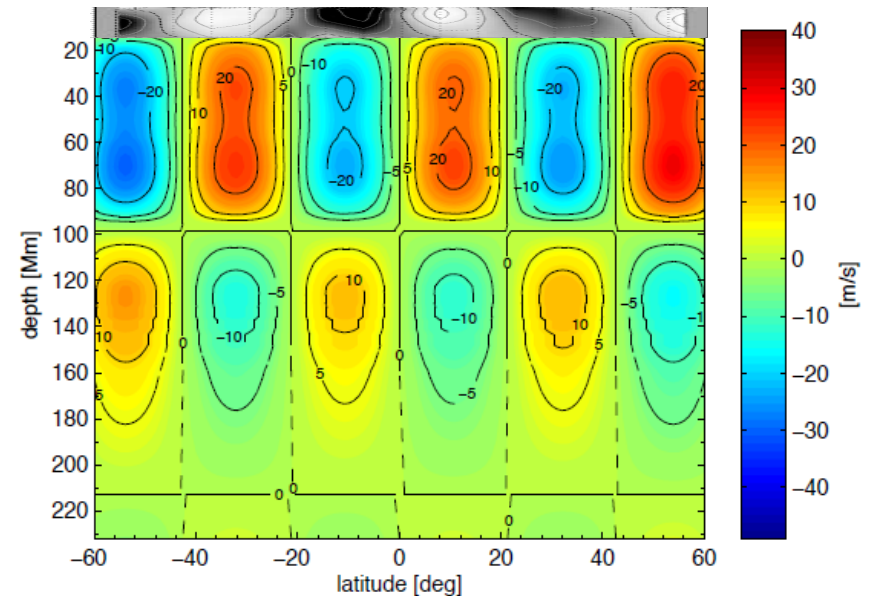
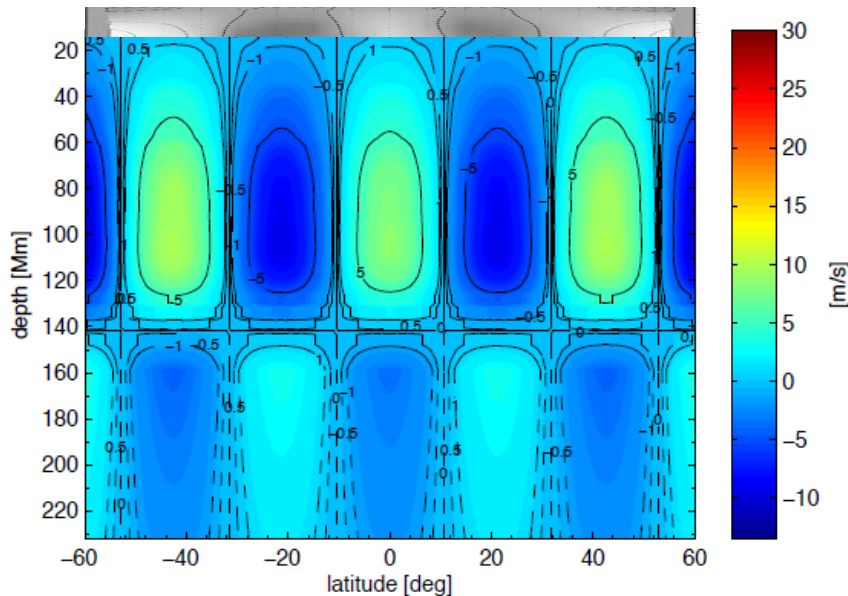
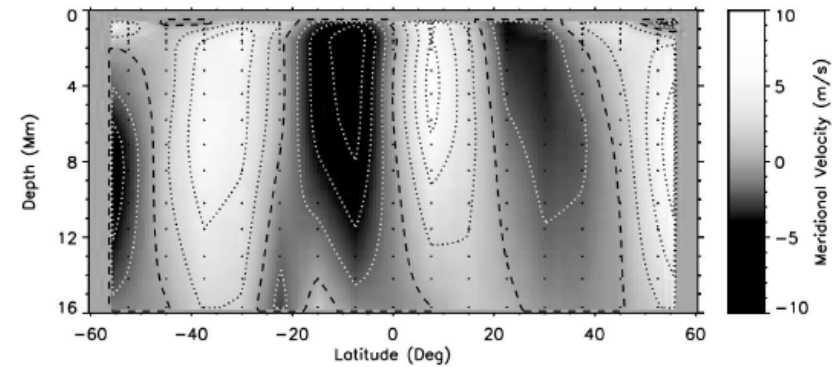
Ring-diagram analysis
(Komm et al., 2005, ApJ)

Global inversion method
(Schad et al., in prep.)

