

# Magnetic reconnection in twisted magnetic fields in solar flares - heating, particle acceleration and observational signatures

**Philippa Browning and Mykola Gordovskyy**  
Jodrell Bank Centre for Astrophysics  
University of Manchester, UK

**Eduard Kontar (Glasgow)**  
**Alexei Kutnetsov (Irkutsk)**  
**Rui Pinto & Nicole Vilmer (Paris)**  
**Alan Hood (St Andrews)**  
**Asad Hussain (Manchester)**



- **Modelling energy release in unstable twisted loops**
- **Observational signatures**
  - Thermal emission and Hard X-rays
  - Turbulent velocities
  - Radio/microwave
- **Energy release in interacting twisted loops**

# Twisted flux ropes in the solar corona

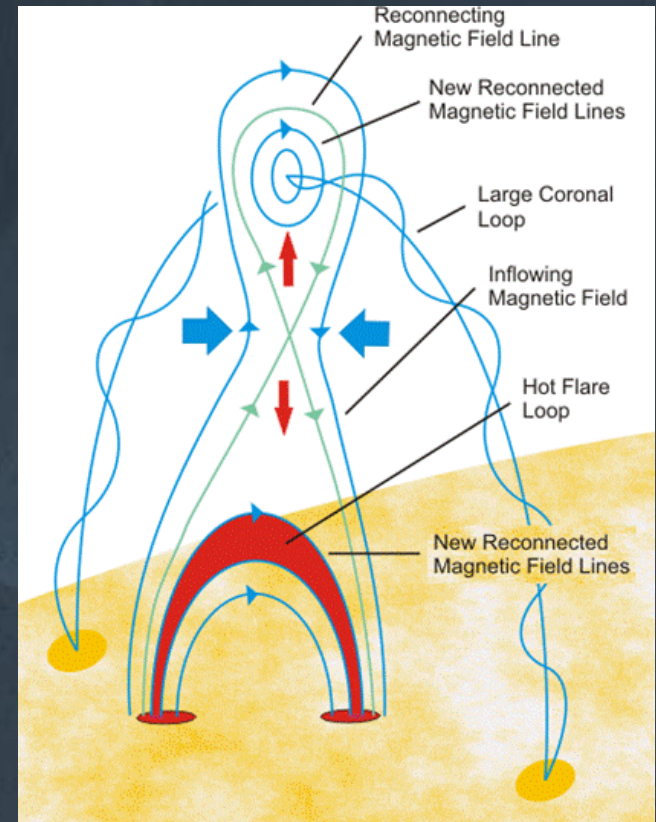
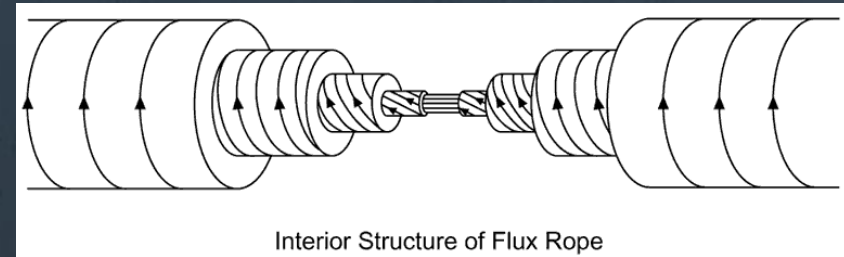
- Twisted magnetic flux ropes are non-potential fields providing reservoirs of free magnetic energy
- Unstable twisted coronal loops can be a good alternative to the standard model, particularly in smaller flares with isolated loops

*e.g. Aschwanden et al. 2009*

- In standard model, flux ropes can form due to magnetic island formation in a reconnecting current sheet with guide field *e.g. Gordovsky et al. 2010*

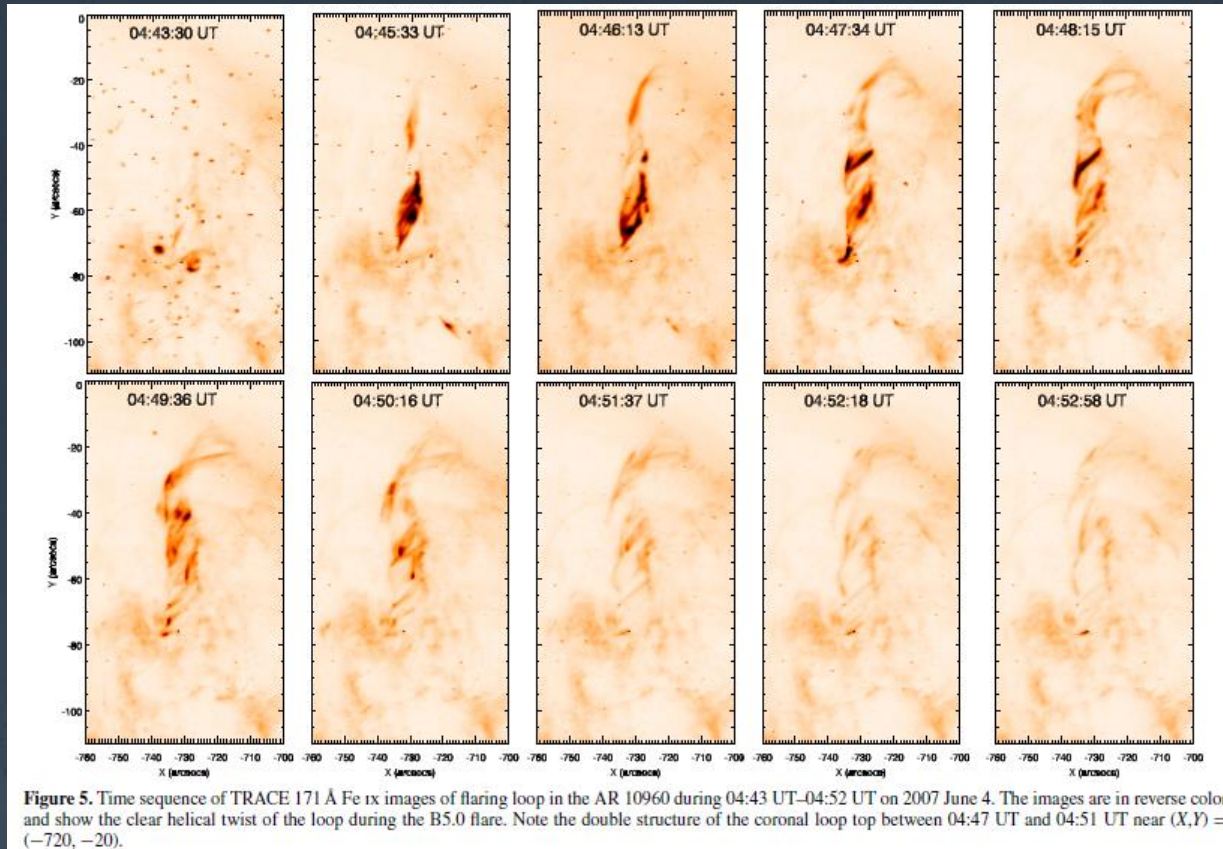
- Twist can be produced by photospheric rotation/shear motions *e.g. Brown et al. 2003*

- Newly emerging flux is expected to have some twist *e.g. simulations by Archontis and Hood*



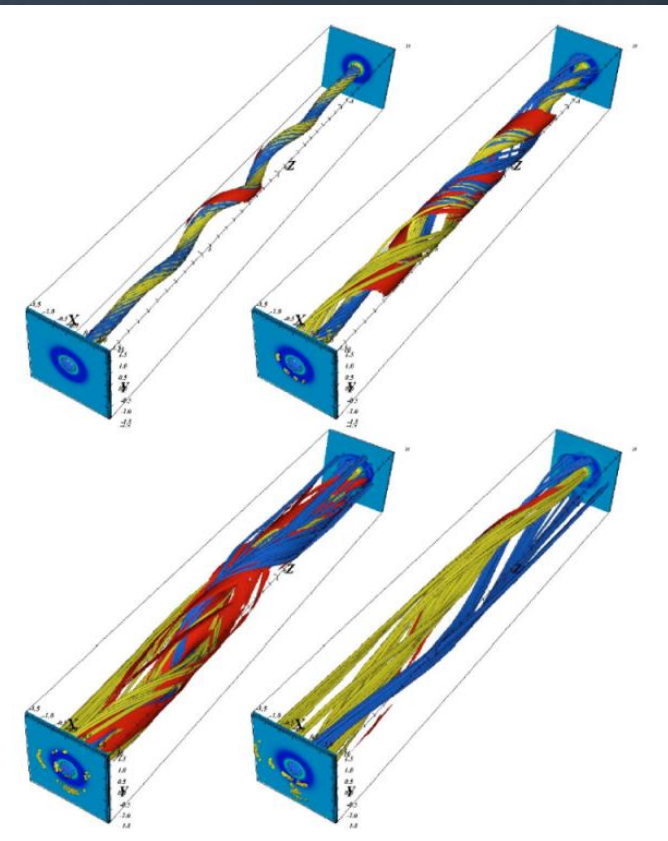
# Twisted loops in flares

from Shrivastava et al 2010 ApJ



# Kink instability modelling

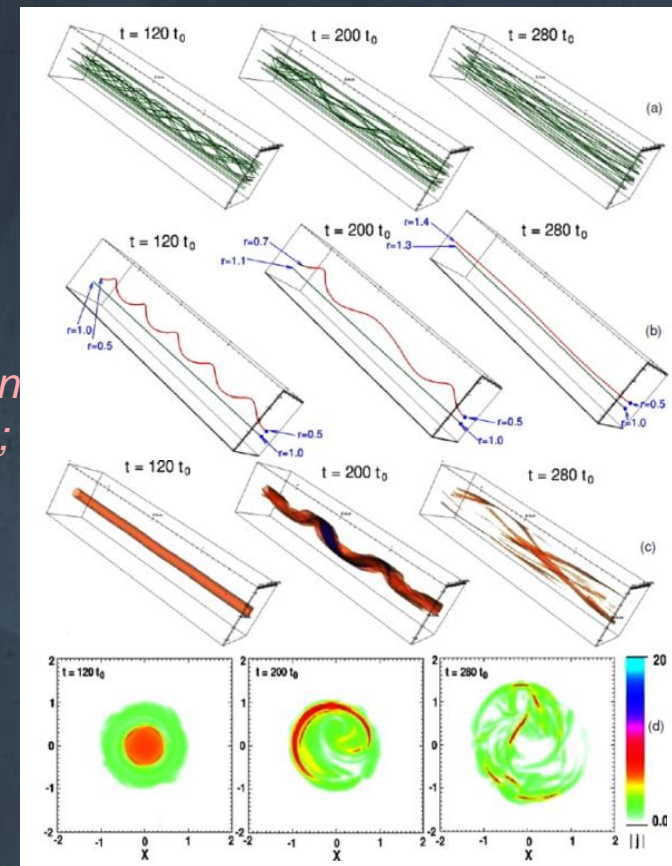
- Ideal kink instability → fragmented current sheet in nonlinear phase
- Internal reconnection and reconnection with untwisted ambient field → untwisting of field, loop expansion
- Fast magnetic reconnection dissipates magnetic energy
- Particle acceleration throughout loop volume



from Hood et al. 2009 A&A

## Kink instability in unstable cylindrical flux ropes

*Baty & Heyvaerts 1996;  
Browning & Van der Linden 2003; Browning et al 2008;  
Hood et al 2009; Botha et al, 2012; Bareford et al 2013, Bareford & Hood 2015*

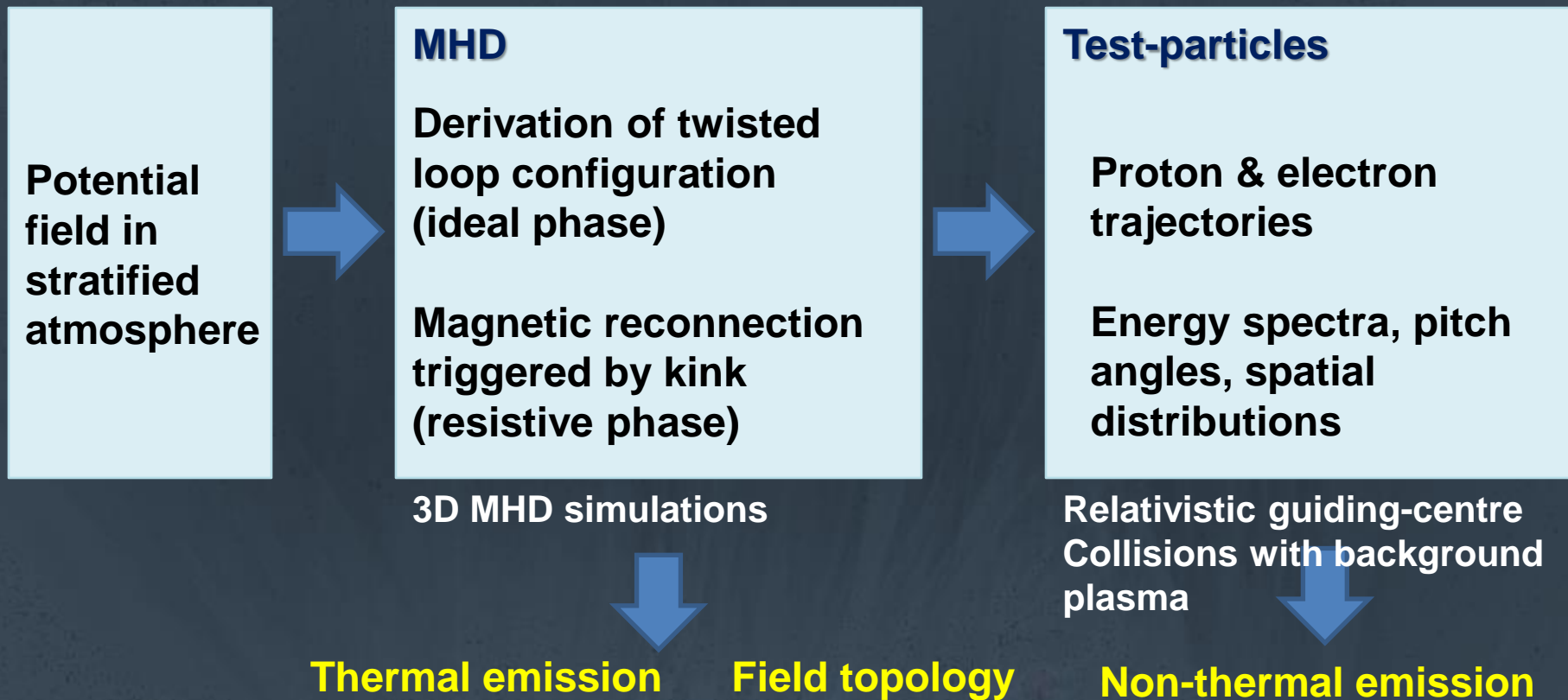


from Gordovskyy & Browning 2011 ApJ

# Observational detection of twisted loops

- Thermal emission (UV & soft X-ray, continuum and spectral lines)?
- Thermal radio?
- Non-thermal Hard X-ray?
- Non-thermal radio?

