

Computational future

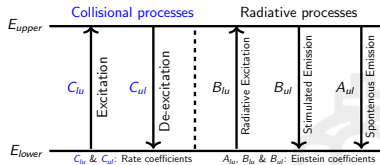
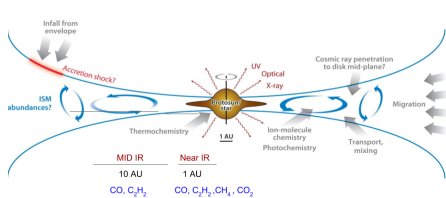
Gerrit C. Groenenboom

Theoretical Chemistry, Institute for Molecules and Materials,
Radboud University Nijmegen, The Netherlands

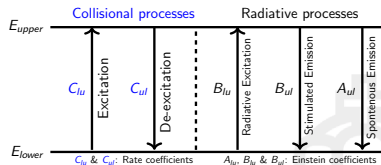
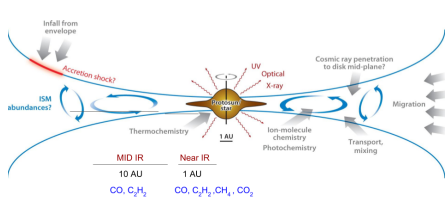
Groningen, 30 August 2024



Radiative transfer modeling - molecular collisions



Radiative transfer modeling - molecular collisions



Databases with **collisional rate constants** and spectroscopic data:

LAMDA Leiden Atomic and Molecular Database

Basecol Ro-vibrational Collisional Excitation Database and Utilities

EMAA Excitation of molecules and atoms for astrophysics

⇒ Data computed by **theoretical chemists/molecular physicist** ...

Also: Reaction rates, photodissociation rates cross section, spectroscopy, CIA





and what do they do ...

- Compute **potential energy surfaces** - **electronic Schrödinger equation**
- Nuclear motion
 - Often: Molecular Dynamics - classical mechanics
 - Here: **Quantum Dynamics/scattering** - **nuclear Schrödinger equation**
- Which problems are hard, which are easy? Will they do my molecule?

Typical rates: $10^{-7} - 10^{-8} \text{ s}^{-1}$

⇒ a little bit of **history**



Proc. R. Soc. A 256, 540-551 (1960)

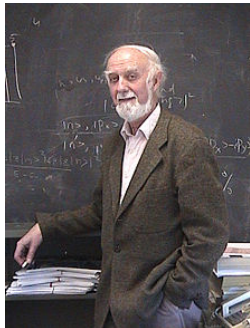
The theory of scattering by a rigid rotator

BY A. M. ARTHURS AND A. DALGARNO

Department of Applied Mathematics, The Queen's University of Belfast

(Communicated by D. R. Bates, F.R.S.—Received 8 February 1960)

A theory of scattering by a rigid rotator in which the coupling between the different energy levels of the rotator is taken into account is formulated and explicit expressions, which do not depend upon the magnetic quantum numbers, are obtained for various elastic and inelastic cross-sections. Several approximations are described, particular attention being paid to the scattering of heavy particles for which it is shown that at low temperatures the orientation-dependent part of the interaction may be more important than the spherically symmetric part. The scattering of low-energy electrons is also investigated and some representative equations are integrated numerically to illustrate the effect of the orientation dependence.



Proc. R. Soc. A 256, 540-551 (1960)

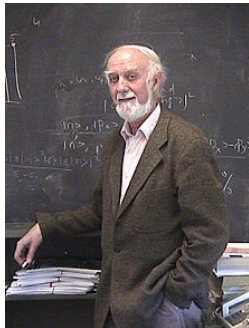
The theory of scattering by a rigid rotator

BY A. M. ARTHURS AND A. DALGARNO

Department of Applied Mathematics, The Queen's University of Belfast

(Communicated by D. R. Bates, F.R.S.—Received 8 February 1960)

A theory of scattering by a rigid rotator in which the coupling between the different energy levels of the rotator is taken into account is formulated and explicit expressions, which do not depend upon the magnetic quantum numbers, are obtained for various elastic and inelastic cross-sections. Several approximations are described, particular attention being paid to the scattering of heavy particles for which it is shown that at low temperatures the orientation-dependent part of the interaction may be more important than the spherically symmetric part. The scattering of low-energy electrons is also investigated and some representative equations are integrated numerically to illustrate the effect of the orientation dependence.



