

UNDERSTANDING DYNAMO MECHANISMS FROM 3D CONVECTION SIMULATIONS OF THE SUN

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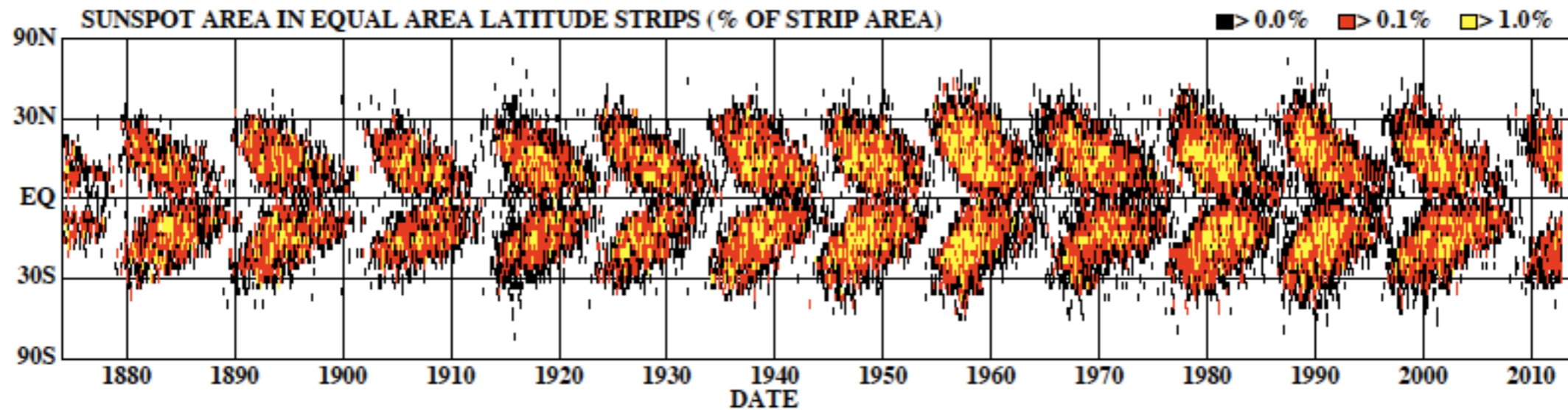
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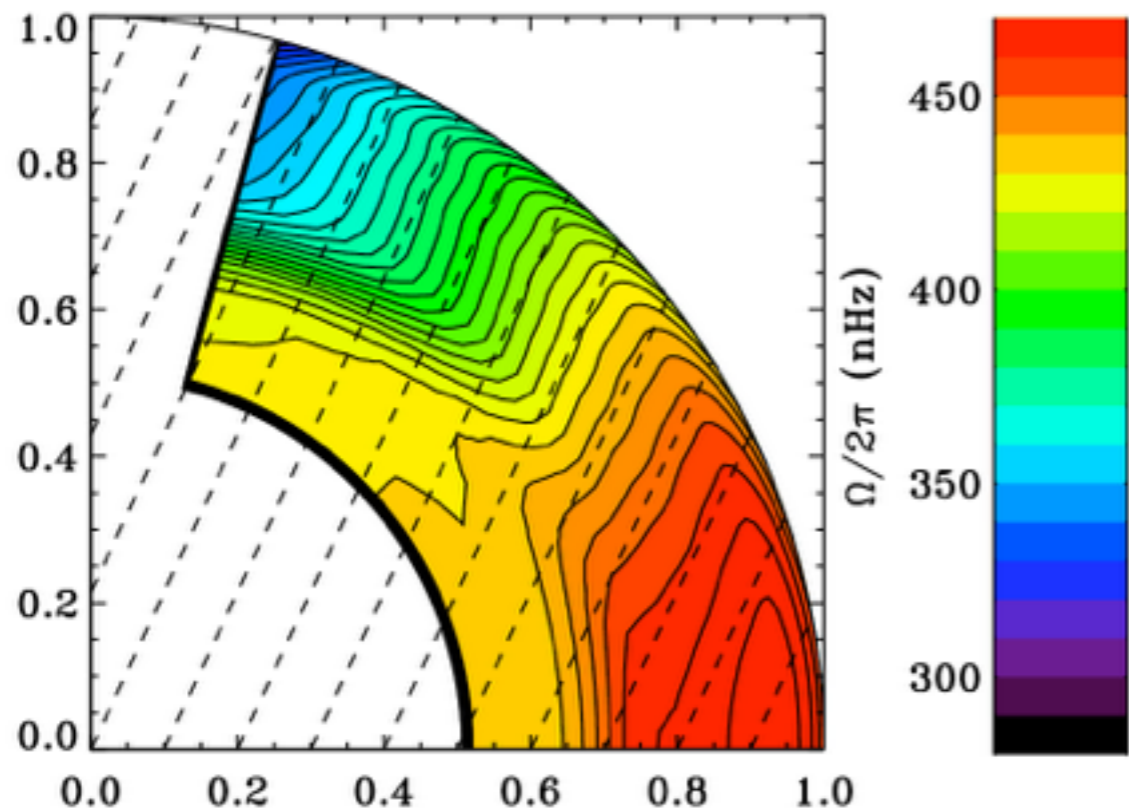
MATTHIAS RHEINHARDT, AALTO UNIVERSITY