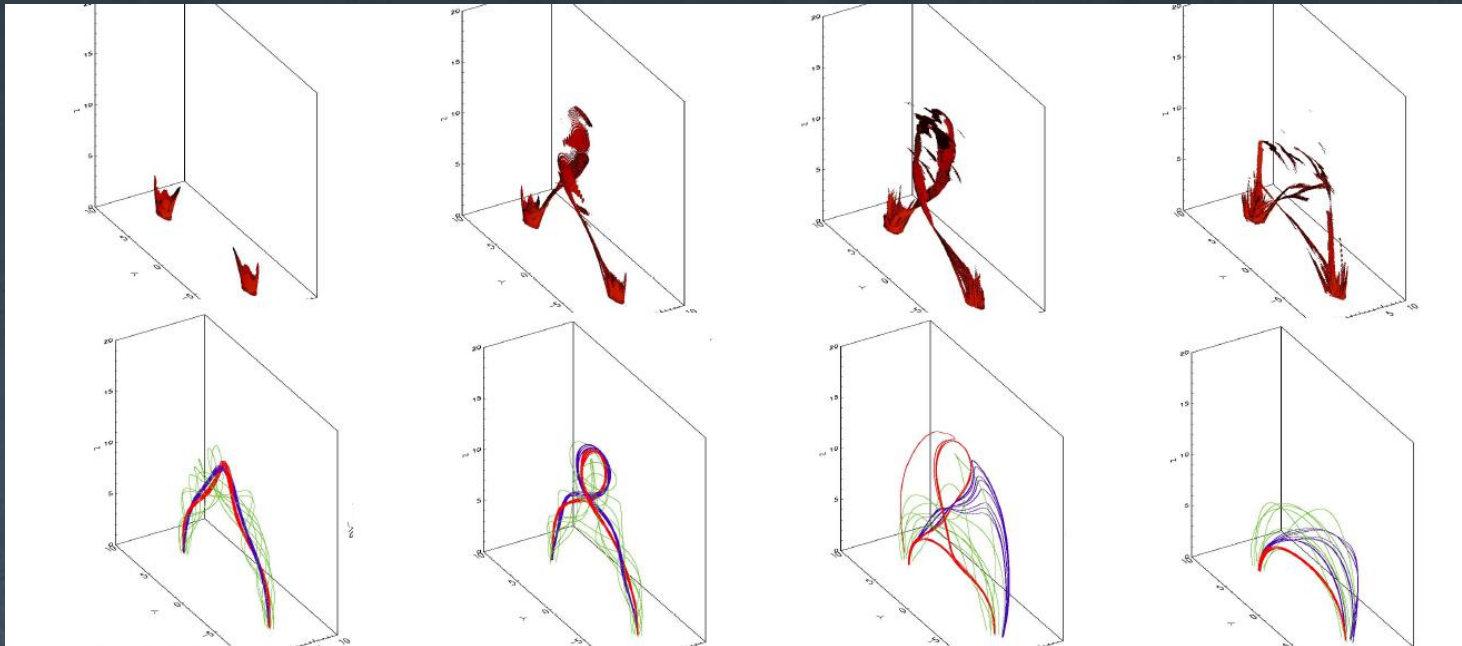
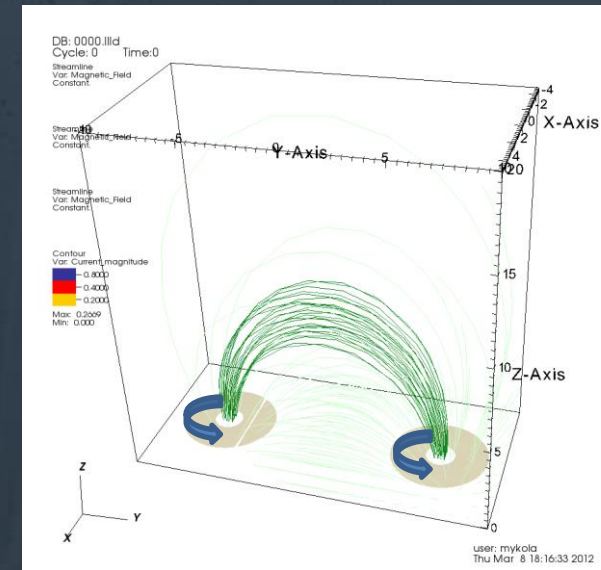
# Coupled MHD and test-particle models



- *Gordovskyy et al A&A 2014* (loop evolution, non-thermal particles & HXR)
- *Bareford et al Sol Phys 2016* (loop geometry effects)
- *Pinto et al A&A 2016* (thermal SXR continuum, non-thermal HXR)
- *Gordovskyy A&A 2016* (EUV lines – non-thermal broadening, shifts)
- *Gordovskyy et al 2017* (thermal and non-thermal microwave)

# Our model

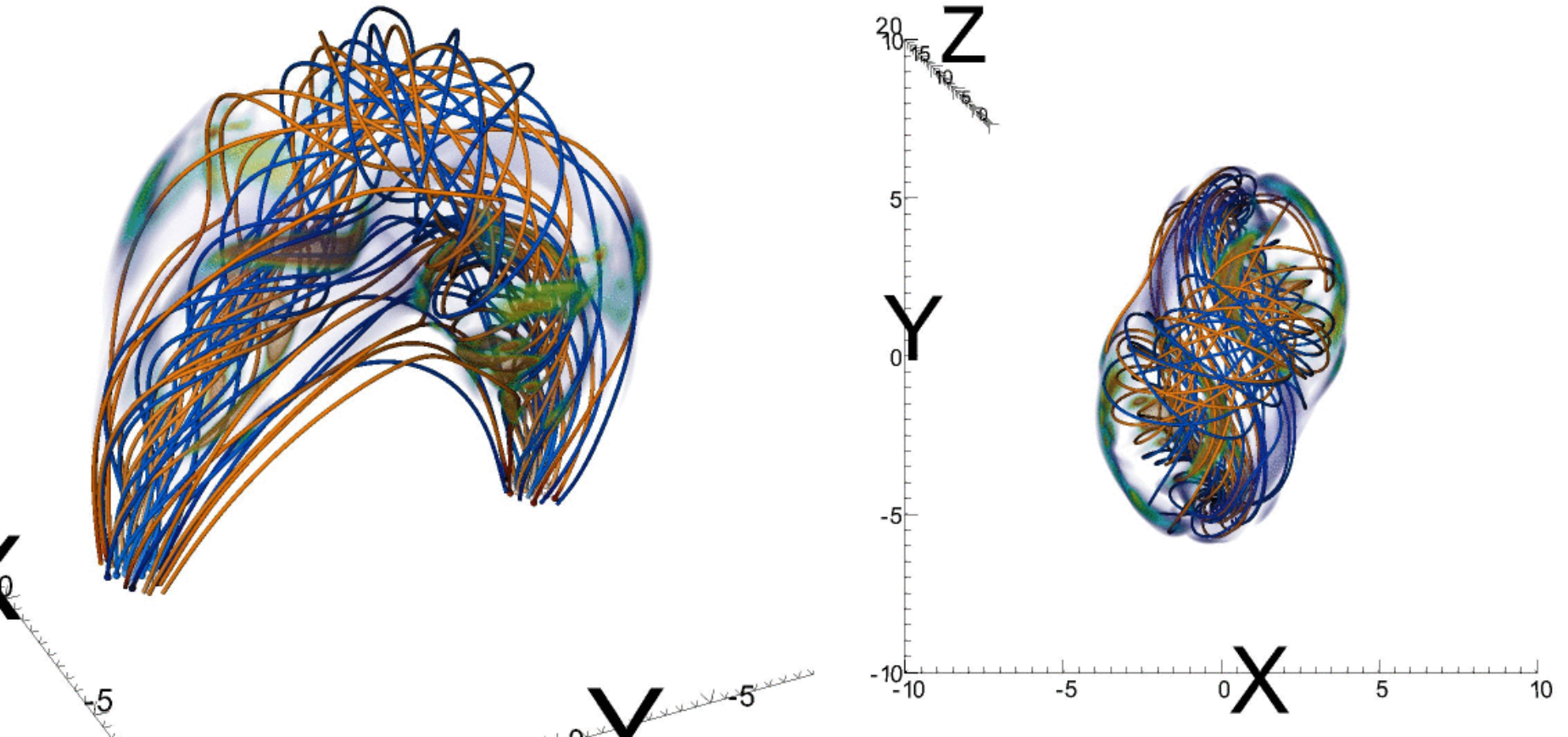
- 3D MHD with LARE3D (*Arber et al 2001*)
- anomalous resistivity triggered by ion acoustic instability
- Gravitationally-stratified atmosphere
- Bipolar magnetic region – localised rotational motions → twisted loop
- Kink instability leads to loop expansion, fragmented currents within loop, large-scale currents and reconnection





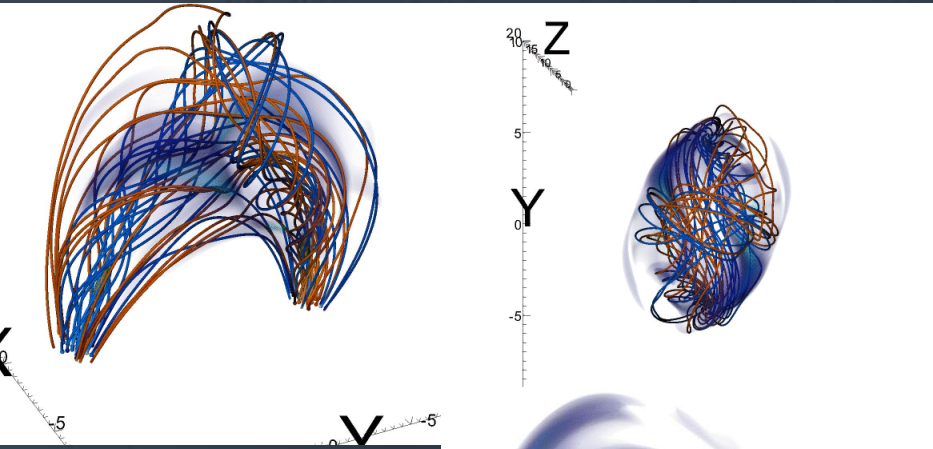
# SXR thermal emission

Small loop (model V), 2keV continuum emission

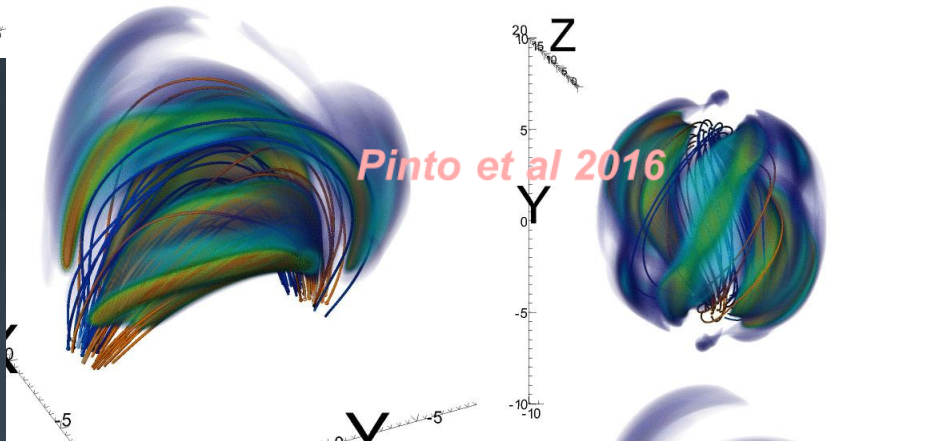


*Pinto et al 2016*

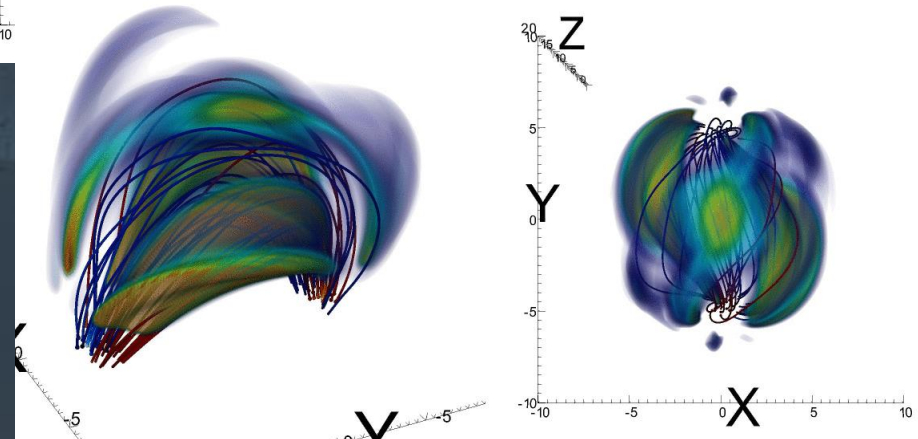
# SXR thermal emission



Emission shows only weak signature of twist



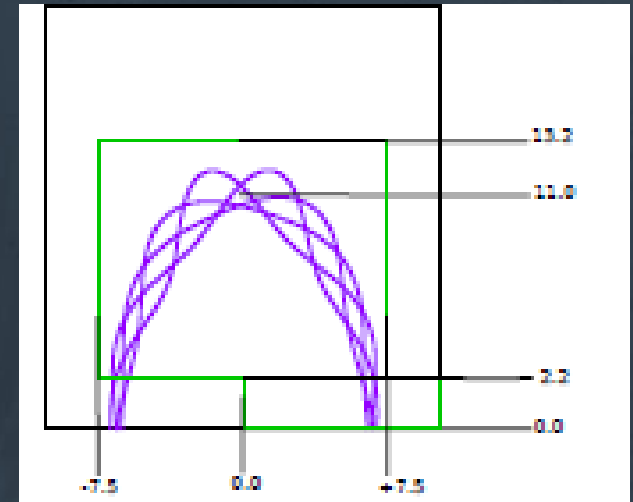
*Pinto et al 2016*



# Thermal emission: velocities

## Observations:

- *Doschek et al. 2008 ApJ*
  - *Doschek et al. 2013 ApJ*
  - *Young et al. 2013 ApJ*
- 
- **Correlation of non-thermal velocity dispersion (line broadening) with temperature: velocity dispersion increases from 20-30 km/s at ~1MK to 100-120km/s at ~15-20MK**
- 
- **Correlation of bulk velocity (line shift) with temperature: increases from 20-30 km/s to 300-350 km/s in the same interval**



