



Laboratoire d'Études Spatiales et d'Instrumentation en Astrophysic

On the relation between observed and predicted global seismic quantities in red-giant stars

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Preliminary definitions & problematic

• Global seismic parameters



 $u_{\max} \quad \begin{array}{l} ext{is the frequency of the} \\ ext{max} & ext{maximum height in the power} \\ ext{spectrum} \end{array}$

 $\nu_{\rm c}$ is the cut-off frequency

 $\Delta
u$ large separation

For gravity modes, we also measure period $\Delta \Pi$ spacing

One can also use the maximum amplitude (i.e. at the v_{max})

)

max

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Preliminary definitions & problematic

• Global seismic parameters

Theory (scaling relations): asymptotic developments, etc...



 $\Delta
u \propto ig\langle
ho ig
angle^{1/2}
onumber
u_{
m max} \propto rac{g}{\sqrt{T_{
m eff}}}
onumber
onumber \Delta\Pi \propto ig\langle
ho ig
angle_{
m core}^{-1/2}
onumber
onumber$

From all those information you can:

✓ obtain «first-order» contraints for modelling

 \checkmark perform ensemble asteroseismology through the use of scaling relations

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 $egin{aligned} &\Delta
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ho ig
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ho ig
angle_{
m core}^{-1/2} \ &A_{
m max} \propto igg(rac{L}{M}ig)^lpha \end{aligned}$

Of course it is not so simple...

 \checkmark In fact, there is always quite a leap between observables and theoretical predictions

 \checkmark This is particular crucial for red-giants

Outline

Period spacing of red-giant stars
 Asymptotic relation for deriving the period spacings

Amplitudes of red-giant stars

✓ Evidence of non-adiabatic effects

 \checkmark The need for both radial velocities and photometric observations

✓ Why red-giant stars are so special? Mixed modes The dispersion relation reads $k_r^2 = \frac{1}{\omega^2 c_s^2} \left(\omega^2 - N^2\right) \left(\omega^2 - S_\ell^2\right)$



There are two propagative cavities (kr² > 0): • upper cavity: acoustic modes can exist $\omega >> S_{\ell}$ and $\omega >> N$

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There are two propagative cavities (kr² > 0): • upper cavity: acoustic modes can exist $\omega >> S_{\ell}$ and $\omega >> N$

• inner cavity: gravity modes

 $\omega << S_{\ell}$ and $\omega << N$

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There are two propagative cavities (kr² > 0): • upper cavity: acoustic modes can exist $\omega >> S_{\ell}$ and $\omega >> N$

• inner cavity: gravity modes $\omega << S_\ell \quad and \quad \omega << N$

• intermediate region: modes are evanescent, it couples the cavities

modes behave as p-mode in the upperlayers and as a g-mode in the inner layers: they are mixed

√ Why red-giant stars are so special? Mixed modes

Mode inertia is a good indicator of the degree of mixing 1





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Mode inertia is a good indicator of the degree of mixing 1





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From an observational point of view:

So the quest for mixed modes
 began

 Mixed modes (p-dominated) have been detected by both CoRoT and Kepler



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 So the quest for mixed modes began

 Mixed modes (p-dominated) have been detected by both CoRoT and Kepler







Bedding, Mosser, Huber, et al. Nature, (2011)

The first issue was to discriminate between Clump stars and RGB stars

RGB (red giant branch): hydrogen shell burning

Red clump stars: helium central burning

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 From mixed modes frequencies one can deduce an «effective» period spacing

 it depends mainly on the core of the giants

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Red clump stars: helium central burning

 From mixed modes frequencies one can deduce an «effective» period spacing

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 it depends mainly on the core of the giants



 Can we use the observed period spacings and compare them to the theory?