Solar Activity

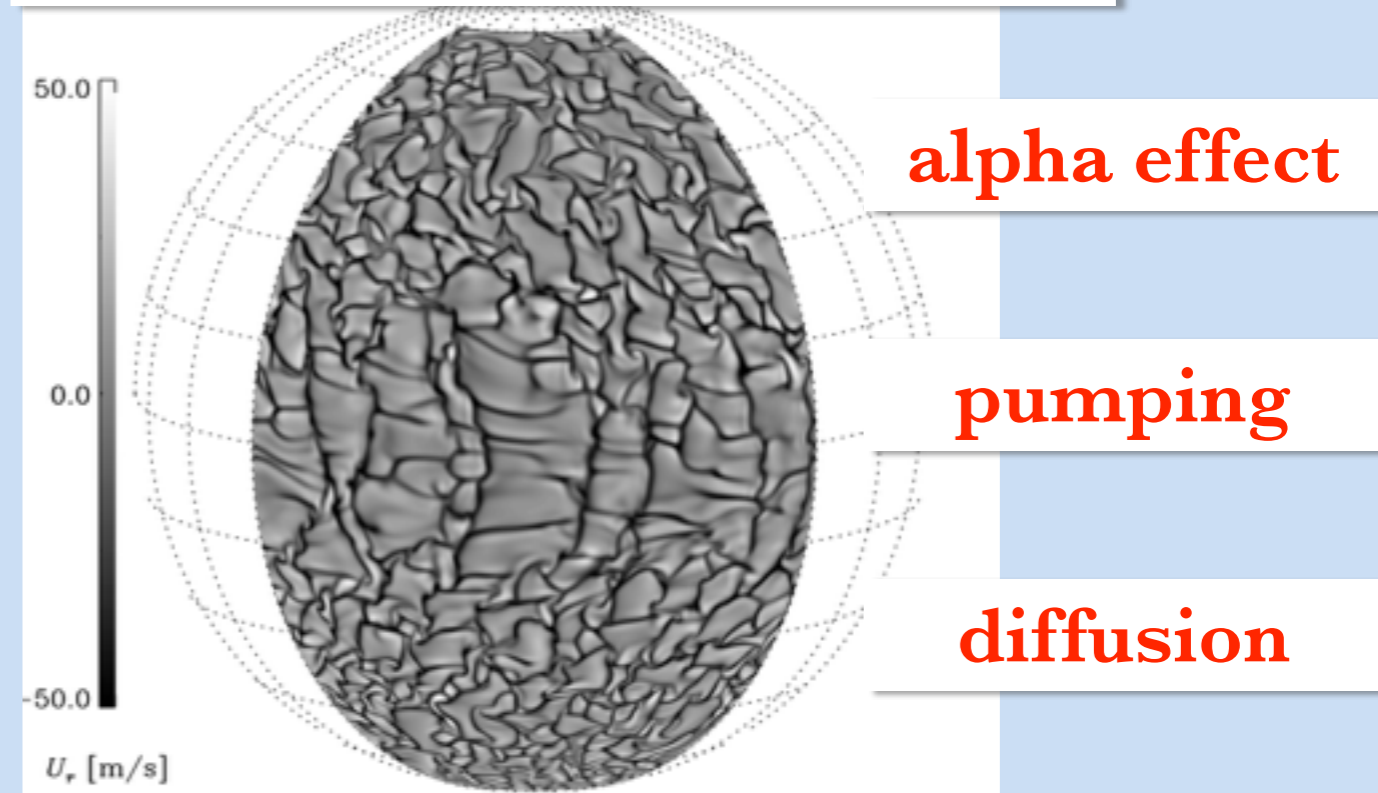
DAILY SUNSPOT AREA AVERAGED OVER INDIVIDUAL SOLAR ROTATIONS



differential rotation



turbulent convective motions



ME **no direct measurements**

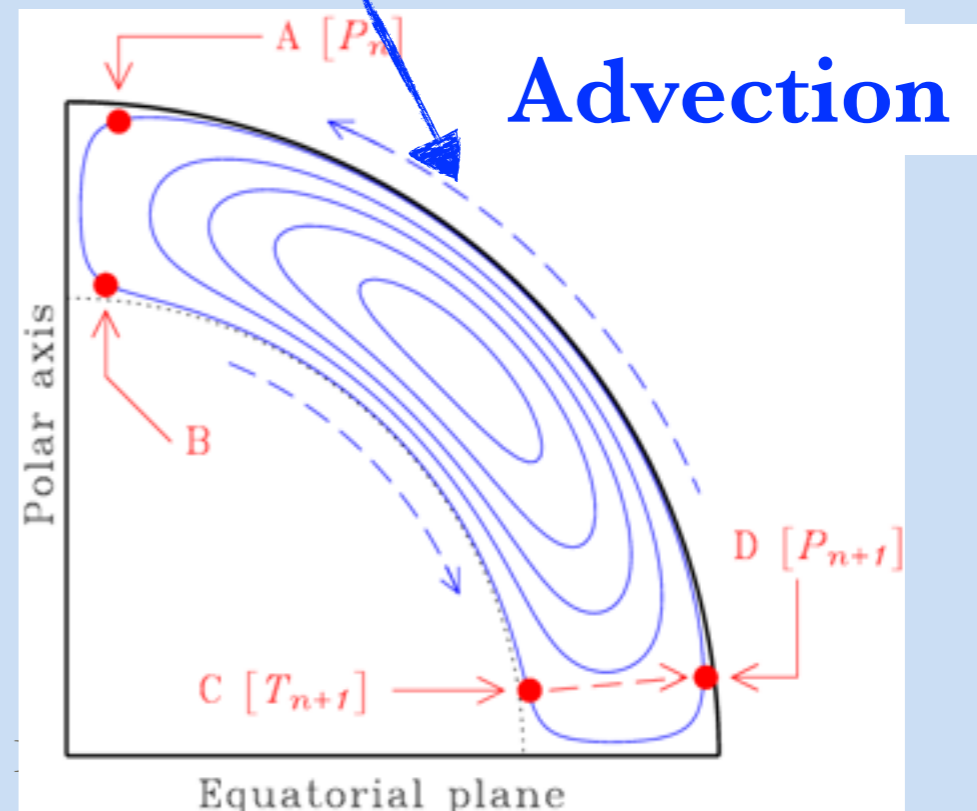
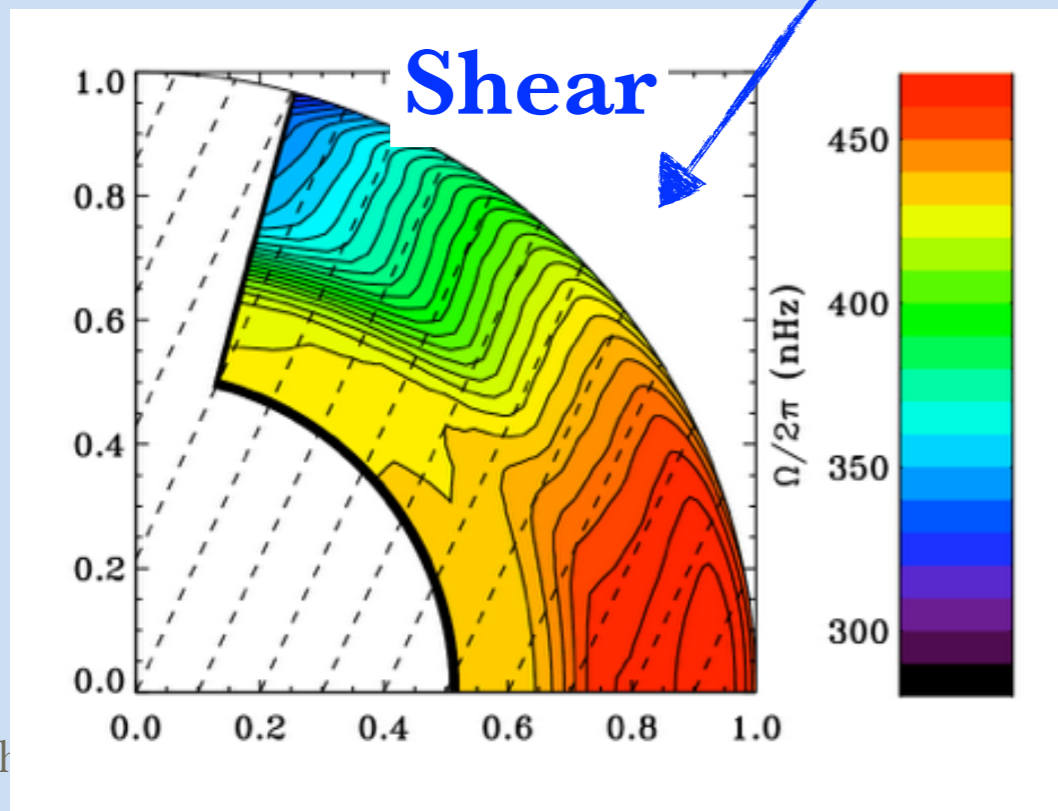
Dynamos

$$\frac{\partial B}{\partial t} = \nabla \times (u \times B) - \nabla \times \eta J$$

$$B = \bar{B} + b' \quad u = \bar{U} + u'$$

$$\frac{\partial \bar{B}}{\partial t} = \nabla \times (\bar{U} \times \bar{B}) + \overline{u' \times b'} - \nabla \times \eta \bar{J}$$

$$\nabla \times (\bar{U} \times \bar{B}) = (\bar{B} \cdot \nabla) \bar{U} - \bar{B} (\nabla \cdot \bar{U}) - (\bar{U} \cdot \nabla) \bar{B}$$