- British mathematician, Harvard University/Center for Astrophysics/ITAMP
- Time-independent Schrödinger equation for collisional cross sections
- Key of 1960 paper: total angular momentum representation.

Proc. R. Soc. A 256, 540-551 (1960)

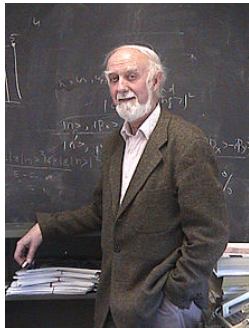
The theory of scattering by a rigid rotator

BY A. M. ARTHURS AND A. DALGARNO

Department of Applied Mathematics, The Queen's University of Belfast

(Communicated by D. R. Bates, F.R.S.—Received 8 February 1960)

A theory of scattering by a rigid rotator in which the coupling between the different energy levels of the rotator is taken into account is formulated and explicit expressions, which do not depend upon the magnetic quantum numbers, are obtained for various elastic and inelastic cross-sections. Several approximations are described, particular attention being paid to the scattering of heavy particles for which it is shown that at low temperatures the orientation-dependent part of the interaction may be more important than the spherically symmetric part. The scattering of low-energy electrons is also investigated and some representative equations are integrated numerically to illustrate the effect of the orientation dependence.



- British mathematician, Harvard University/Center for Astrophysics/ITAMP
- Time-independent Schrödinger equation for collisional cross sections
- Key of 1960 paper: total angular momentum representation.

This turns computation of rates into many independent jobs:

- Total energies E , e.g., 100 values between 1 cm^{-1} and 1000 cm^{-1}
- Total angular momentum: J , e.g., 0-100
- **Parity** odd/even inversion symmetry

⇒ Example: $100 \times 100 \times 2 = 20\,000$ jobs



WORKSHOP PARTICIPANTS

Workshop participants. *Left to right:* L. Thomas, N. Harvey, J. Light, D. Secrest, G. Parker, A. Askar, D. Malik, R. Gordon, M. Alexander, R. Walker, M. Redmon, A. Allison, F. Krogh, H. Rabitz, B. Johnson, H. Shin, T. Schmalz, Y. Yung, D. Truhlar, W. Lester, Jr.

Workshop NRCC Proceedings No. 5, Argonne National Lab, 1979

- **B. Robert Johnson**, [Log-derivative and renormalized numerov algorithm](#)
- **Millard H. Alexander**, [HIBRIDON](#)

Hibridon: a program suite for time-independent non-reactive quantum scattering calculations
M.H. Alexander, P.J. Dagdigian, H.-J. Werner, J. Kłos, B. Desrousseaux, G. Raffy, and F. Lique
Comp. Phys. Commun., 289, 108761 (2023)
<https://github.com/hibridon/hibridon>

Not at this workshop: **Sheldon Green** - [MOLSCAT](#) code - now on Github by **Jeremy Hutson**

Collinear He+H₂(*v*) → He+H₂(*v'*), Secret & Johnson

"Simplest coupled channels problem"



Hamiltonian for He+H₂

$$\hat{H} = \underbrace{\hat{H}_0}_{\text{H}_2 \text{ Hamiltonian}} - \underbrace{\frac{\hbar^2}{2\mu} \frac{\partial^2}{\partial R^2}}_{\text{kinetic energy}} + \underbrace{V(R, r)}_{\text{interaction potential}}$$

Time-independent Schrödinger equation, energy E

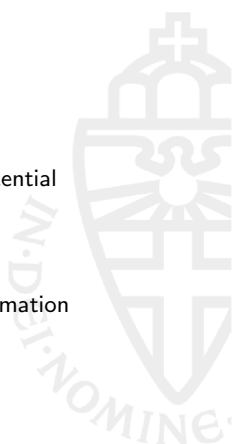
$$\hat{H}\Psi(R, r) = E\Psi(R, r)$$

Vibrational wave functions for H₂ — harmonic oscillator approximation

$$\hat{H}_0\phi_v(r) = \epsilon_v\phi_v(r), \quad v = 0, 1, 2$$

Coupled channel expansion

$$\Psi(R, r) = \sum_{v=0}^2 \phi_v(r) \underbrace{c_v(R)}_{\text{Need to solve}}$$