Radial flow



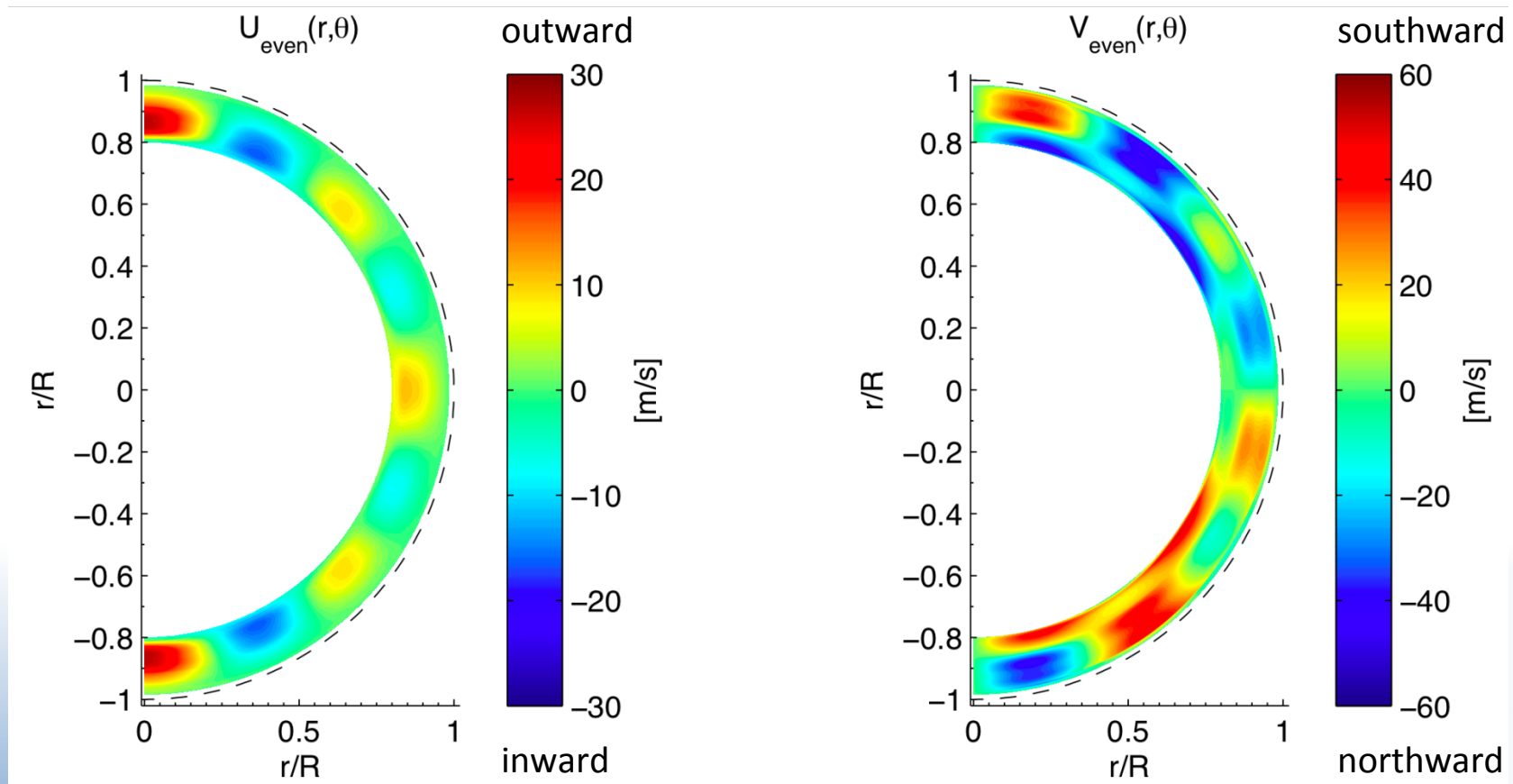
Horizontal flow



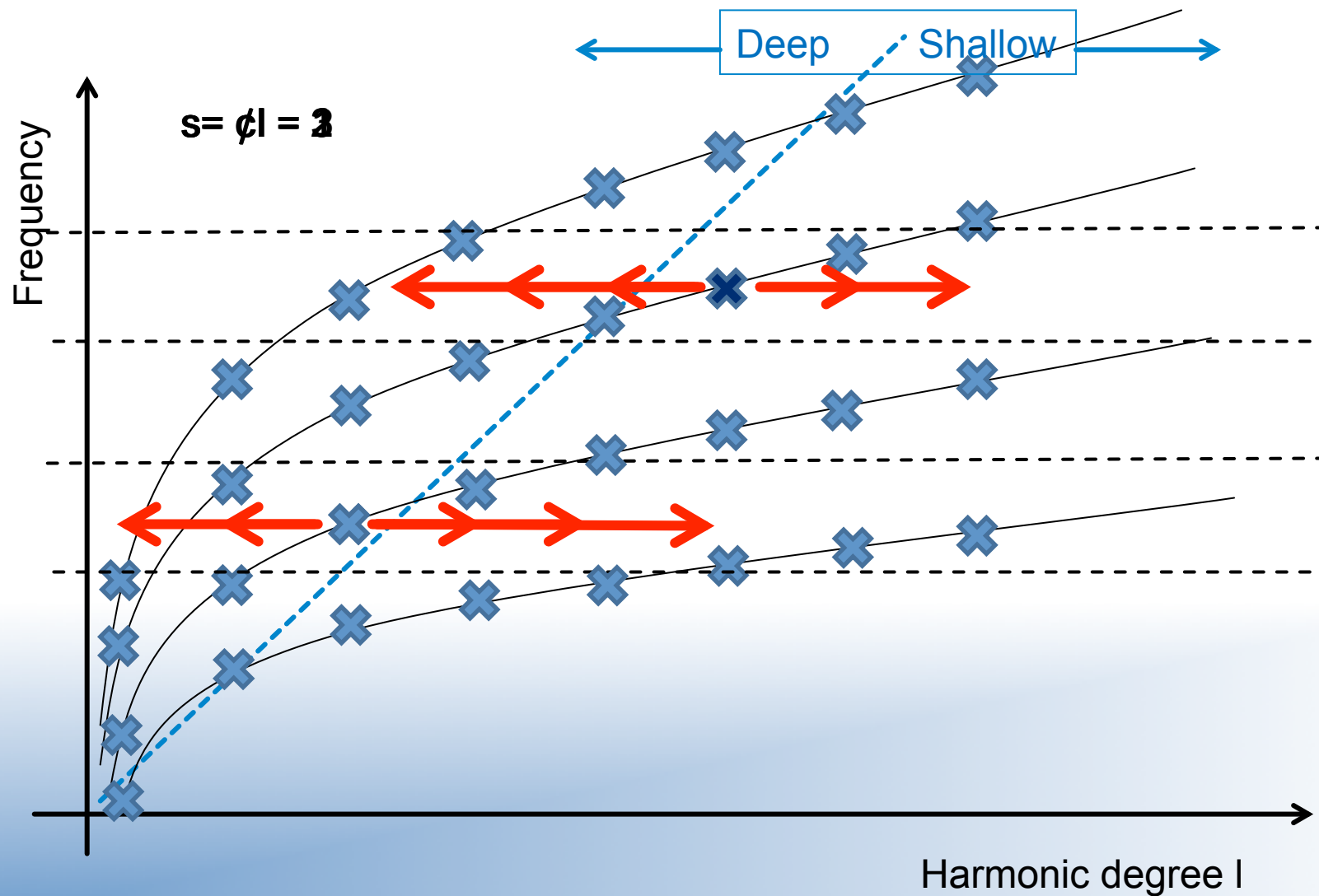
Ring-diagram analysis (0.6–16 Mm depth; 1 year average);
Global method (13–271 Mm depth; 6 years average)

Result of banded flow pattern?