Electromotive force

$$\mathcal{E} = \mathbf{a} \cdot \bar{\mathbf{B}} + \mathbf{b} \cdot \nabla \bar{\mathbf{B}} + \dots$$

$$\mathcal{E}_i = a_{ij} \bar{B}_j + b_{ijk} \partial_j \bar{B}_k + \dots$$


$$\mathcal{E} = \alpha \cdot \bar{\mathbf{B}} + \gamma \times \bar{\mathbf{B}} - \beta \cdot (\nabla \times \bar{\mathbf{B}}) - \delta \times (\nabla \times \bar{\mathbf{B}}) - \kappa \cdot (\nabla \bar{\mathbf{B}})^{(S)}$$

Test-field method

Schrinner et al. 2005, 2007, 2012

$$\frac{\partial \bar{\mathbf{B}}}{\partial t} = \nabla \times (\bar{\mathbf{u}} \times \bar{\mathbf{B}} + \overline{\mathbf{u}' \times \mathbf{B}'}) - \nabla \times \eta \nabla \times \bar{\mathbf{B}},$$

$$\mathcal{E} = \alpha \cdot \bar{\mathbf{B}} + \gamma \times \bar{\mathbf{B}} - \beta \cdot (\nabla \times \bar{\mathbf{B}}) - \delta \times (\nabla \times \bar{\mathbf{B}}) - \kappa \cdot (\nabla \bar{\mathbf{B}})^{(S)}$$

$$\frac{\partial \mathbf{B}'}{\partial t} = \nabla \times (\mathbf{u}' \times \bar{\mathbf{B}}^T + \bar{\mathbf{u}} \times \mathbf{B}' + \mathbf{u}' \times \mathbf{B}' - \overline{\mathbf{u}' \times \mathbf{B}'}) - \nabla \times \eta \nabla \times \mathbf{B}'$$

i	1	2	3	4	5	6	7	8	9
$\bar{B}_{Tr}^{(i)}$	see talk by Fred Gent					0	ϑ	0	0
$\bar{B}_{T\vartheta}^{(i)}$	0	1	0	0	r	0	0	ϑ	0
$\bar{B}_{T\varphi}^{(i)}$	0	0	1	0	0	r	0	0	ϑ