Johnson's Renormalized Numerov algorithm

Coupled channel expansion

$$\Psi(R, r) = \sum_{v=0}^{v_{\max}} \phi_v(r) c_v(R), \quad \text{expansion coeffs: } \mathbf{c}(R) = \begin{pmatrix} c_0(R) \\ c_1(R) \\ c_2(R) \end{pmatrix}$$

Coupled channels equation

$$\frac{\partial^2}{\partial R^2} \mathbf{c}(R) = \mathbf{W}(R) \mathbf{c}(R)$$

Main steps in computation

- 1 Choose grid $\mathbf{c}(R_i)$ for $R_i = R_0 + i\Delta_R$ with $i = 1, 2, \dots, n_R$.
- 2 Compute matrix elements of interaction potential

$$V_{v'v}(R_i) = \int_0^{\infty} \phi_{v'}^*(r) V(R_i, r) \phi_v(r) dr.$$

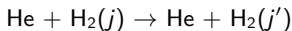
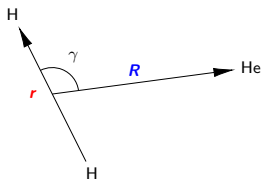
- 3 Solve CC Eq. with Johnson's Renormalized Numerov method

$$\mathbf{c}(R_{i-1}) = \mathbf{Q}_i \mathbf{c}(R_i)$$

$$\mathbf{Q}_{i+1} = \left[2 + \frac{10\Delta_R^2 \mathbf{W}_i}{12} - \left(1 - \frac{\Delta_R^2 \mathbf{W}_{i-1}}{12} \right) \mathbf{Q}_i \right]^{-1} \left(1 - \frac{\Delta_R^2 \mathbf{W}_{i+1}}{12} \right).$$



Atom-diatom rotationally inelastic scattering



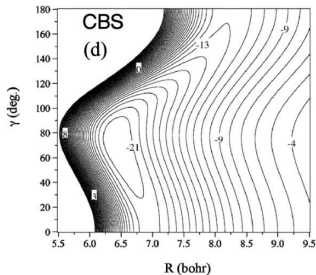
Rigid rotator: fix bondlength $r = |\mathbf{r}|$

Channel functions for rigid rotator approximation:

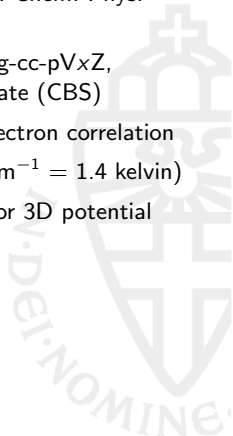
$$|(j)JM_J\rangle = \sum_{m_j=-j}^j \sum_{m_l=l}^l Y_{j,m_j}(\hat{r}) Y_{j,m_l}(\hat{R}) \underbrace{\langle jm_j m_l | JM_J \rangle}_{\text{Clebsch-Gordan coefficients}}$$

- j - rotational quantum number H_2 - we care about this one
- J - total angular momentum, good quantum number (great!)
- M_J - projection - results independent of this quantum number (great!)
- m_j and m_l - summed over (thanks Alex Dalgarno!)
- $p = (-1)^{j+l}$, parity good quantum number (great!)
- l - “end-over-end rotation” increases number of channels (bummer)

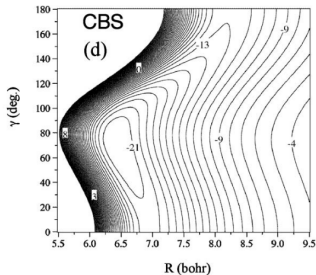
Example He+CO potential



- Peterson and McBane, J. Chem. Phys. **123**, 084314 (2005)
- One-electron basis: d-aug-cc-pVxZ, $x = D, T, Q \rightarrow$ extrapolate (CBS)
- CCSD(T) - treatment electron correlation
- $D_e = 22.34 \text{ cm}^{-1}$ ($1 \text{ cm}^{-1} = 1.4 \text{ kelvin}$)
- about 1000 geometries for 3D potential