**Model –  
calculate line-of-sight averages of  $\langle v \rangle$  and velocity dispersion in different parts of loop**

# Predicted LOS velocity dispersion

Large scale loop  
Strong field

Length 80 Mm

Time

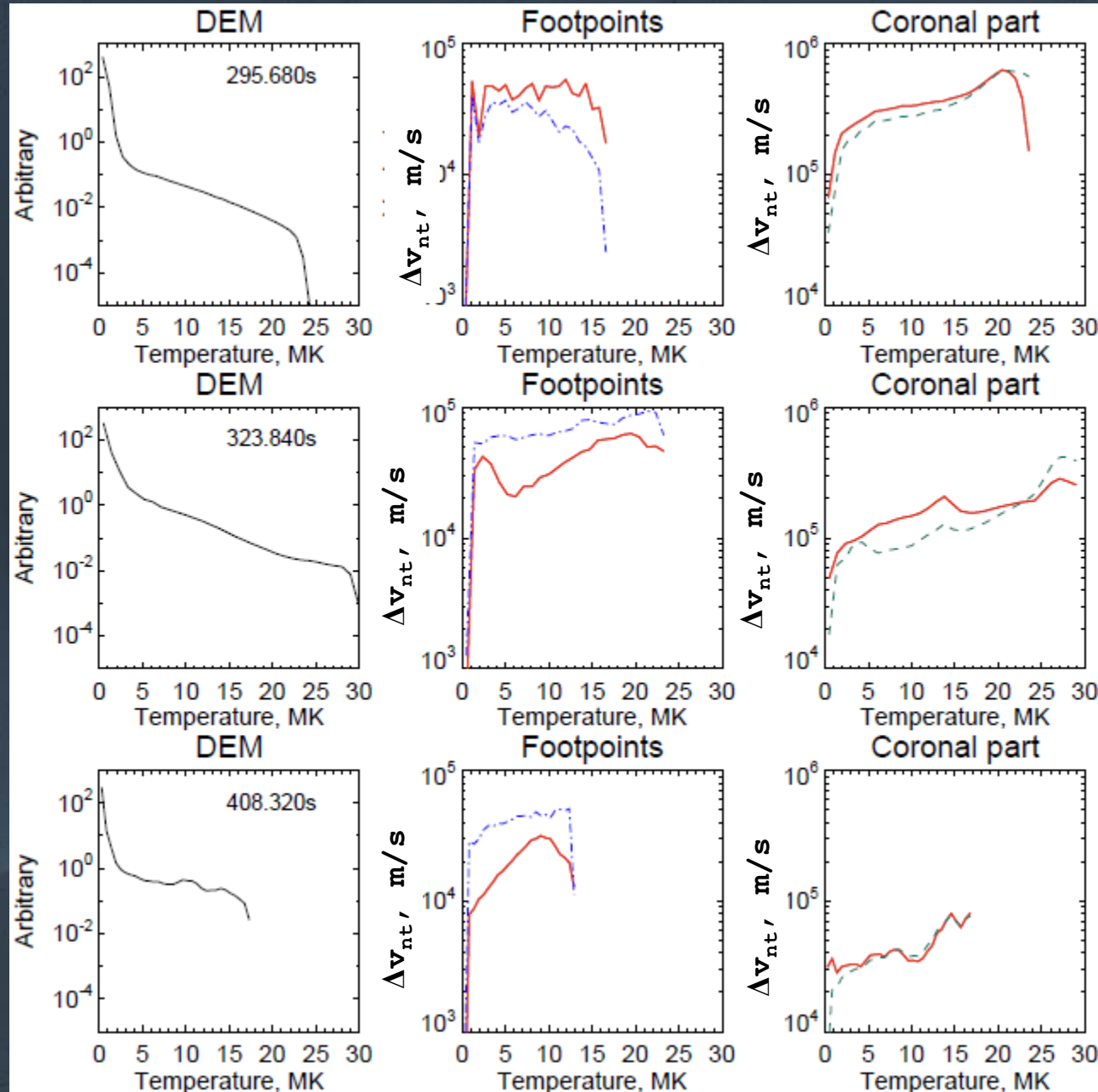
Footpoint field  
1500 G



Increasing  
velocity  
dispersion with  
temperature

Probably typical  
of all reconnection  
and not distinctive  
characteristic of  
twisted loops

*Gordovsky et al 2016*

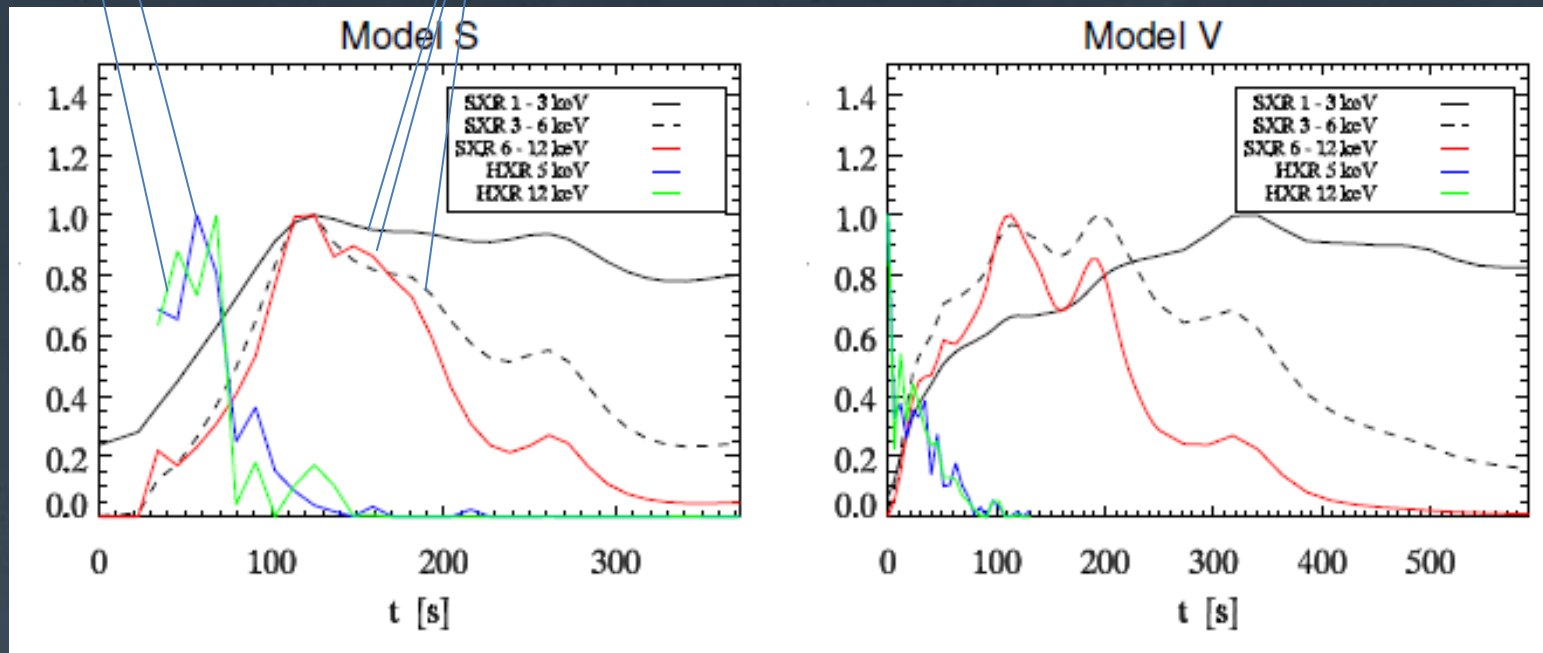


# Thermal (SXR) v. non-thermal (HXR) emission: Light curves

Neupert effect in small loops

HXR

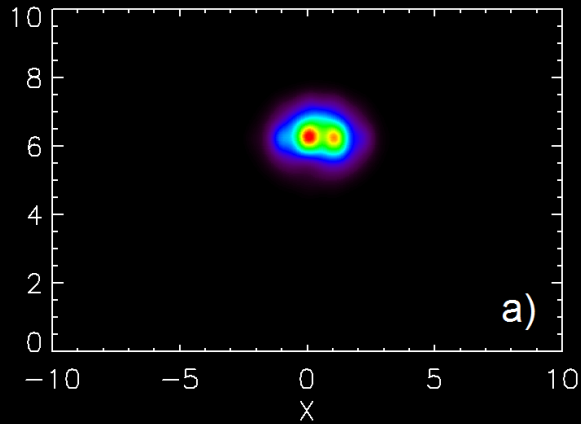
SXR



*Pinto et al 2016*

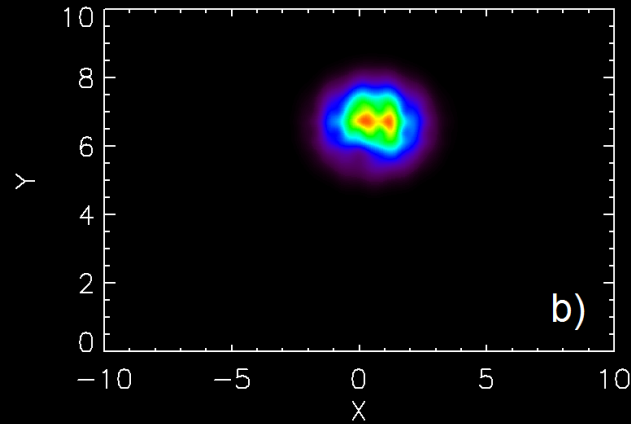
# Non-thermal HXR emission

Onset of  
magnetic  
reconnection



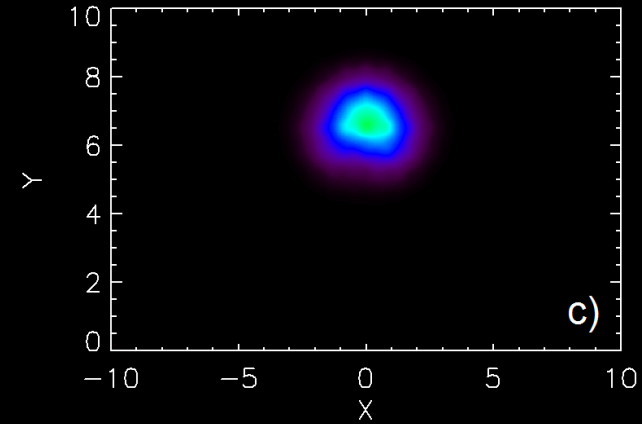
FWHM=1.9

Maximum  $dE/dt$   
in MHD model



FWHM=2.2

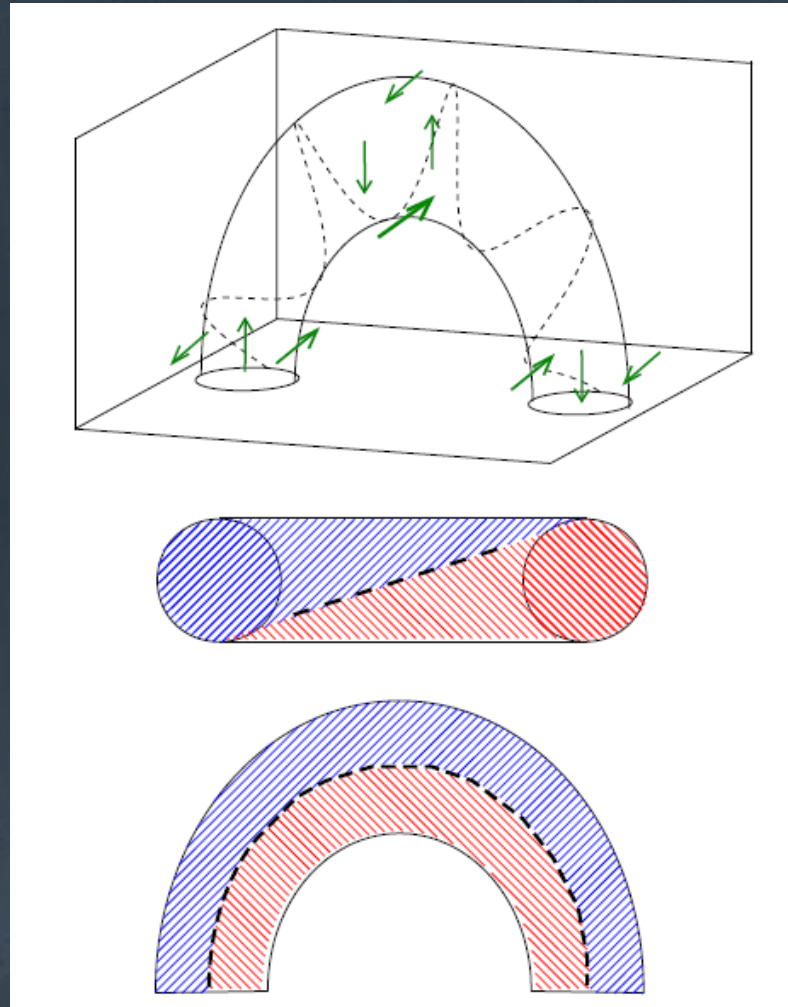
Decay phase



FWHM=2.5

- Loop cross-section increases by 50-100%
- FWHM of HXR sources increases by 20-30% (with RHESSI resolution)
- *Kontar et al 2011 ApJ* – quantitatively and qualitatively similar effect