The Simulation

Global convective dynamo simulations

$$\frac{\partial A}{\partial t} = u \times B + \eta \nabla^2 A$$

$$\frac{D \ln \rho}{Dt} = -\nabla \cdot u$$

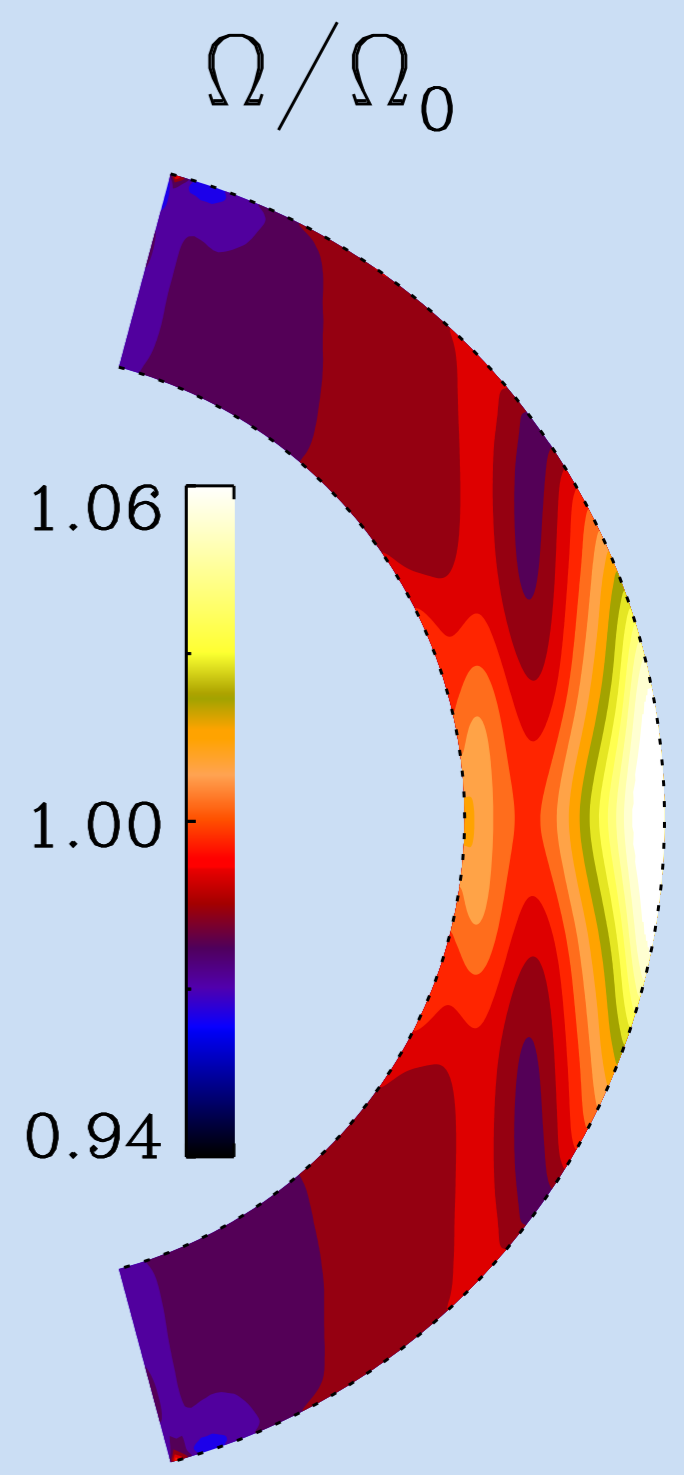
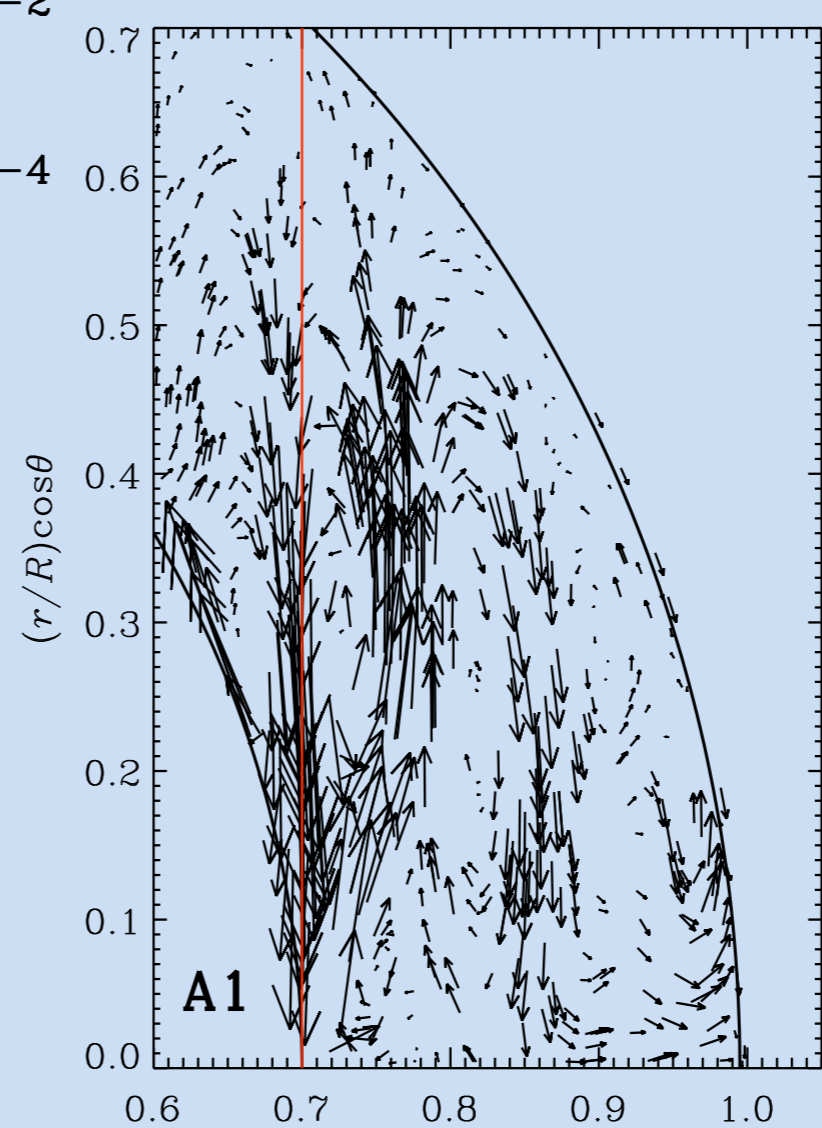
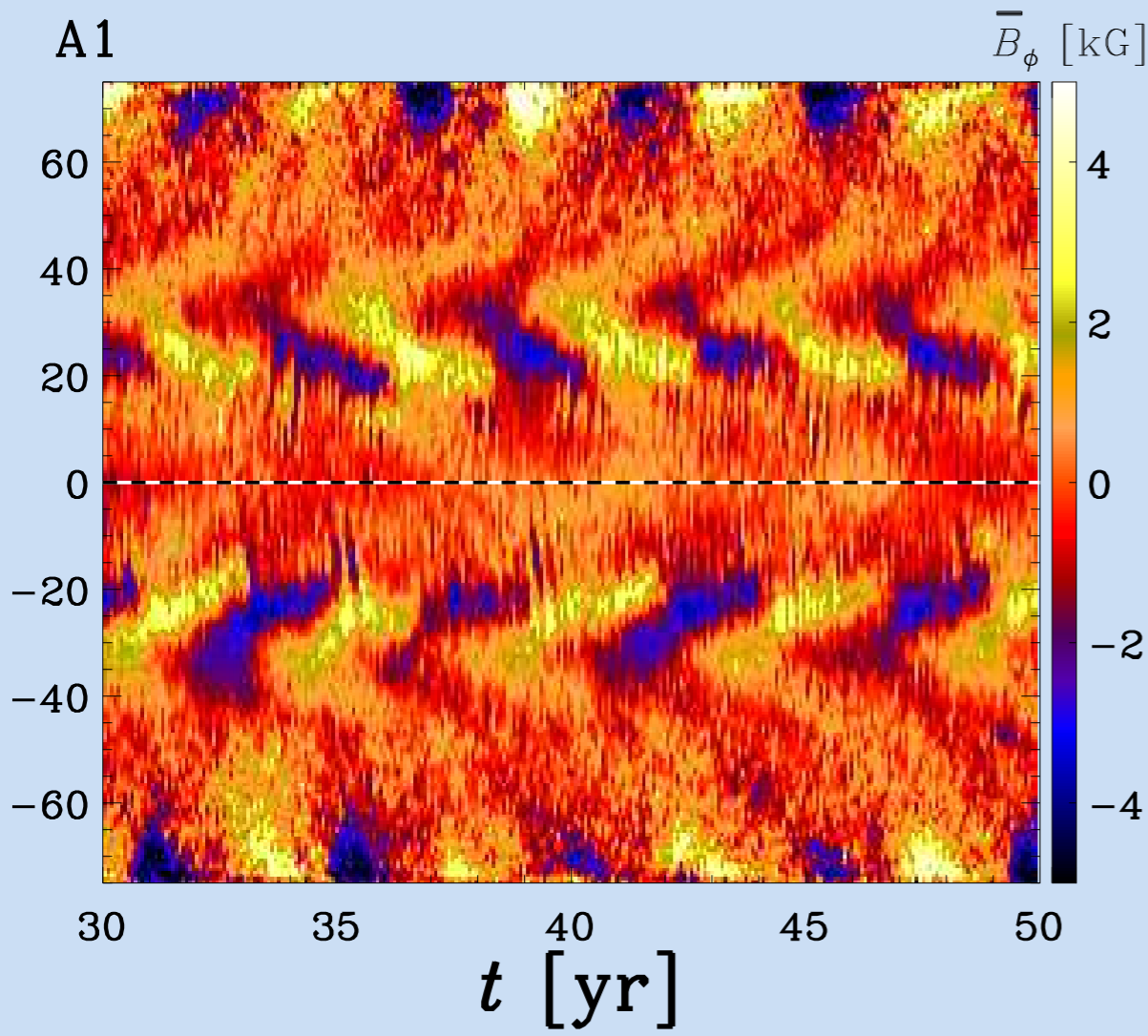
$$\frac{Du}{Dt} = g - 2\Omega_0 \times u + \frac{1}{\rho} (J \times B - \nabla p + \nabla \cdot 2\nu \rho S)$$

$$T \frac{Ds}{Dt} = \frac{1}{\rho} \nabla \cdot (K \nabla T + \chi_t \rho T \nabla s) + 2\nu S^2 + \frac{\mu_0 \eta}{\rho} J^2 - \Gamma_{\text{cool}}(r),$$



- high-order finite-difference code
- scales up efficiently to over 60.000 cores
- compressible MHD