Example He+CO potential

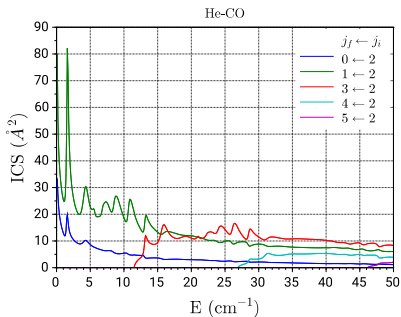


- Peterson and McBane, J. Chem. Phys. **123**, 084314 (2005)
- One-electron basis: d-aug-cc-pVxZ, $x = D, T, Q \rightarrow$ extrapolate (CBS)
- CCSD(T) - treatment electron correlation
- $D_e = 22.34 \text{ cm}^{-1}$ ($1 \text{ cm}^{-1} = 1.4 \text{ kelvin}$)
- about 1000 geometries for 3D potential

“Rotational excitation of CO by He impact”, Thomas, Kraemer, Dierksen, Chem. Phys. **51**, 131 (1980)

- CISD/TZ (configuration interaction)
- 50% uncertainty in well depth
- Coupled channels and IOS (infinite order sudden) with MOLSCAT

Cross section $\text{He} + \text{CO}(j_i = 2) \rightarrow \text{He} + \text{CO}(j_f)$



Calculation in 2024 (by me)

- 500 energies, E step of 0.02 cm^{-1}
 - CO rotational states: $j_{\text{max}} = 10$
 - Partial waves: $J_{\text{max}} = 10$
 - Radial propagation steps: $n_R = 80$
 - Number of coupled channels: 10-120
 - (5-60 in parity adapted basis)
- CPU time: about **one minute** on PC (6 cores) [Scilab script]
 - **Experimental observation** of $j = 0 \rightarrow 1$, five resonances for $E = 0 - 25 \text{ cm}^{-1}$, A. Bergeat *et al.*, Nature Chemistry 2015

Scaling of CPU time with # coupled channels (n)

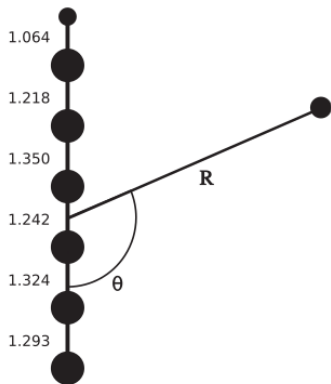
- Matrix multiplication/inversion: CPU time: $t \propto n^3$
- CPU time on PC (single core) for $n = 1000$: $t \approx 0.1$ sec.
- Main frame **year 1980**: same operation about **one hour!**

Rotational energy

$$E_j = B j(j + 1) - \dots$$

	B (cm ⁻¹)	#($E_j < 1000$)
H ₂ (¹ Σ_g^+)	61	4/2=2
CO(¹ Σ^+)	1.7	24
C ₆ H(² Π)	0.046	147×2=294

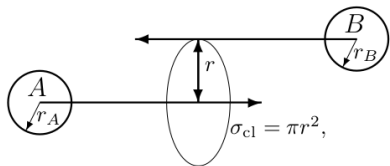
How about l ?



Walker, Lique, Dawes, MNRAS **473**, 1407 (2018)

Estimating l , impact parameter

If molecules were hard spheres



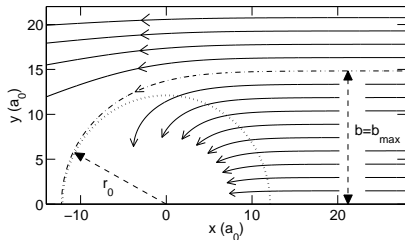
“End-over-end” angular momentum of complex:

$$l = \mu v r$$

$$\mu = \frac{m_A m_B}{m_A + m_B}$$

$$E = \frac{1}{2} \mu v^2$$

With long range interactions



Example:

- $\mu = m_{He}$
- $E = 1000 \text{ cm}^{-1}$
- $b = 10 a_0$
- $\Rightarrow l = 80$