# Circular microwave polarisation pattern as a detection tool?



Polarisation  
pattern –  
Disc

Limb

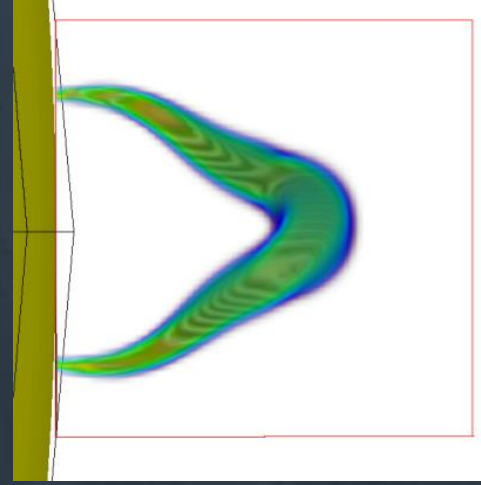
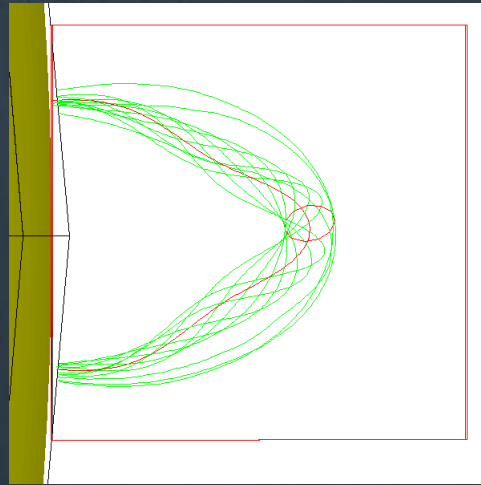
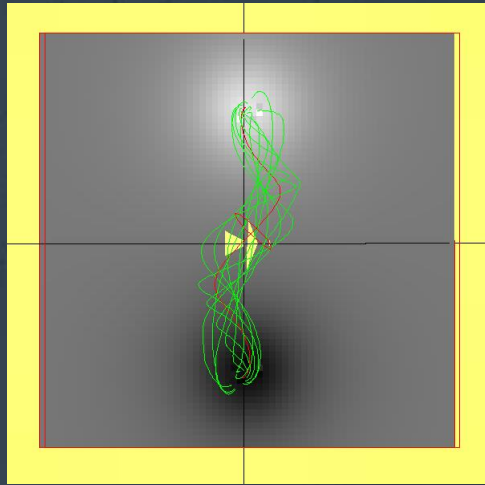
- *Kuznetsov et al. 2016, Sharykin & Kuznetsov 2016*
- *Gordovskyy et al, A&A under review*

# Synthetic microwave emission from twisted loops

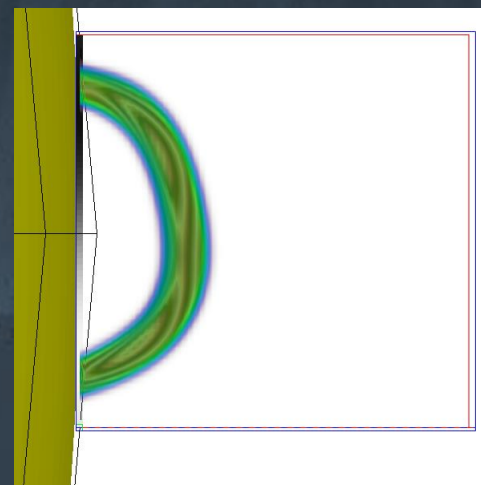
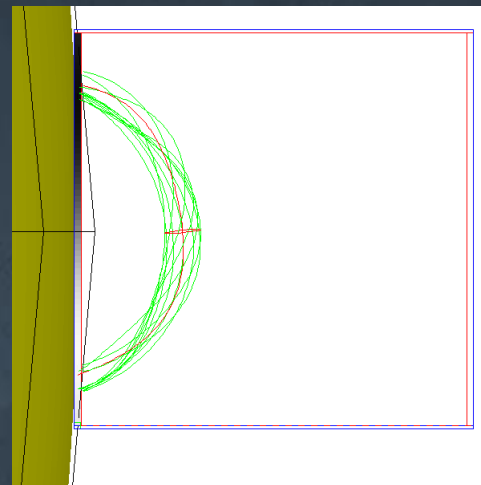
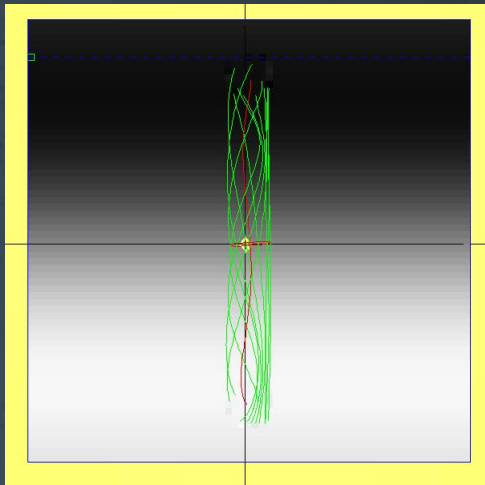
- Is this pattern visible in “real” loop (time duration, optical thick effects, more complex field.....)?
- GX Simulator  
*Fleishman & Kuznetsov (2010), Nita et al. (2015)*
- Magnetic field and thermal plasma density & temperature taken directly from MHD model
  - magnetic field 100-1000 G at foot-points
  - ambient plasma  $n \approx 10^9 \text{cm}^{-3}$ ,  $T \approx 1 \text{MK}$
  - hot plasma in the reconnecting loop  $n \approx 10^9 \text{cm}^{-3}$ ,  $T \approx 10\text{-}30 \text{MK}$
- Non-thermal electron population approximated by a single power law fitting test-particle simulations
  - $n_b \approx 5 \cdot 10^7 \text{cm}^{-3}$ ,  $E_{\text{low}} = 3 \text{keV}$ ,  $E_{\text{up}} = 3 \text{MeV}$ ,  $\gamma = 2.0\text{-}4.0$



# Magnetic field convergence effect



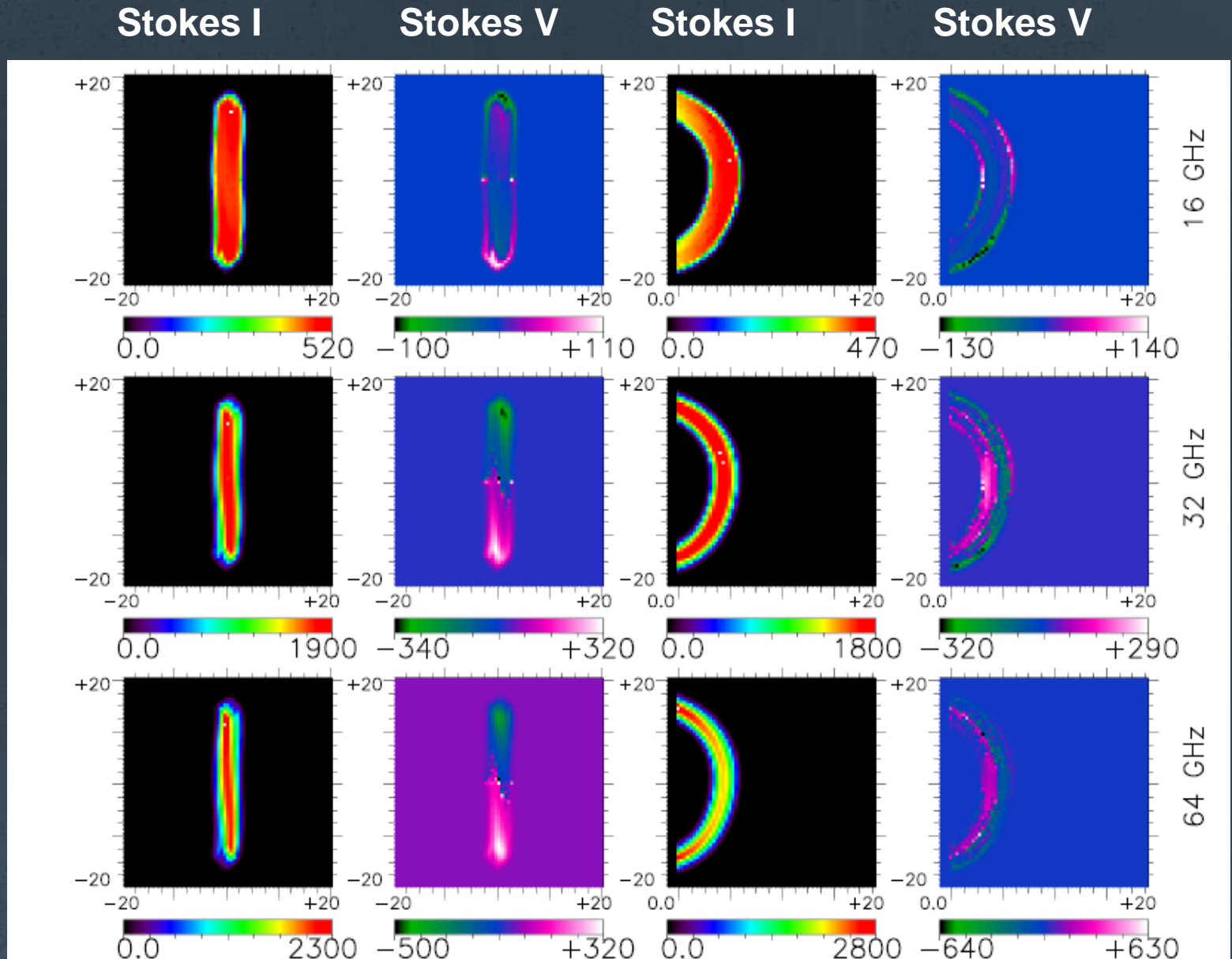
$$B_{\text{top}}/B_{\text{fp}} = 10$$



$$B_{\text{top}}/B_{\text{fp}} = 2$$

# MW polarisation in weakly converging loop

At start  
of  
energy  
release



# MW polarisation: frequency variation

4GHz

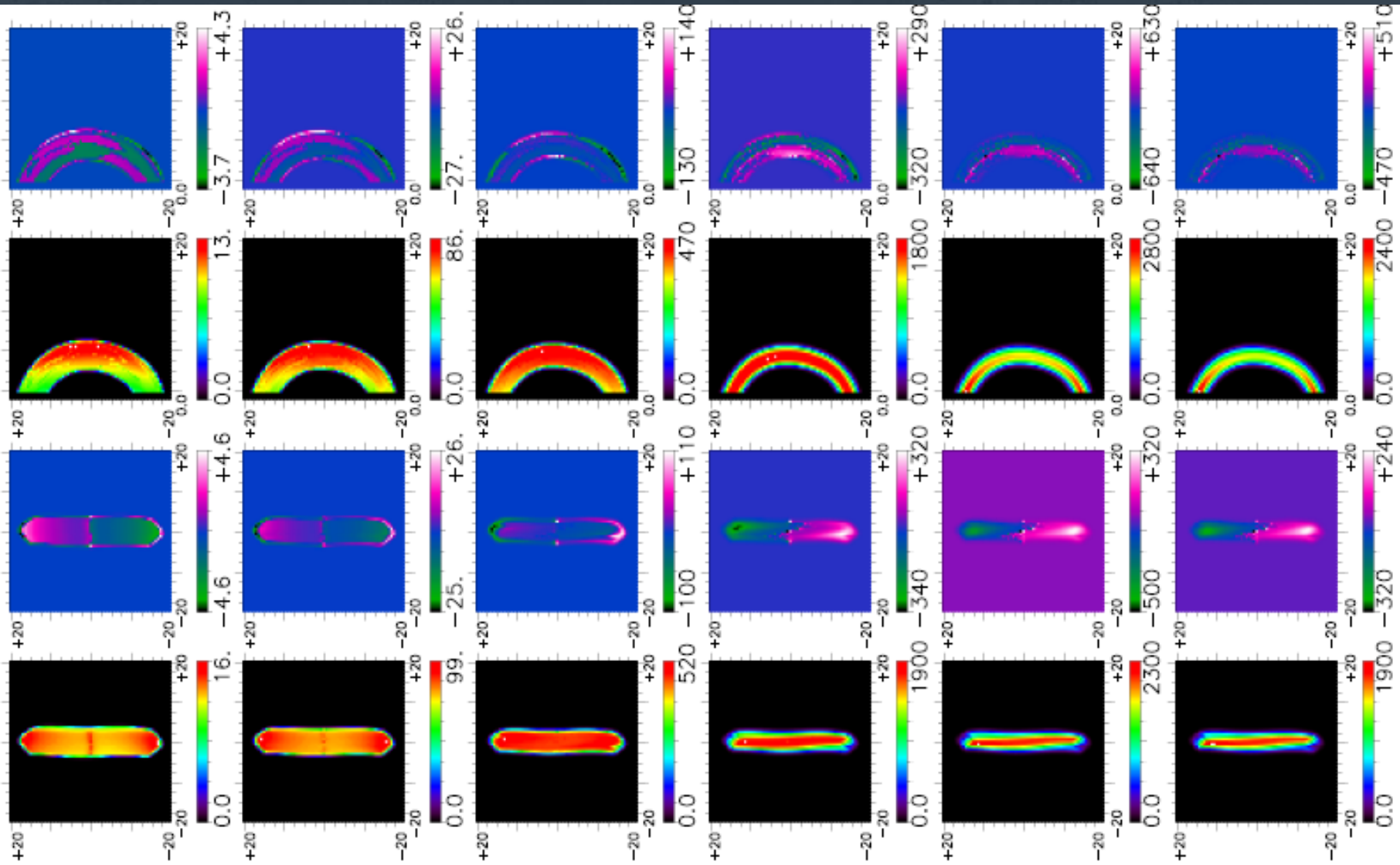
8GHz

16GHz

32GHz

64GHz

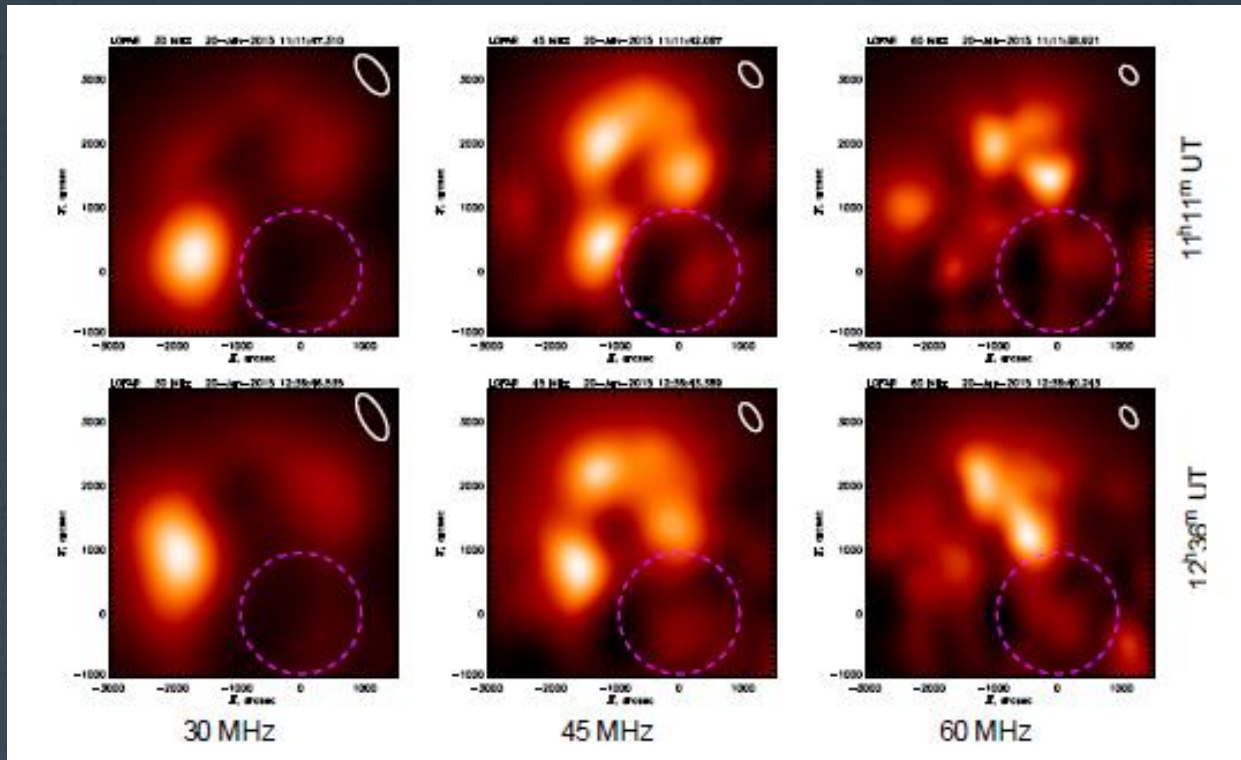
95GHz



- **Thermal microwave: circular polarisation gradient across the loop should be visible, especially at the limb, although the intensity should be low**
  