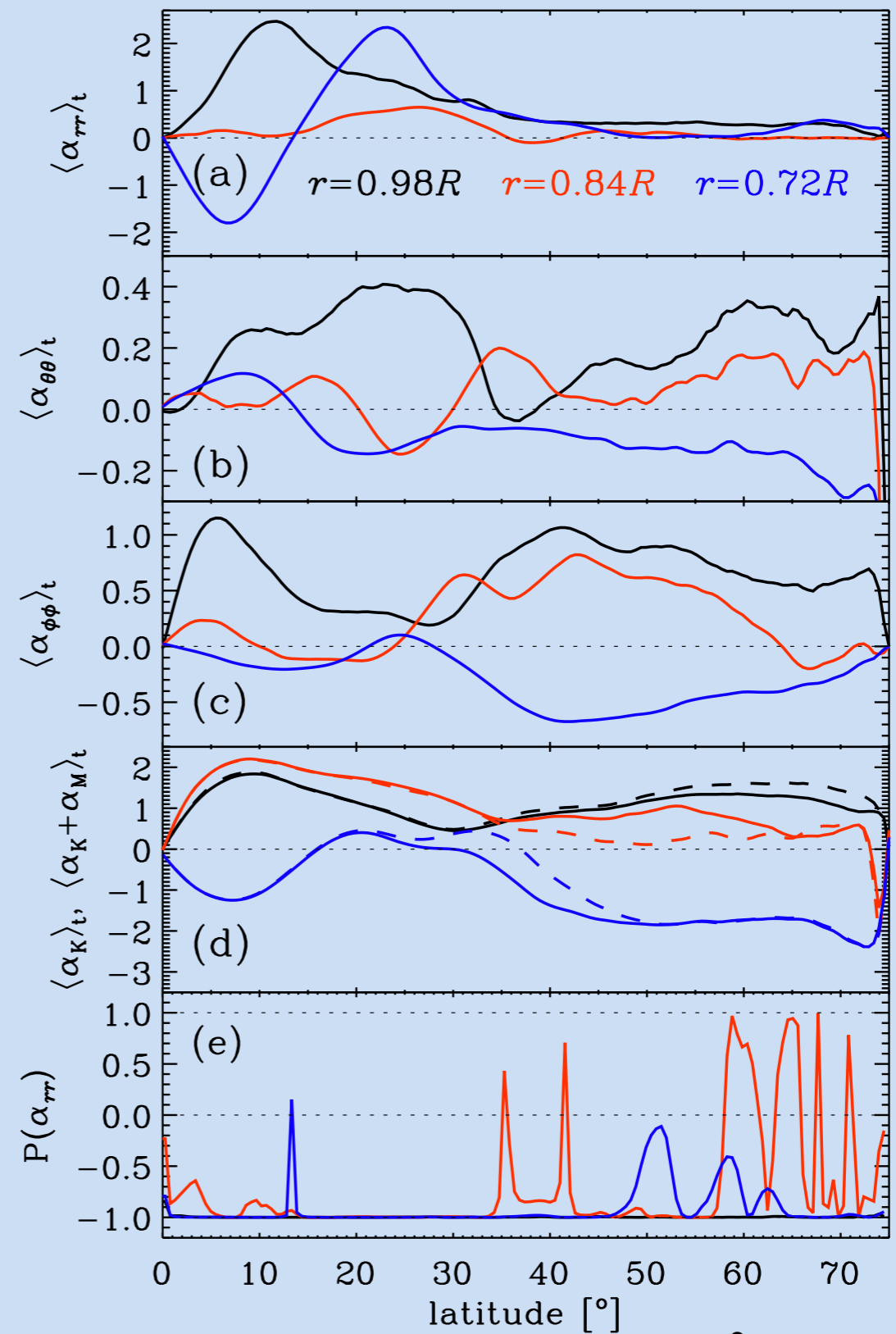
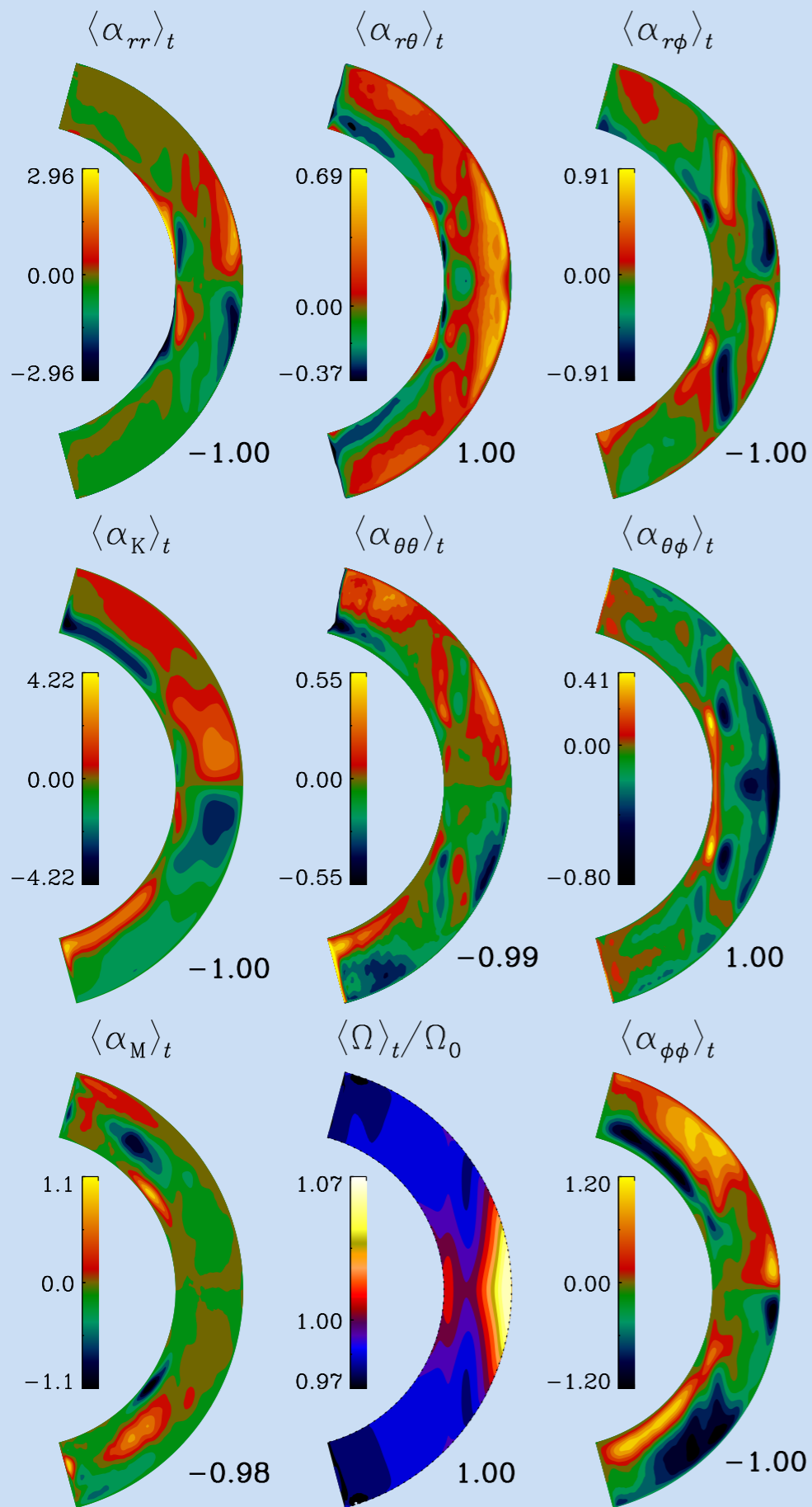
<https://github.com/pencil-code/pencil-code/>