Coupled states approximation

- Introduce body-fixed angular momentum projection quantum number: Ω
- Unitary transformation from I to Ω :

$$|(jI)JM_J\rangle = \sum_{\Omega} |j\Omega)JM_J\rangle U_{\Omega,I}^{(j,J)}$$

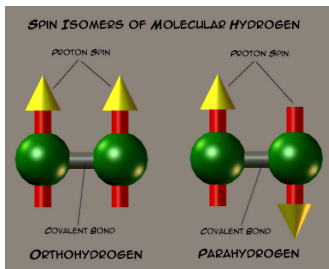
- Potential matrix elements independent of Ω
- Kinetic energy coupling $T_{\Omega,\Omega\pm 1}$ (Coriolis coupling)
- **Coupled States Approximation**: neglect $T_{\Omega,\Omega\pm 1}$ coupling
- \Rightarrow often good approximation at high collision energy

Many papers in the 1970's, Sheldon Green, Don Kouri, R. T Pack



What about collisions with H₂ and HD instead of He?

Ortho- and para- H₂



https://en.wikipedia.org/wiki/Spin_isomers_of_hydrogen

- Para-H₂: channels: $j = 0, 2, 4, \dots$
- Ortho-H₂: channels: $j = 1, 3, \dots$
- For $j = 0$, quadrupole moment = 0
- $m_{\text{H}_2}/m_{\text{He}} = 1/2 \Rightarrow$ de Broglie wavelength $\times \sqrt{2}$

Switching to HD?

- No more symmetry: $j = 0, 1, 2, \dots$
- Rotational constant: $B = 46 \text{ cm}^{-1}$ (vs 61 cm^{-1} for H₂)



Including vibrations: He+CO(v, j)

C. Cecchi-Pestellini, E. Bodo, N. Balakrishnan, A. Dalgarno, ApJ, **571**, 1015 (2002)

TABLE 3
CO ROTATIONAL ENERGIES

J	$\Delta\epsilon_{v,J}/\text{cm}^{-1}$
1.....	3.85
2.....	11.54
3.....	23.07
4.....	38.45
5.....	57.68
6.....	80.75
7.....	107.66
8.....	138.42
9.....	173.03
10.....	211.48
11.....	253.77
12.....	299.91
13.....	349.90
14.....	403.73

TABLE 1
CO VIBRATIONAL ENERGIES

v	$\Delta\epsilon_v/\text{cm}^{-1}$
0.....	0.0
1.....	2144.7
2.....	4263.3
3.....	6356.0
4.....	8422.6
5.....	10463.1
6.....	12477.7

How did they handle this?

- 4 CPU Compaq Digital Alpha
- Adapted MOLSCAT code
- Collision energies: 5 to 600 cm^{-1}
- Rotational transitions: coupled channels
- Vibrational transitions: infinite-order sudden (IOS)

Similar to $He+CO(v,j)$?



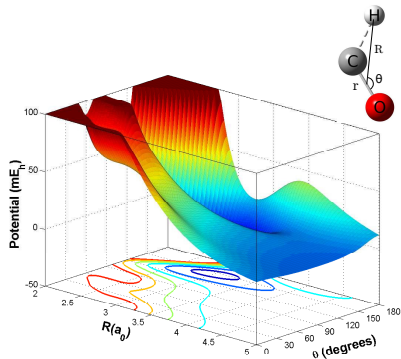
Similar to He+CO(v, j)?

- Sensitivity analysis: H+CO(v, j) key - Ewine van Dishoeck (Leiden)
- Dutch Astrochemistry Network I - funding for laboratory astrochemistry

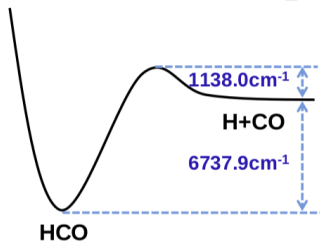


, Ph project Lei Song on H+CO(v, j), 2012-2016

H+CO potential, 3D



- 3-D potential, with MOLPRO
- 4400 *ab initio* points
- UCCSD(T), d-aug-cc-pVnZ
- extrapolated from $n = 3, 4, 5$ to CBS
- $D_e = 0.8 \text{ eV} \sim 6500 \text{ cm}^{-1}$



Lei Song, A. van der Avoird, GCG, J. Phys. Chem. A, **117**, 7571 (2013)