- **Non-thermal microwave: circular polarisation gradient across the loop should be visible, however**
  - **- the life-time of that pattern would be ~30-60s**
  - **- visibility of the pattern would depend on loop orientation**
  - **- visibility would depend on the magnetic field convergence**

# LoFAR observations of giant loop

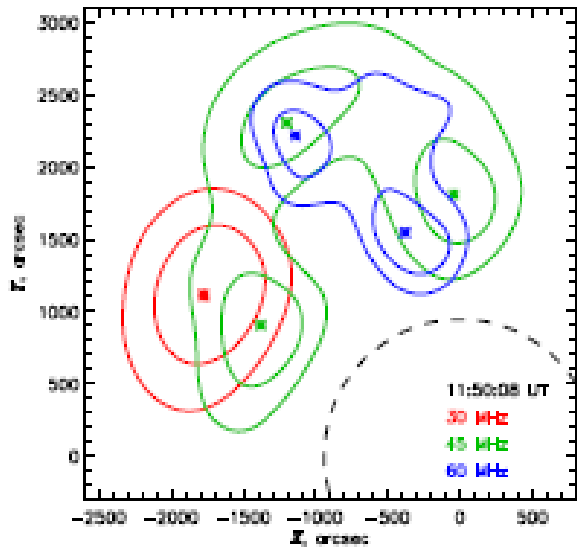
*Gordovsky, Kontar, Kuznetsov and Browning Ap J Lett submitted*



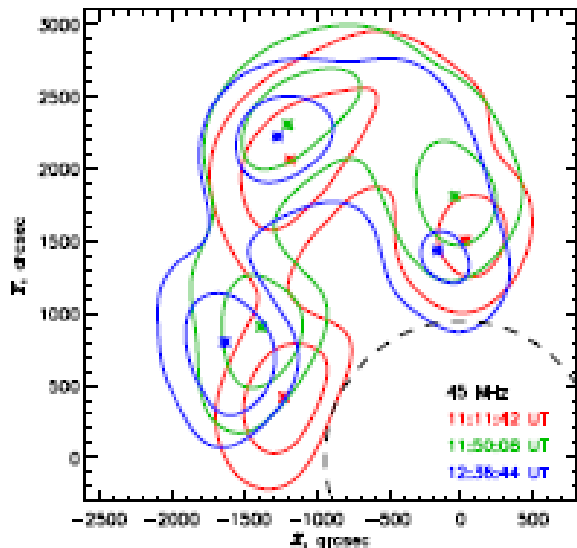
- Imaging of large-scale  $\sim 1 R_{\odot}$  loop in radio 30 – 60 MHz
- Narrow frequency range, high brightness temperature
- Plasma emission associated with energetic electrons (frequency  $\sim \sqrt{\text{density}}$ )

$30 \text{ MHz} \leftrightarrow 1.1 \times 10^7 \text{ cm}^{-3}$ ,  $60 \text{ MHz} \leftrightarrow 4.4 \times 10^7 \text{ cm}^{-3}$ ,

## Frequency variation



(a)



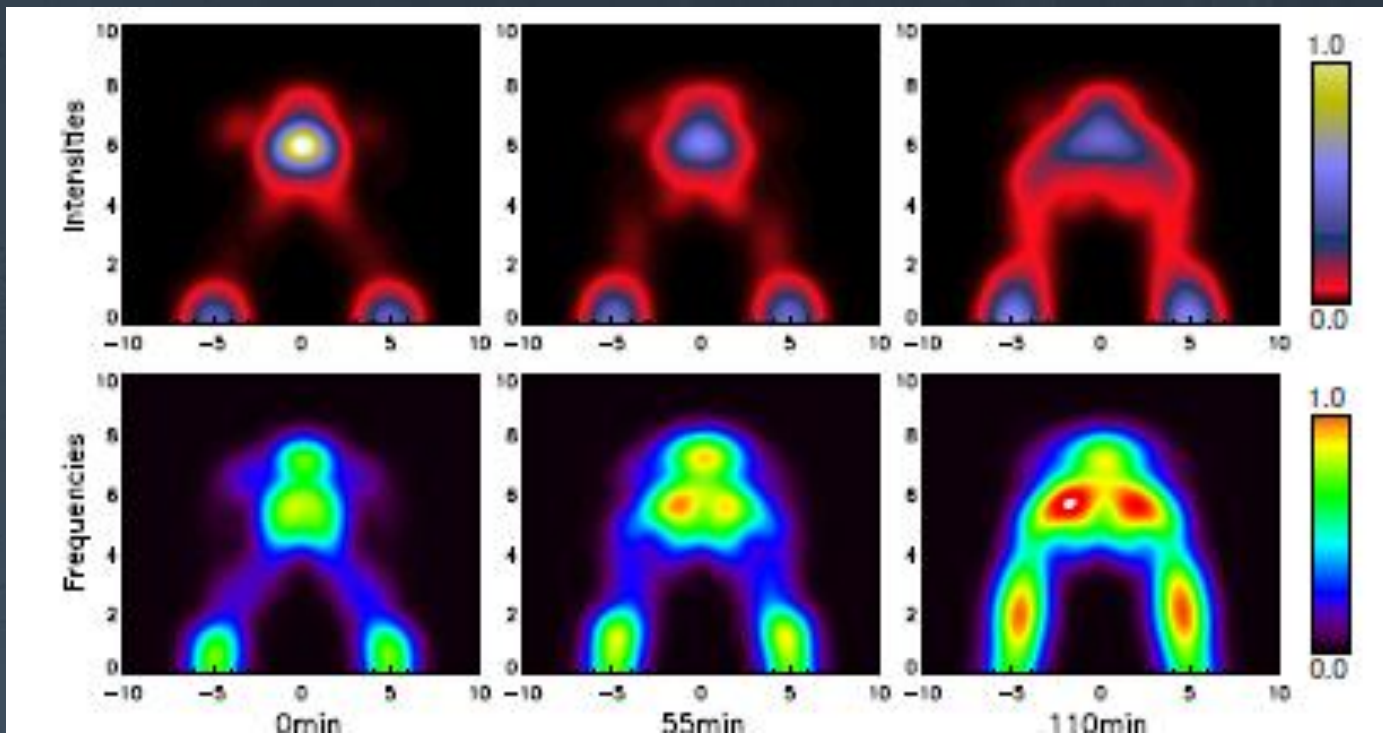
(c)

- Distinct loop structure
- Three main sources with weak background throughout loop
- Density does not match ambient coronal stratification – but strong gradients across loop
- Some asymmetry – density increases by factor of  $\sim 4$  from left to right
- No substantial expansion/motion in time

## Time variation

# Model

- 3D MHD simulation - large-scale twisted loop with strong convergence
- Use local Ohmic dissipation rate  $j \cdot E$  as proxy for particle acceleration
- Line-integrate - average frequency is local frequency (density) weighted by intensity
- Good agreement with many aspects of observations



Intensity

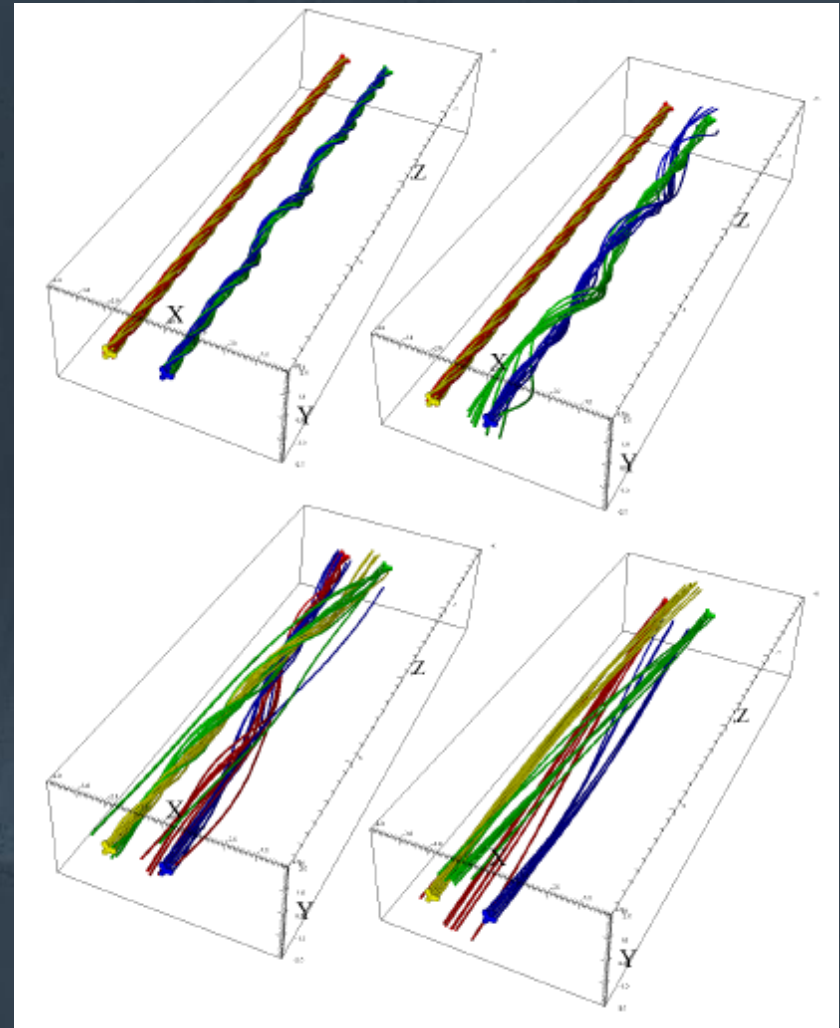
Average  
frequency

# Disruption of neighbouring stable flux rope

*Tam, Hood, Browning and Cargill A&A 2015*

Consider adjacent zero-net-current loops

- If the loops are sufficiently close, an unstable loop may trigger relaxation in a neighbouring stable loop
- In this case the two loops merge into a single (very weakly twisted) loop
- Releases energy stored in stable loop (as well as unstable loop)





Stable  
loop

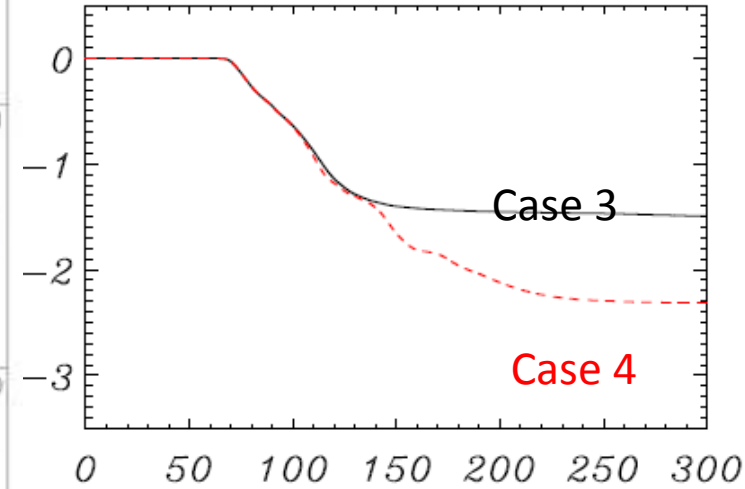
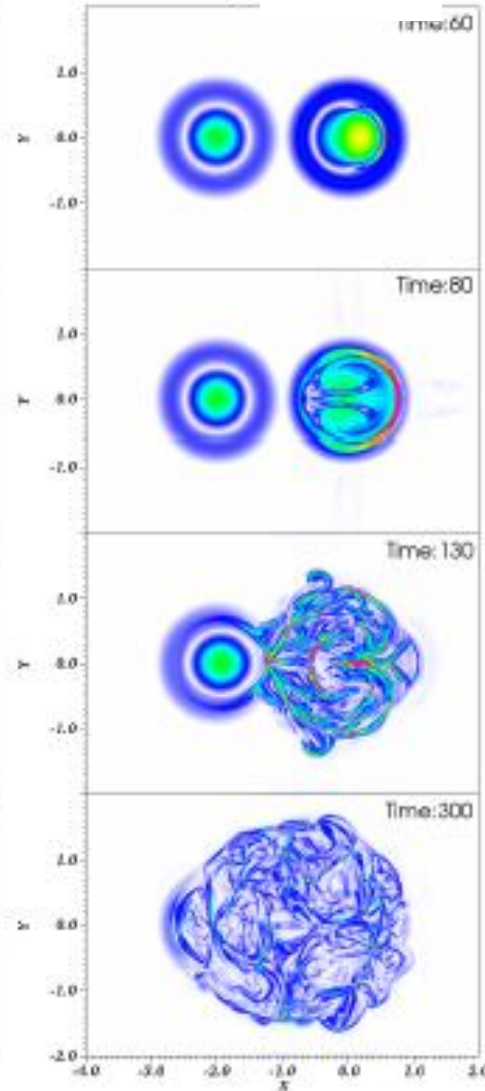
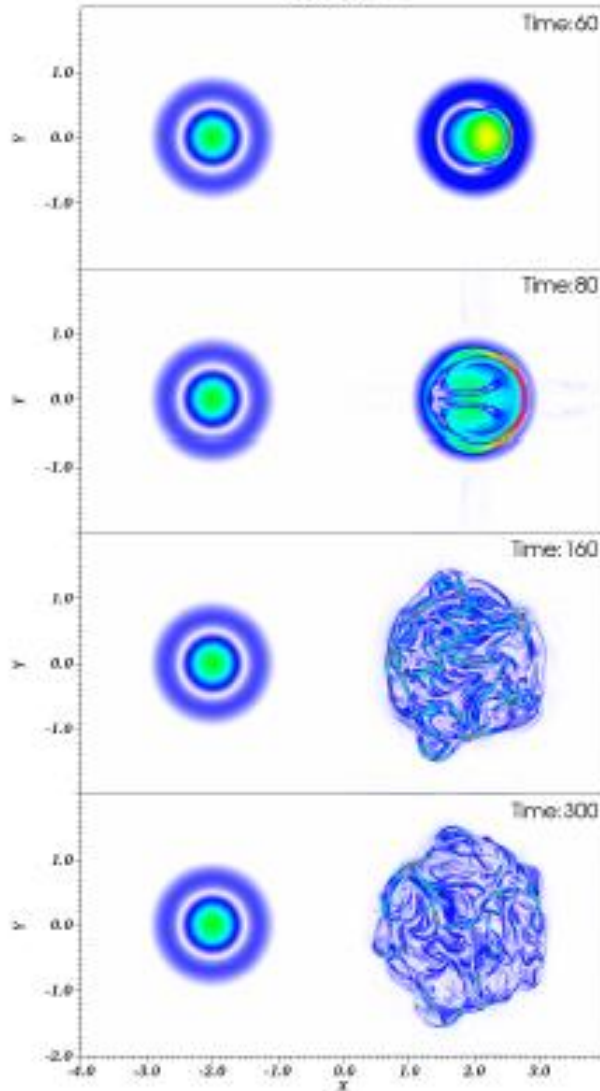
Unstable  
loop