Käpylä et al.
2012, 2013, 2016, 2017

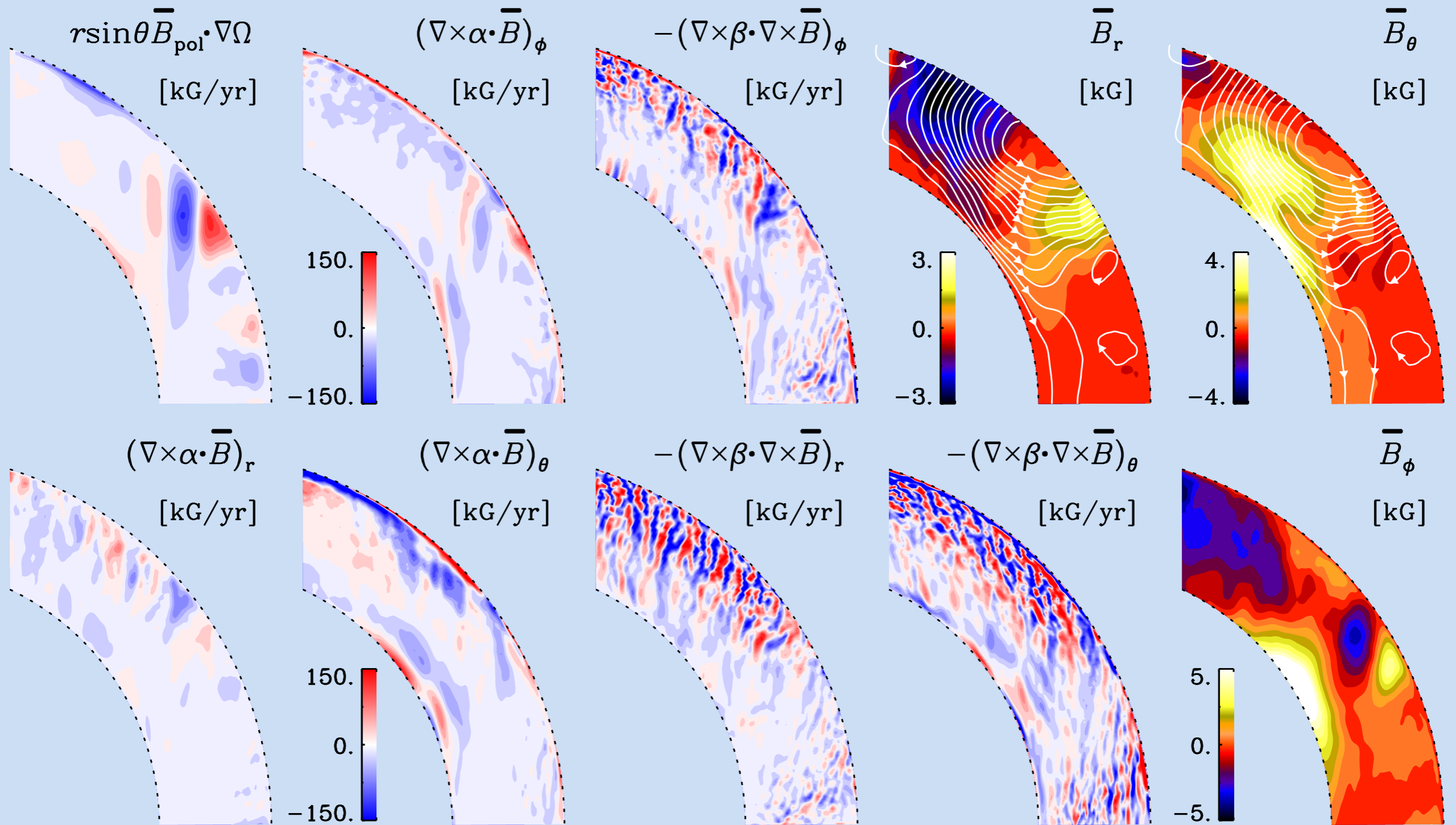
Warnecke et al.
2014, 2016

Results



$$P(\alpha_{ij}) = \frac{(\alpha_{ij}^{es})^2 - (\alpha_{ij}^{ea})^2}{(\alpha_{ij}^{es})^2 + (\alpha_{ij}^{ea})^2},$$