Quality of potential: HCO bound states ($J = 0$)

(ν_1, ν_2, ν_3)	Experimental	Calculated	Deviation
(0,0,1)	1080.76	1079.6	-0.11%
(0,1,0)	1868.17	1871.6	0.18%
(0,0,2)	2142	2145.2	0.15%
(1,0,0)	2434.48	2428.5	-0.25%
(0,1,1)	2942	2950.0	0.27%
(0,0,3)	3171	3186.2	0.48%
(1,0,1)	3476	3464.1	-0.34%
(0,2,0)	3709	3719.0	0.27%
(0,1,2)	3997	4013.1	0.40%
(0,0,4)	4209 ^a	4195.7	-0.32%
(1,1,0)	4302	4301.5	-0.01%
(1,0,2)	4501 ^a	4478.7	-0.50%
(2,0,0)	4570 ^a	4546.0	-0.53%
(0,2,1)	4783.2	4797.2	0.29%

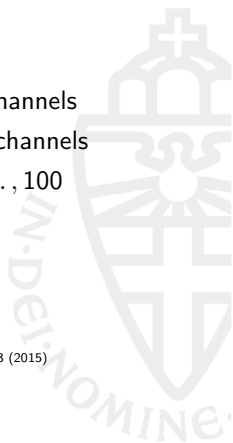
Lei Song, A. van der Avoird, GGC, J. Phys. Chem. A, **117**, 7571 (2013)

Coupled channels scattering calculation $H+CO(v, j)$

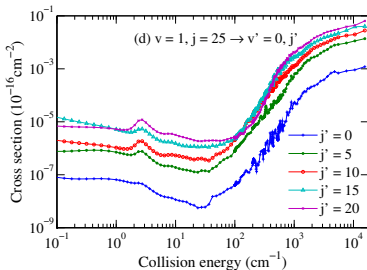
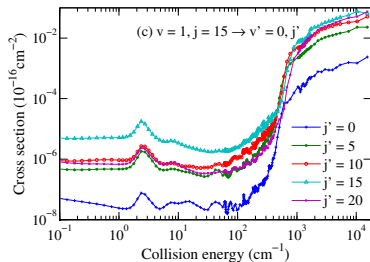
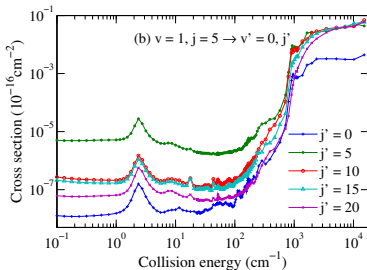
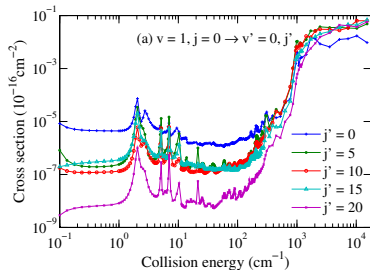
	Big basis	small basis
v	j_{\max}	j_{\max}
0	79	61
1	70	52
2	59	42
3	45	22
4	45	
5	30	
6	30	
7	25	
8	25	

- **Big basis** \sim **11,000** channels
- **Small basis** \sim **4,300** channels
- Partial waves $J = 0, \dots, 100$
- R -grid: : 3 - 40 a_0

Lei Song, N. Balakrishnan, A. vd Avoird, T. Karman, and GCG, J. Chem. Phys., **142**, 204303 (2015)



$\text{H} + \text{CO}(v, j) \rightarrow \text{H} + \text{CO}(v', j')$ - several months later



Lei Song, N. Balakrishnan, A. vd Avoird, T. Karman, and GCG, *J. Chem. Phys.*, **142**, 204303 (2015)

Boltzmann average of cross section

$$r_{vj \rightarrow v'j'}(T) = \left(\frac{8k_B T}{\pi \mu} \right)^{1/2} \frac{1}{(k_B T)^2} \int_0^\infty \sigma_{vj \rightarrow v'j'}(E) e^{-E/k_B T} E dE$$