Stable  
loop

Unstable  
loop

Case 3

Case 4



Magnetic  
energy versus  
time

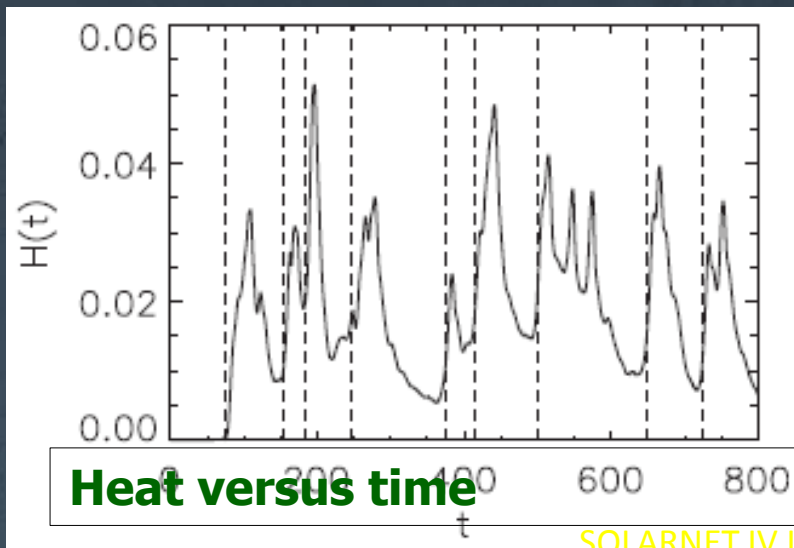
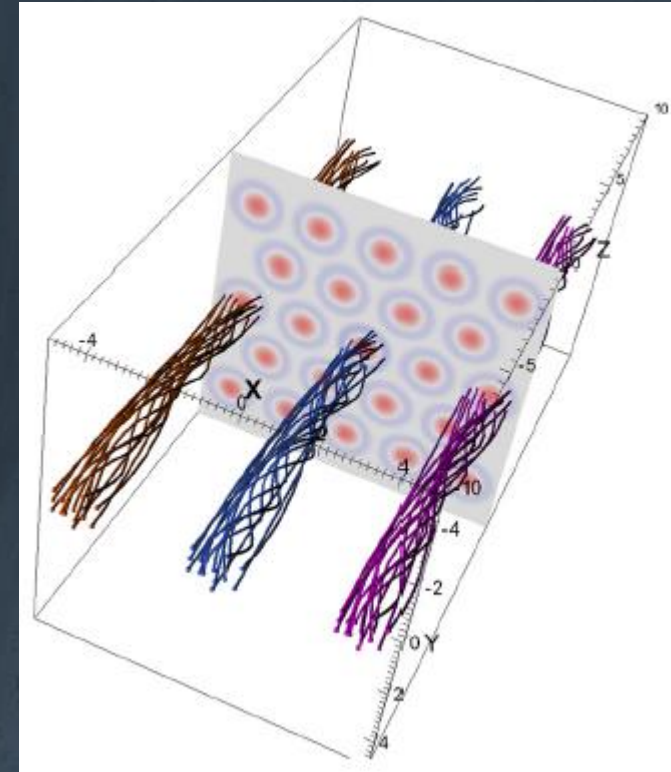
# Avalanche triggered by one unstable flux rope

→ Under certain conditions can have an avalanche of heating events

*Cf Lu and Hamilton 1981, Charbonneau et al 2001*

- First demonstration of avalanche - as in “cellular automaton” models - using 3D magnetohydrodynamics
- One unstable loop, 22 stable loops

*Hood, B et al Ap J 16*

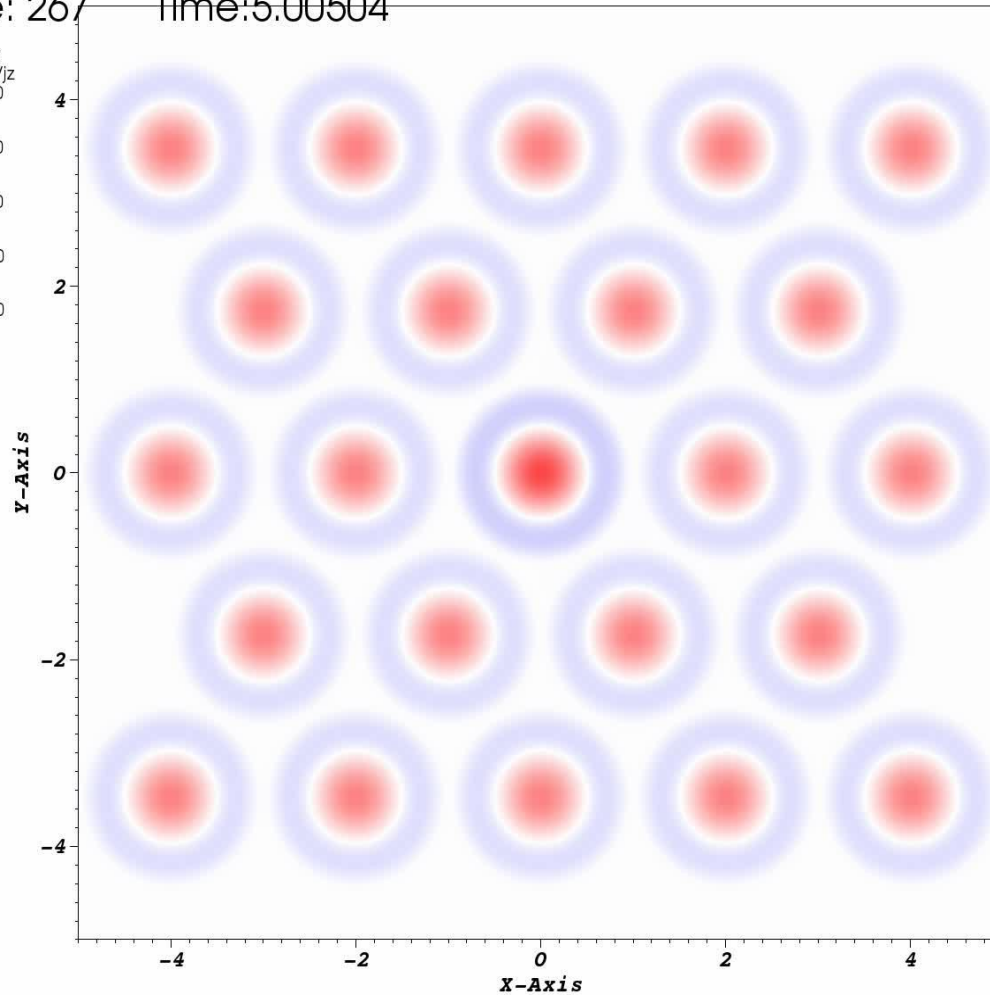


# Multiple loops - avalanche

DB: 0001.cfd

Cycle: 267 Time:5.00504

Pseudocolor  
Var: current/jz  
5.000  
2.500  
0.000  
-2.500  
-5.000  
Max: 3.599  
Min: -0.9177



Current in mid plane

One unstable loop,  
22 stable loops

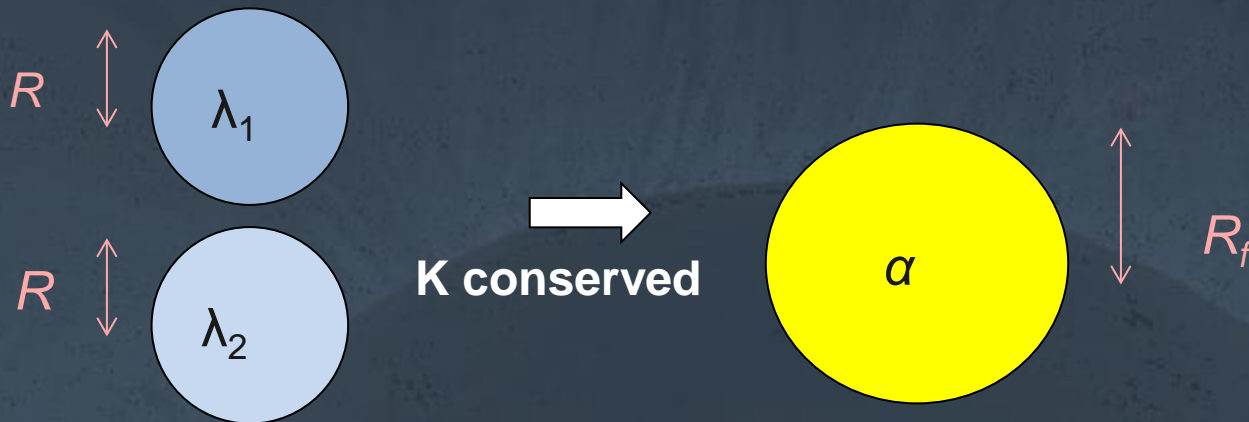
# Taylor relaxation model of avalanche

- Initial force-free flux ropes (following Tam et al 2015)
- Energy  $W$ , magnetic helicity  $K$  – superposition of flux ropes
- Relax conserving helicity - individually OR merge into single cylindrical flux rope with constant- $\alpha$

$$B_z = B_1 J_0(\alpha r), \quad B_\theta = B_1 J_1(\alpha r)$$

$$\sum_{\text{flux ropes}} K(\lambda_i, R) = K_{\text{constant-}\alpha}(\alpha, R_f) \rightarrow \text{value of } \alpha$$

- **Magnetic energy difference converted into heat**

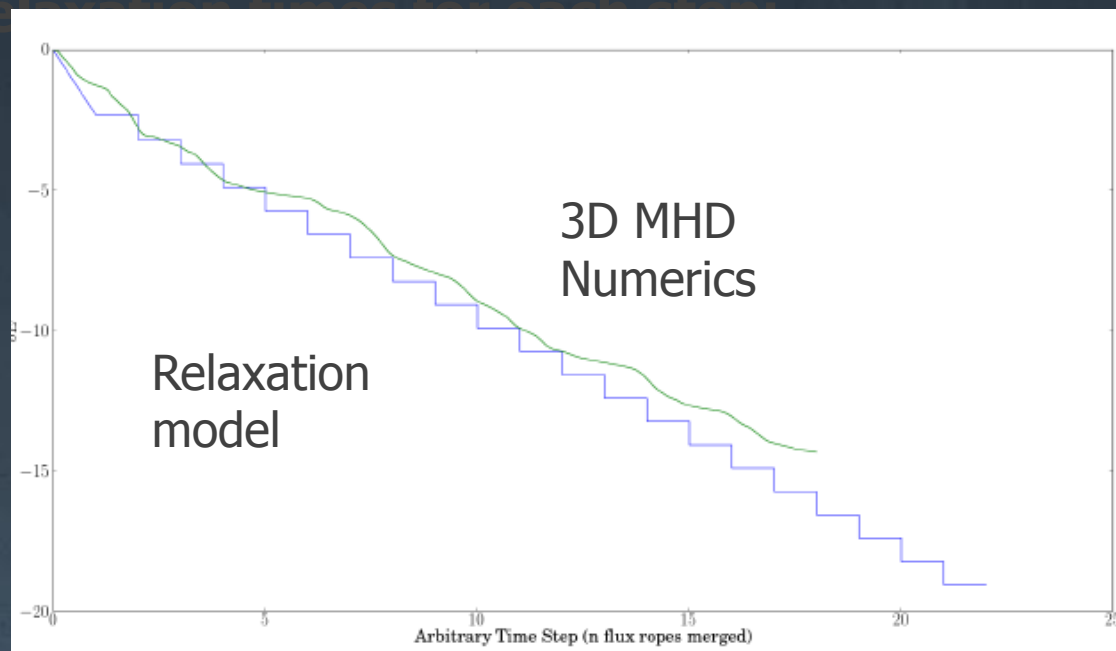


*Hussein et al, 2016  
A&A in press*

# Multi-rope relaxation model

- Very good agreement with MHD simulations
- Numerical simulations are very demanding of computational resource but relaxation model is quick and easy to calculate → use relaxation model to explore wide parameter space (different loop sizes, twists, etc)
- Apply relaxation model to 23 loop avalanche (assuming equal relaxation times for all loops)