Magnetic field generation



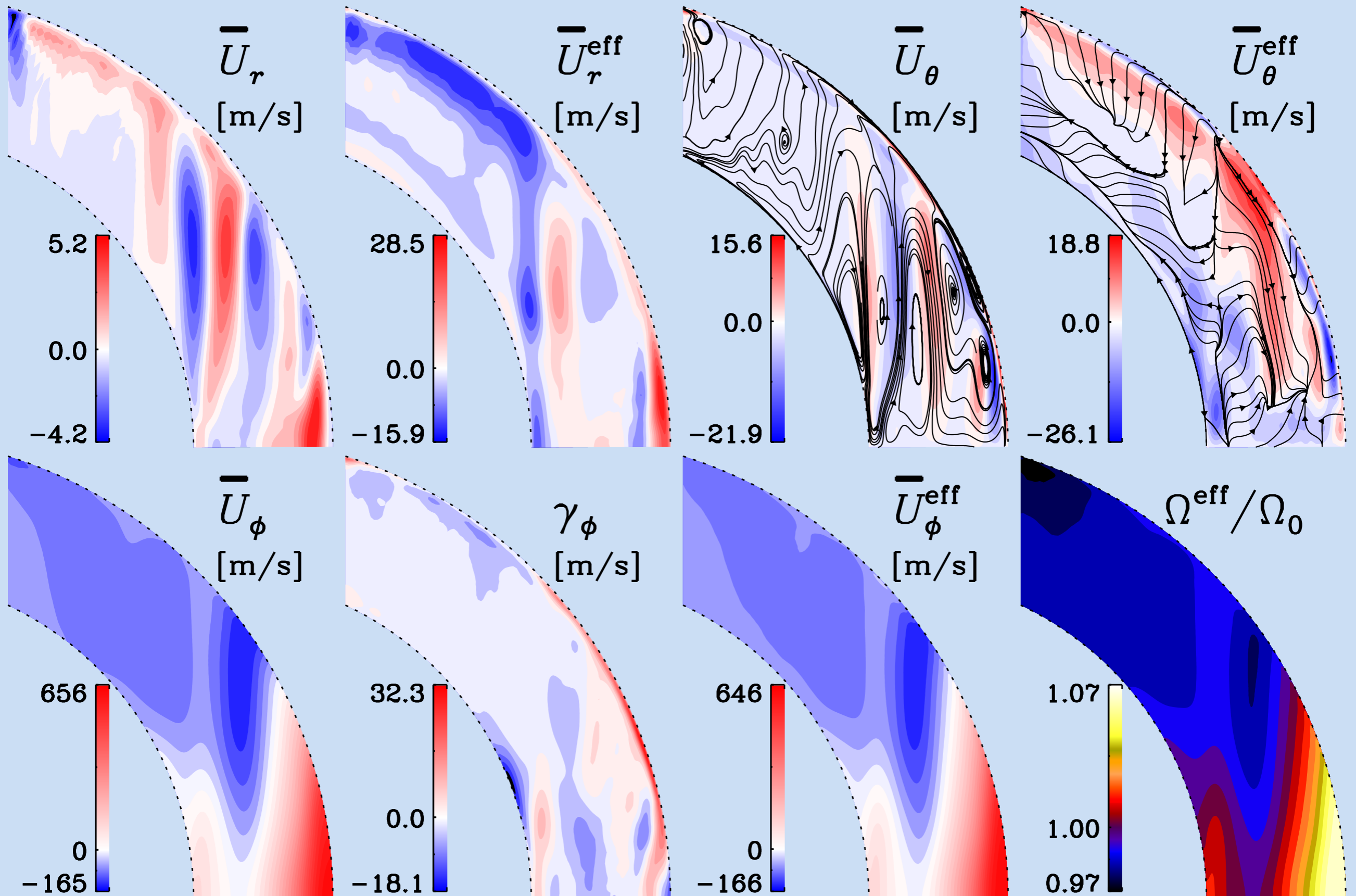
Turbulent pumping

$$\mathcal{E} = \alpha \cdot \bar{\mathbf{B}} + \gamma \times \bar{\mathbf{B}} - \beta \cdot (\nabla \times \bar{\mathbf{B}}) - \delta \times (\nabla \times \bar{\mathbf{B}}) - \kappa \cdot (\nabla \bar{\mathbf{B}})^{(S)}$$

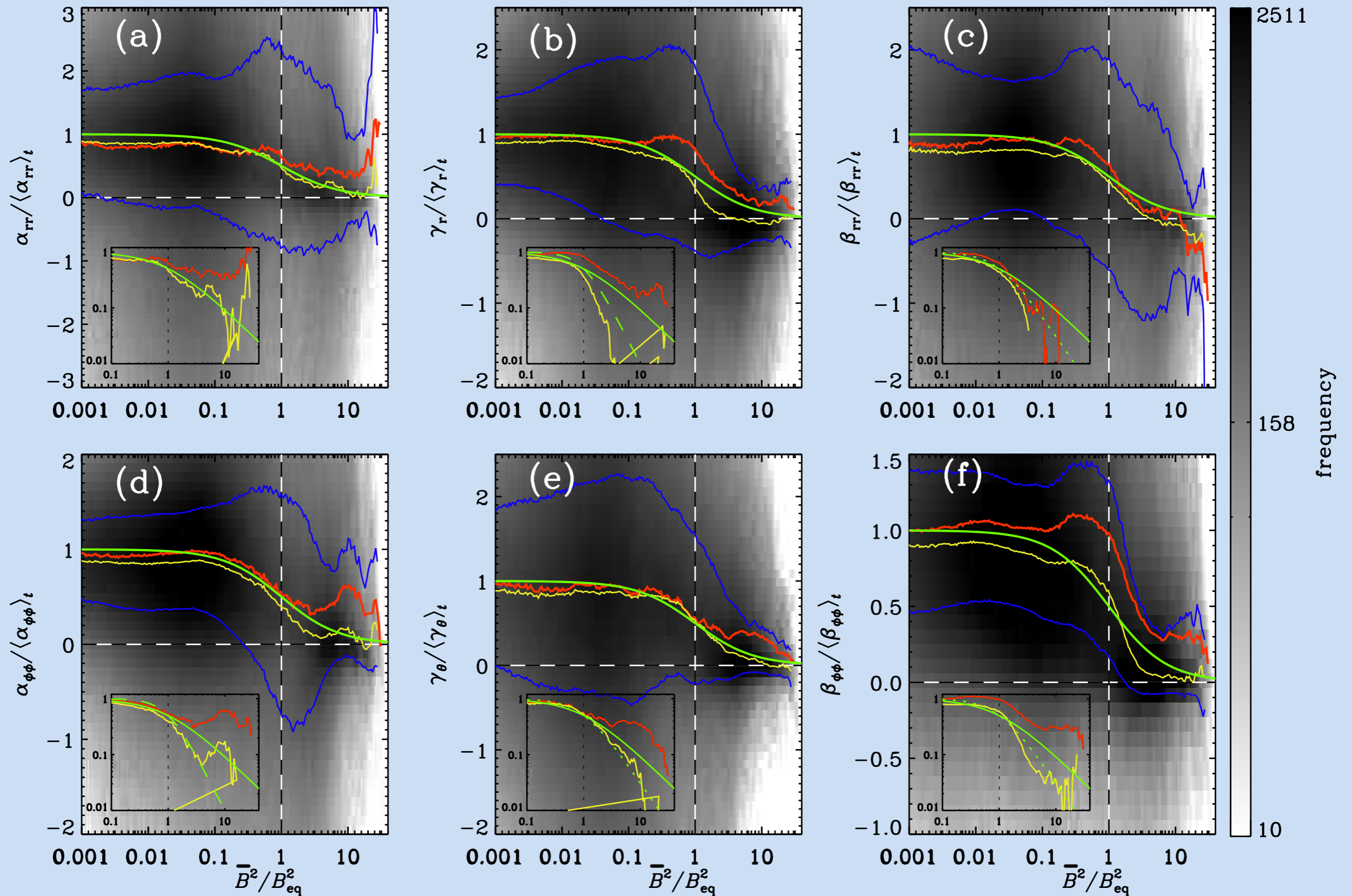
$$\frac{\partial \bar{\mathbf{B}}}{\partial t} = \nabla \times (\bar{\mathbf{u}} \times \bar{\mathbf{B}} + \overline{\mathbf{u}' \times \mathbf{B}'}) - \nabla \times \eta \nabla \times \bar{\mathbf{B}},$$