Vibrational quenching

$$r_{v \rightarrow v'}(T) = \frac{\sum_{j,j'} g_j e^{-E_{vj}/k_B T} r_{vj \rightarrow v'j'}(T)}{\sum_j g_j e^{-E_{vj}/k_B T}}$$

Extrapolation used in astrochemical applications

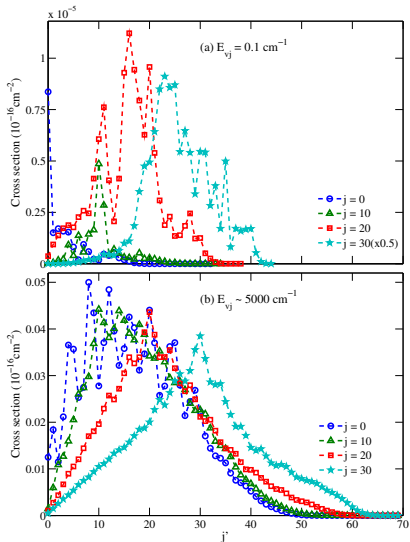
$$r_{vj \rightarrow v'j'}(T) = P_{vv'}(T) r_{0,j \rightarrow 0,j'}(T)$$

with

$$P_{vv'}(T) = \frac{r_{v \rightarrow v'}(T) \sum_j g_j e^{-E_{vj}/k_B T}}{\sum_j g_j e^{-E_{vj}/k_B T} \sum_{j'} r_{0,j \rightarrow 0,j'}(T)}$$

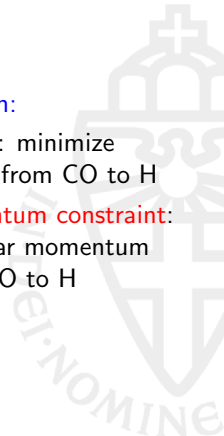


CO(v, J') product distributions

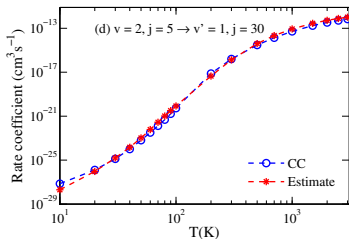
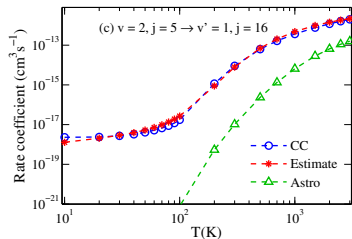
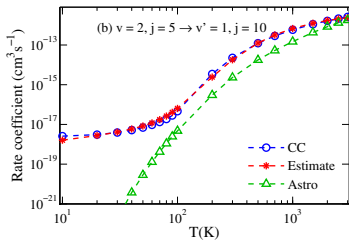
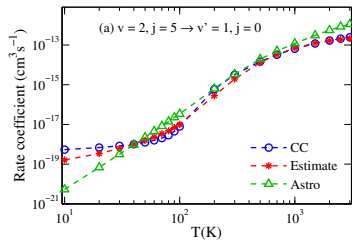


Competition between:

- **Energy gap law:** minimize energy transfer from CO to H
- **Angular momentum constraint:** minimize angular momentum transfer from CO to H



Collision rates H+CO: $r_{v=2,j \rightarrow v'=1,j'}(T)$

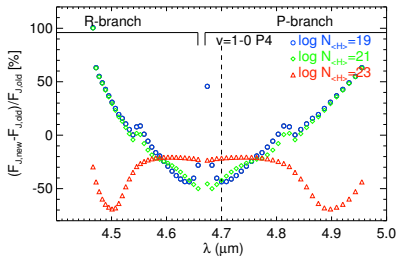


- Song, Balakrishnan, Walker, Stancil, Thi, Kamp, vd Avoird, and GCG, *Astrophys. J.*, 2015
- Astro: W. F. Thi, Inga Kamp, et al., *A&A*, 551, A49 (2013)
- Extrapolation formula works well starting with $(v = 1, j) \rightarrow (v' = 0, j')$ data.