Energy



Time

# Summary

- Twisted magnetic fields store free magnetic energy which may be released by reconnection in large-scale currents and fragmented current sheets triggered by ideal kink instability
- Modelled by coupled 3D MHD simulations and test-particles
- Predict spatial and temporal variations of thermal and non-thermal emission
- Thermal emission shows only weak twist
- Polarisation patterns of microwaves may be used to detect twisted fields in certain circumstances
- Large-scale loop structure observed by LoFAR, consistent with electron acceleration in large twisted loop
- Instability in one twisted thread may trigger energy release in many neighbouring stable threads
- First demonstration of heating avalanche using 3D MHD simulations, also relaxation model

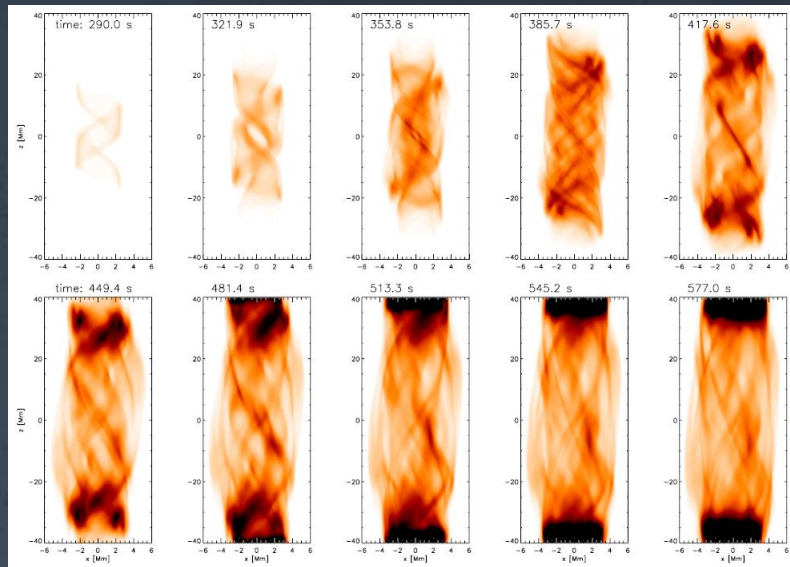
# Kink instability modelling and observables

- Field topology and thermal emission

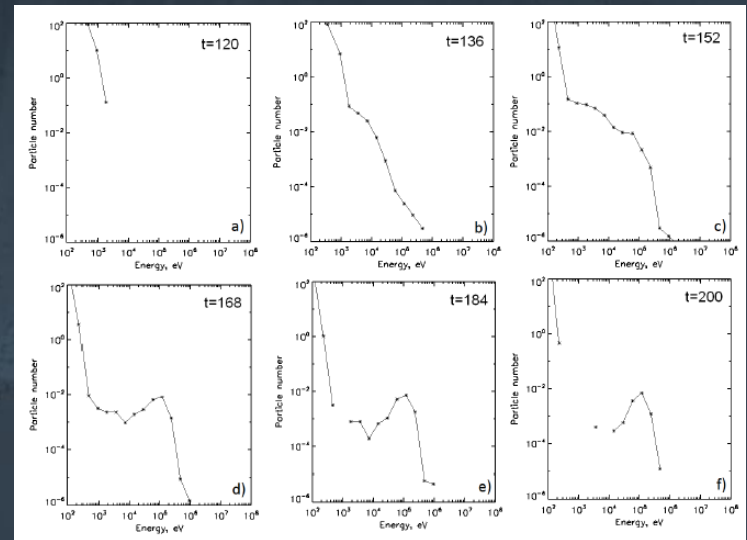
(e.g. Arber et al 1999; Botha et al 2012; Srivastava et al 2014; Pinto et al. 2015; Gordovskyy et al. 2016)

- Non-thermal particle spectra, HXR

(Gordovskyy & Browning 2012; Gordovskyy et al 2013, 2014; Pinto et al 2015)



from Botha et al. 2012 ApJ



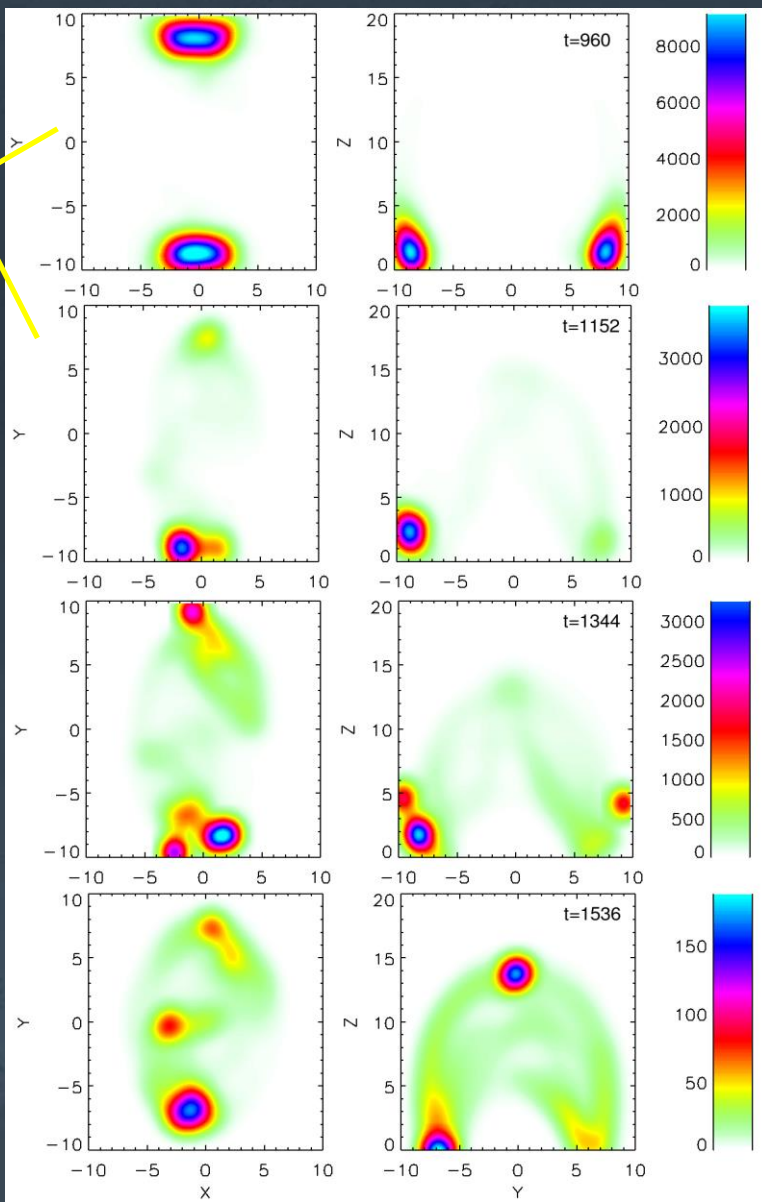
from Gordovskyy et al 2013 SolPhys

# Non-thermal HXR emission

Low-density case,  $\epsilon=10\text{keV}$

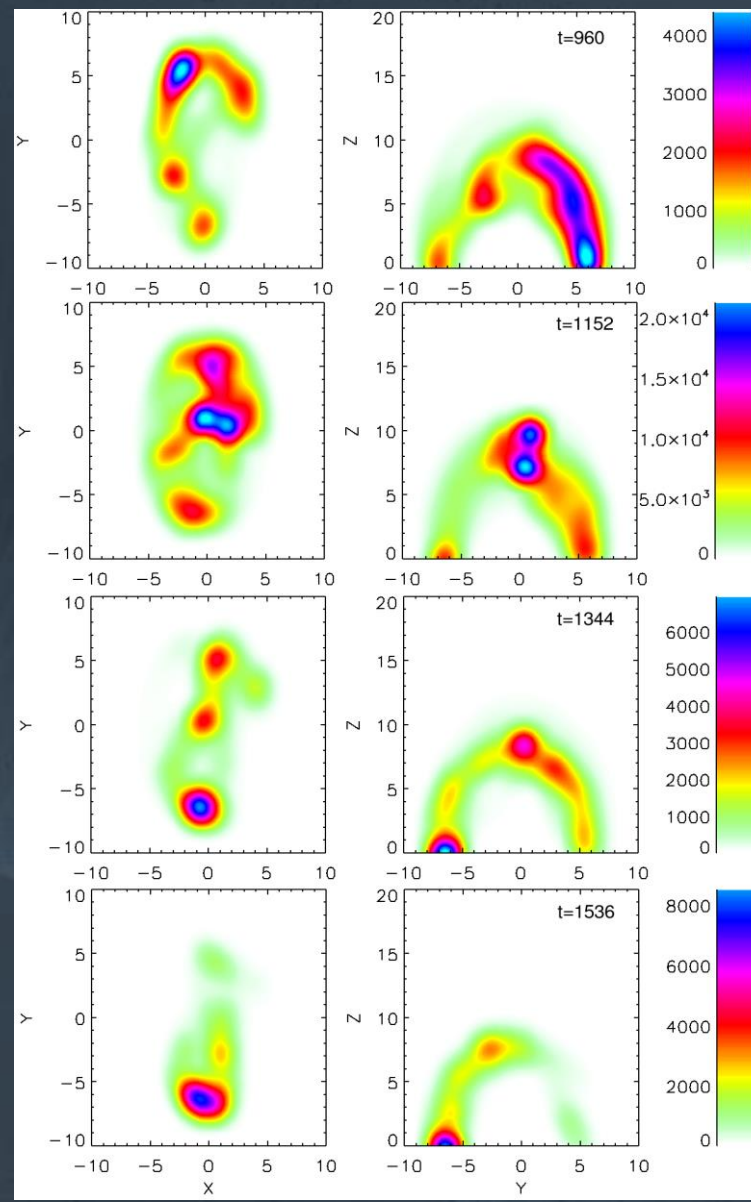
High-density case,  $\epsilon=10\text{keV}$