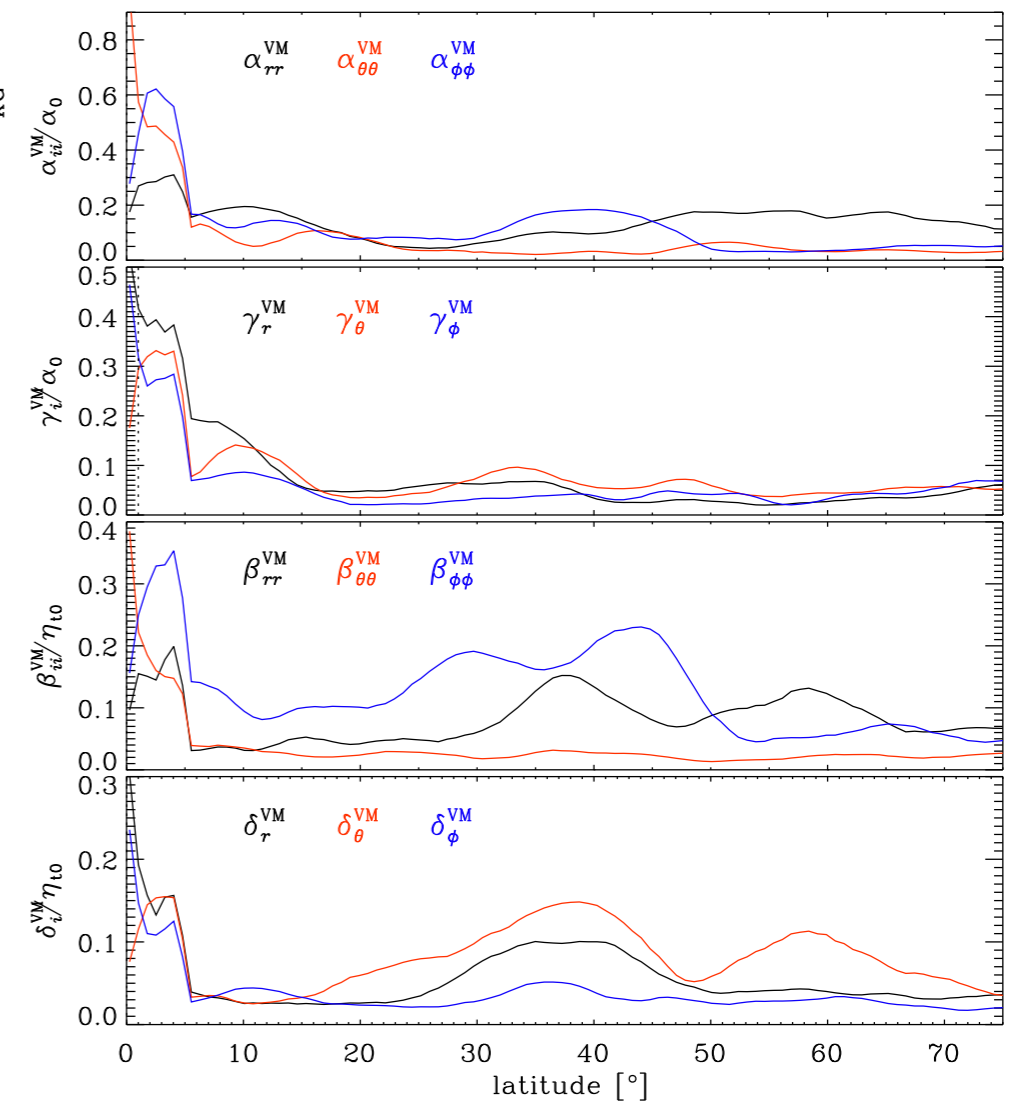
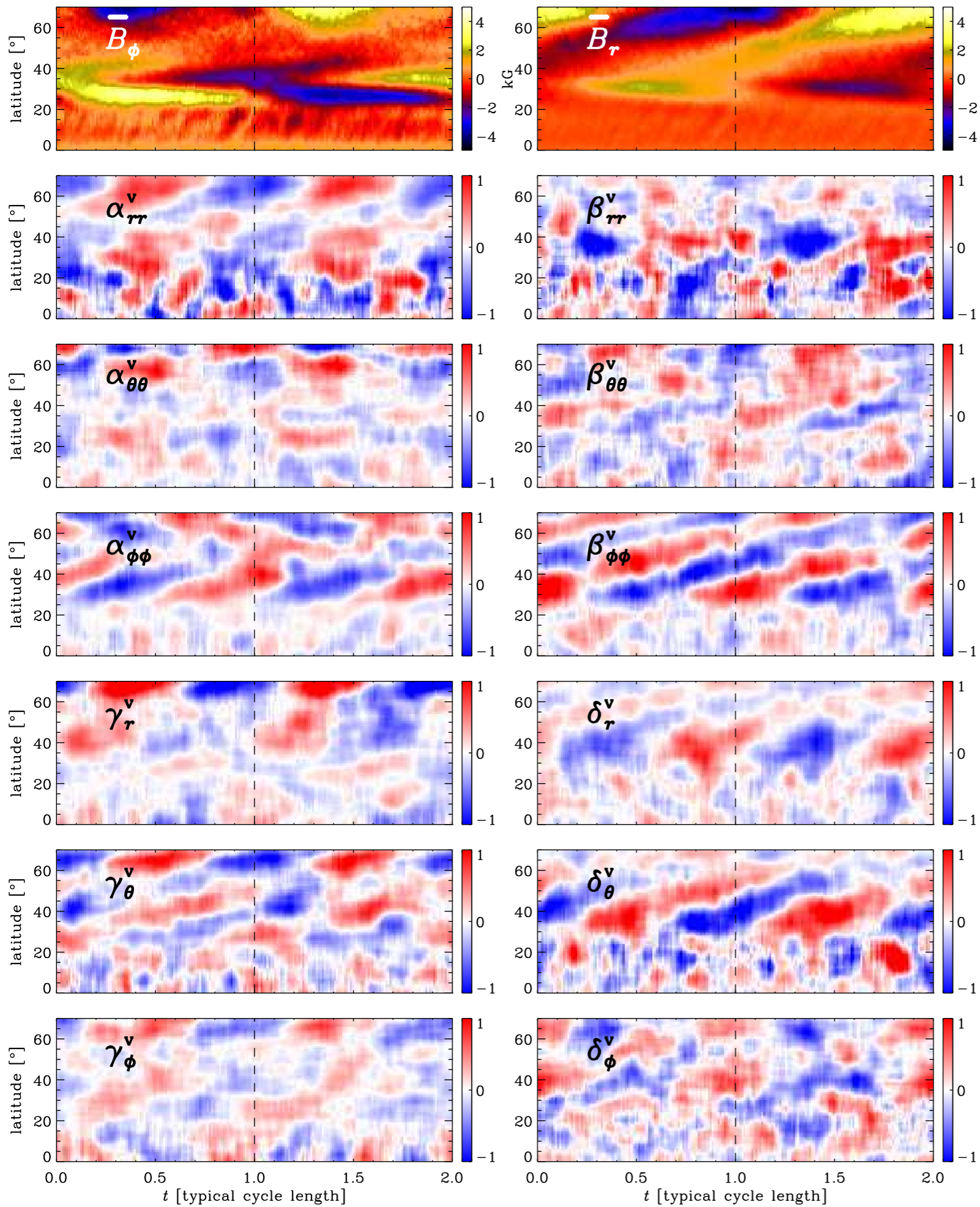
$$\partial_t \bar{\mathbf{B}}^{\text{pol}} = \nabla \times \left[\dots + \left(\gamma^{\text{pol}} + \bar{\mathbf{U}}^{\text{pol}} \right) \times \bar{\mathbf{B}}^{\text{pol}} \right] \quad (16)$$

$$\partial_t \bar{\mathbf{B}}^{\text{tor}} = \nabla \times \left[\dots + \left(\gamma^{\text{pol}} + \bar{\mathbf{U}}^{\text{pol}} \right) \times \bar{\mathbf{B}}^{\text{tor}} + \left(\gamma^{\text{tor}} + \bar{\mathbf{U}}^{\text{tor}} \right) \times \bar{\mathbf{B}}^{\text{pol}} \right] \quad (17)$$



Magnetic quenching





$$\alpha = \langle \alpha \rangle_t + \alpha^v.$$

$$\alpha_{ij}^v = \sqrt{\langle \alpha_{ij}^{v2} \rangle_t},$$

Conclusions

- Test-field method is one way to understand dynamo simulations.
- Alpha deviates from helicity expression.
- Complicated mixture of dynamo effects.
- Turbulent pumping changes significantly the eff. flow.
- Quenching does not depends analytical on B
- Strong cyclic variations of coefficients