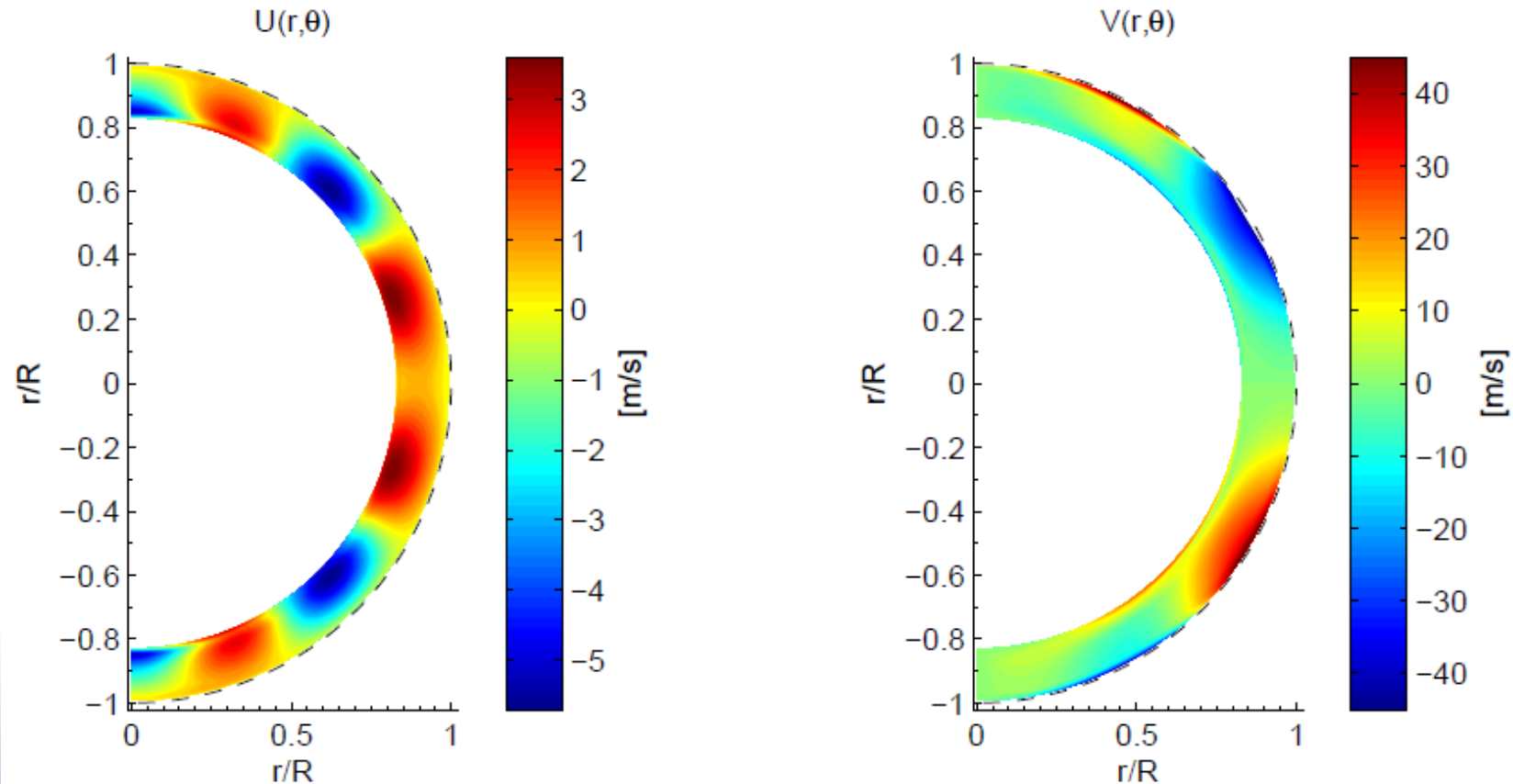
Overall Result: Inversion result including all significant flow components in a depth range of 13 – 141 Mm



Clear evidence for multiple cells in depth and latitude



Ignoring Leakage



Knowledge about leakage matrix is crucial.
Ignoring leakage results in plausible but systematically wrong results

- Cross-spectral analysis is a new global helioseismic tool
 - Possible on the Sun (not necessarily on stars) because of long time series needed
- Measurements of flows in the solar interior:
 - Rotation
 - Meridional Flow
- Important Input:
 - Leakage Matrices -> to be part of the peak-bagging tool?
- To Do: Studying systematic effects
 - Leakage Matrices
 - Center-to-limb variation of MDI line ($\bar{m}_{nl}(z)$; eigenfunctions needed)