Onset of magnetic reconnection



Maximum  $dE/dt$  in MHD model

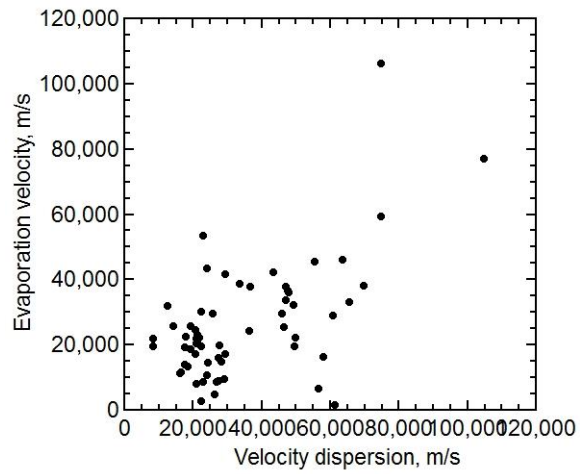
Decay phase



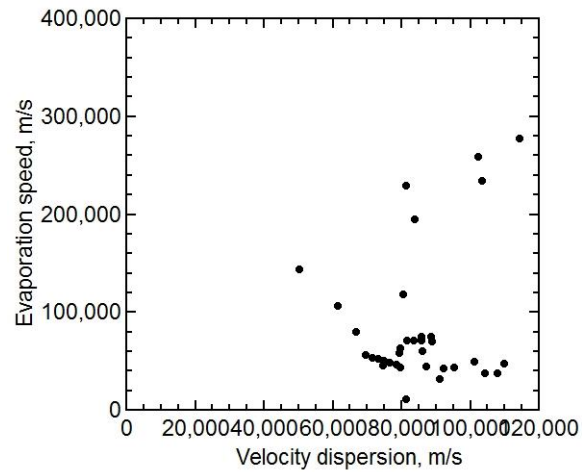


# LOS velocity dispersion v. bulk LOS velocity

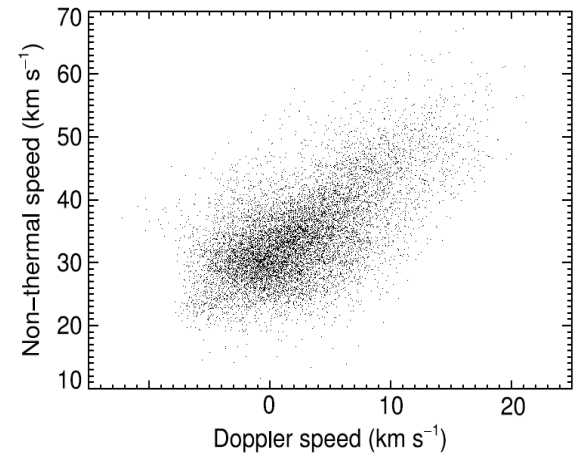
## Model V



## Model Z

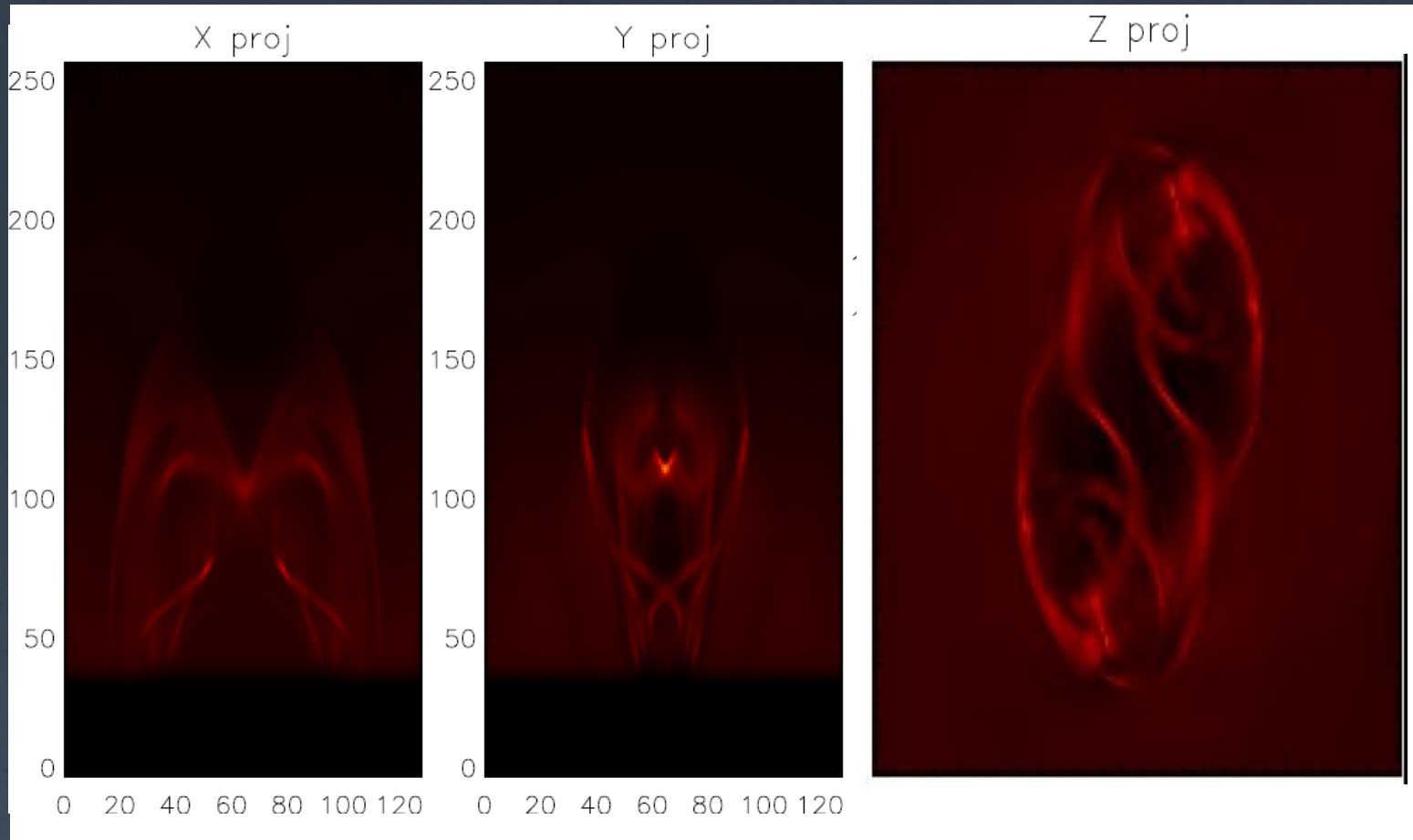


## Observed (after Doschek et al. 2008)



# SXR thermal emission (LOS integrated)

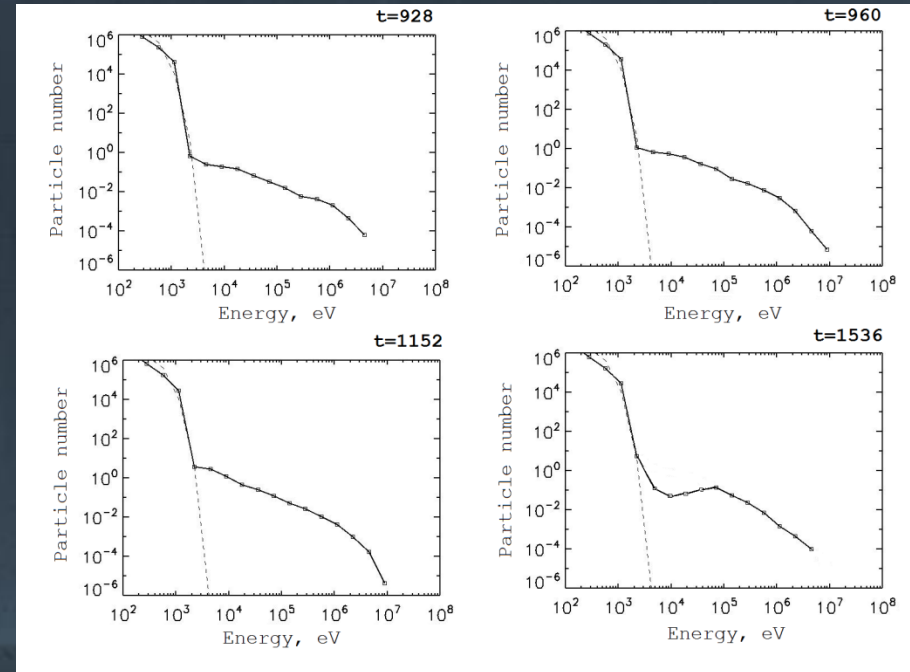
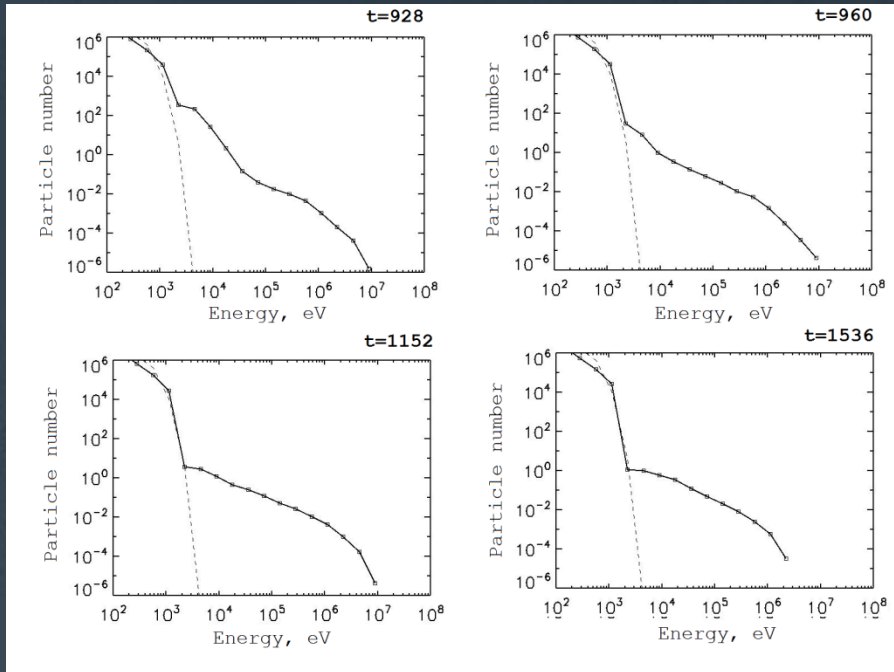
2keV emission (near onset, large loop)



# Particle kinetics: electron energy spectra

Low density loop

High density loop



Time evolution of electron energy spectrum

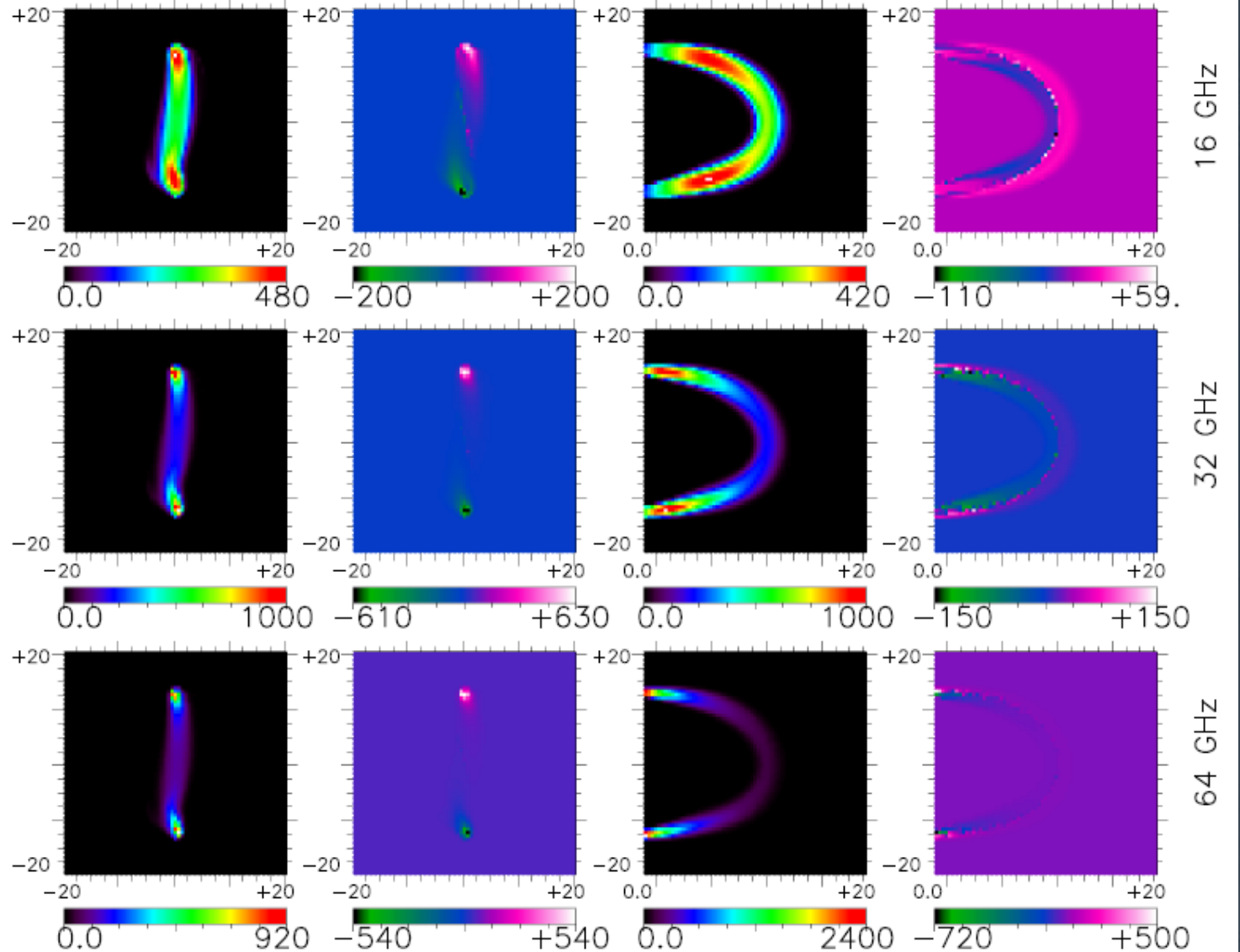
# MW polarisation in strongly converging loop

Stokes I

Stokes V

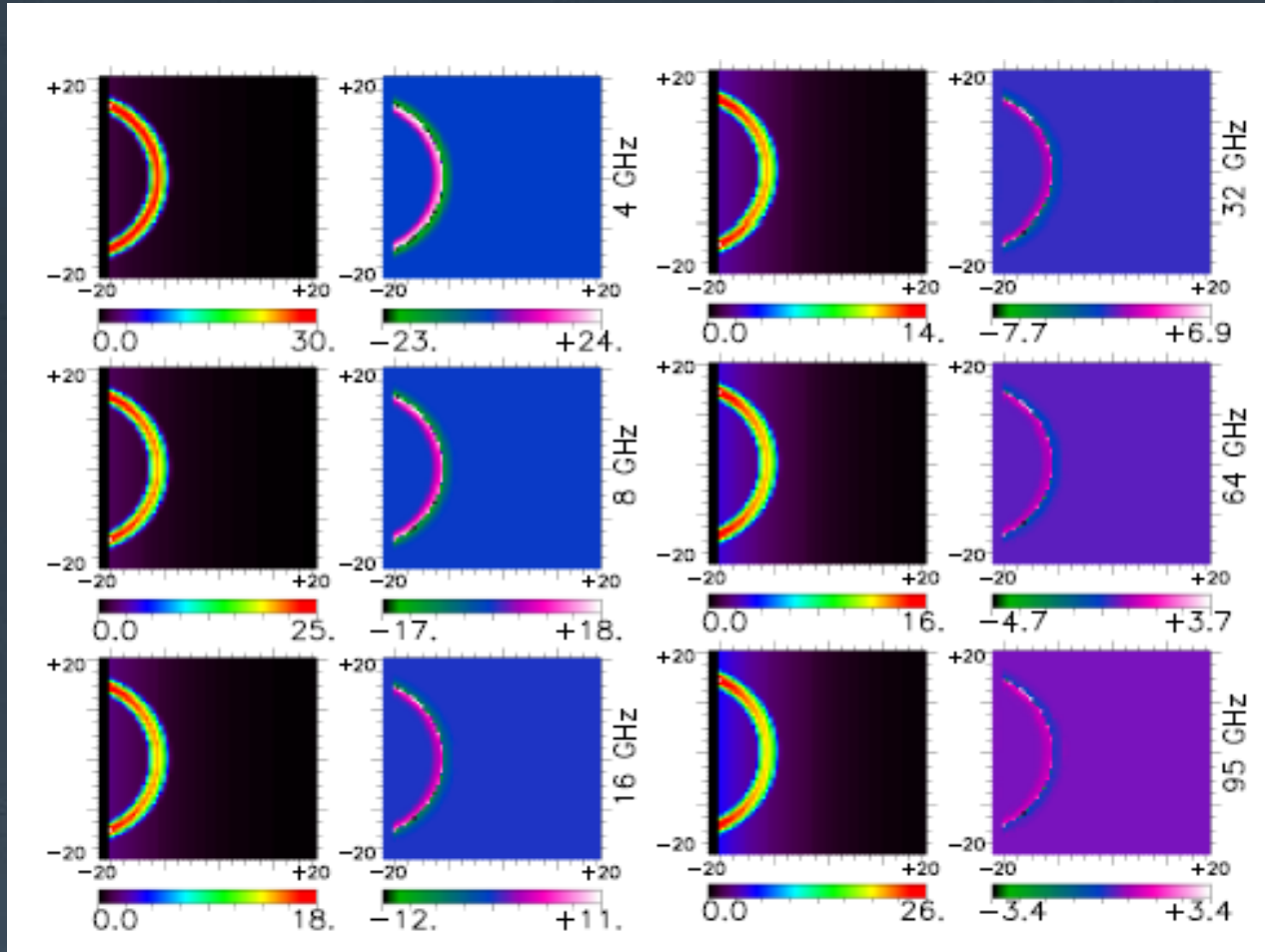
Stokes I

Stokes V



At start  
of  
energy  
release

# Purely thermal MW emission



Weakly converging loop, footpoint field 360 G, time of peak temperature (about +30s after onset of reconnection)

# Benchmarking relaxation model

- To determine final radius  $R_f$ , assume magnetic pressure continuous at flux rope surface (other constraints also considered)

$$B_z^2(R_f) + B_\theta^2(R_f) = B_z^2 \text{ initial} = B_0^2 \left(1 - \frac{\lambda^2}{7}\right)$$

	Final Radius	$\Delta E_{\text{MHD}}$	$\Delta E_{\text{Taylor}}$
Case 1*	0.999 (each)	-3.031	-2.608
Case 2	1.412	-3.069	-3.164
Case 3*	0.999	-1.5	-1.304
Case 4	1.413	-2.3	-2.29

\*Case 3 only considers one flux rope (in the Taylor model) relaxing since the first rope is stable, case 1 is twice of case 3.

- Good agreement between relaxation model and numerics