H+CO: application to astrophysical models

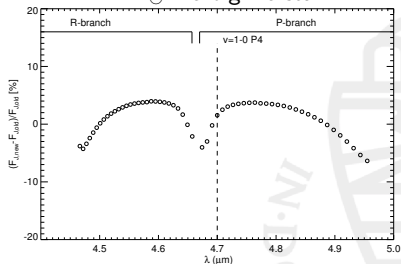
Slab model

$T = 200\text{K}$, $\log n_{H, H_2} = 9$, $\log n_{H_e} = 8$, $\log n_e = 5$



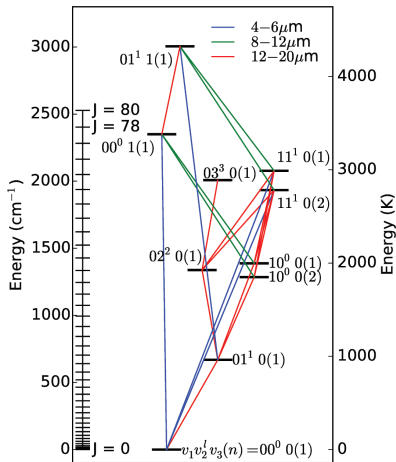
Disk model

$2.2 M_{\odot}$ Herbig Ae star



Lei Song, N. Balakrishnan, Kyle Walker, Phillip Stancil, [Wing-Fai Thi](#), [Inga Kamp](#), Ad van der Avoird, GCG, *Astrophys. J.*, **813**, 96 (2015)

CO₂ vibrational transitions



CO₂ vibrations

	(cm ⁻¹)	mode
ν_1	1333	symmetric stretch
ν_2	667	bend
ν_3	2349	asymmetric stretch

States: (ν_1, ν_2^l, ν_3)

Vibrational angular momentum:

$$l = \nu_2, \nu_2 - 2, \dots, -\nu_2$$

Symmetric stretch not IR active,
but: $2 \times 667 = 1334$

Fermi resonance: (10⁰)-(02⁰)

"CO₂ infrared emission as a diagnostic of planet-forming regions of disks",

A. D. Bosman, S. Bruderer, E. F. van Dishoeck, A&A 601, A36 (2017)



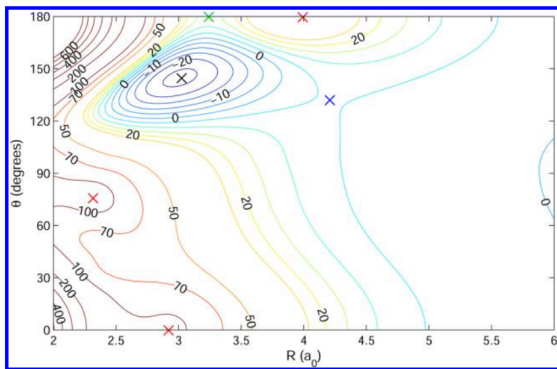
Taha Selim, **DAN II**
Chemical Physics part of the problem
PhD defense October, **2024**



Jelena Boskovic (**DAN III**)
with Floris van der Tak, RUG
Astrochemical applications
Started March **2024**



Conical intersections and Renner-Teller coupling in HCO



- **Red crosses** - conical intersections (Lowest $\sim 11\,000\text{ cm}^{-1}$)
- **Green cross** - Renner Teller coupling at linear geometry, 1747 cm^{-1}

Lei Song, Ad van der Avoird, and Gerrit C. Groenenboom *J. Phys. Chem. A*, **117**, 7571 (2013)

Ongoing project Kim Steenbakkens, Sandra Brünken:

Spectroscopy of CCH^+ - Renner Teller coupling, ground state $^3\Pi$, low lying $^3\Sigma^-$

Conclusion - computational future

- Steady progress electronic structure methods
 - electron correlation: CCSD(T), CCSDT(Q), ...
 - basis set: aug-cc-pVTZ, aug-cc-pVQZ, ...
 - Open-shell systems, conical intersections \Leftarrow hard
- Fitting higher dimensional potentials: machine learning, AI methods
 \Rightarrow not always the biggest problem, but lots of interest
- Dynamics: rotation-vibration, larger molecules \Leftarrow ongoing efforts
- Nonadiabatic dynamics \Leftarrow hard
- Software advances - data science techniques \Leftarrow many opportunities
- Price/flop keeps going down, great - but: power usage is concern
- Theoretical advances, bonus - similar to hardware improvements

Thank you! Questions?