 No: asymptotic period spacings concerns gravity modes (not p-dominated mixed modes)

 We need some theoretical framework for relating «effective» and asymptotic value of the period spacing



Following the Unno et al. (1989) approach to get the phase relation for asymptotic mixed modes

Goupil et al. derived a asymptotic relation for mixed modes

$$\nu = \nu_p + \frac{\Delta \nu}{\pi} \arctan \left[q \tan \pi \left[\frac{1}{\Delta P \nu} - \epsilon_g \right] \right]$$

Goupil, Belkacem, Marques, Mosser et al. in prep

frequency of 'pure' p modes

$$\nu = \nu_p + \frac{\Delta \nu}{\pi} \arctan \left[q \tan \pi \left[\frac{1}{\Delta P \nu} - \epsilon_g \right] \right]$$

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large separation

$$\nu = \nu_p + \frac{\Delta \nu}{\pi} \arctan\left[q \tan \pi \left[\frac{1}{\Delta P \nu} - \epsilon_g\right]\right]$$

coupling factor

Goupil, Belkacem, Marques, Mosser et al. in prep

$$u =
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Mosser et al. in prep
REAL period spacing

$$u = \nu_p + \frac{\Delta \nu}{\pi} \arctan\left[q \tan \pi \left[\frac{1}{\Delta P \nu} - \epsilon_g\right]\right]$$

Goupil, Belkacem, Marques, Mosser et al. in prep

B. Mosser, M.J. Goupil, K. Belkacem et al. (2012)



Through a fit of the Period échelle diagram

Period spacing

frequency versus mode period modulo the period spacing

Period spacing vs large separation



period spacing from models computed with the asymptotic relation

good agreement with the observations

Mosser et al. (2012)



We have direct constraints on the core of stars

For more insight into the period spacings of redgiant stars see the recent paper of J. Montalban (Montalban et al. 2013, arXiv:1302.3173)

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Amplitudes of red giants

Principle: driven and damped oscillator

Mode amplitude is a balance between mode driving and damping

$$\frac{\mathrm{d}}{\mathrm{d}t}A^2 = \mathcal{P} - \eta A^2$$





Amplitudes of red giants: observations

From ground-based observations we have only a few stars

Observations vs Theory: not satisfactory for red giants

• From space-borne missions CoRoT and Kepler





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RGB stars

Red Clump stars

 stars with «depressed» dipolar modes So what do we know about mode amplitudes in red giants and related scaling relations?

So what do we know about mode amplitudes in red giants and related scaling relations?

Theory:

• Kjeldsen & Bedding (1995)

Derived from Christenssen-Dalsgaard & Frandsen (1983) calculations

Based on a equipartition principle (Goldreich & Keeky 1977)

• Houdek et al (1999)

Based on Balmforth (1992)'s formulation, closely derived from Goldreich & Keeky (1977) theory. Assumes implicitly a Gaussian eddy-time correlation, no thermal forcing

• Samadi et al (2007)

Based on Samadi & Goupil (2001 formalism), which is a generalization of Goldreich & Keeky (1977) theory. Thermal forcing; free choice for eddy-time correlation

Obsevations (empirical approaches):

• Kjeldsen & Beeding (2011)

Phenomenological approach. Assumes that the squared amplitude is proportional to the power in the granulation spectrum

• Huber et al (2011)

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Purely empirical. Coefficients fitted to match the observations

 $A \propto \left(\frac{L}{M}\right)$

 $A \propto \eta^{-0.5} \; \frac{L}{M^{1.5} T^{2.25}}$

 $A \propto \frac{L^s}{M^t T^{r-1}_{\ c}}$

 $A \propto \left(rac{L}{M}
ight)^s$

So what do we know about mode amplitudes in red giants and related scaling relations?

In short: the overall picture is (at least) not clear for red-giant stars!

Samadi et al. (2012)

$$\left(\frac{\delta L}{L}\right)_{\rm max} \propto \eta_{\rm max}^{-0.5} \left(\frac{L}{M}\right)^{1.55} \Delta \nu$$

Based on theoretical modelling and a set of 3D hydrodynamical simulations

Samadi et al. (2012)

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Based on theoretical modelling and a set of 3D hydrodynamical simulations Belkacem et al. (2012)

$$\eta_{
m max}^{-1} \propto T_{
m eff}^{-10.8} \, g^{0.3}$$

Based on full non-adiabatic calculations

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m max}^{-0.65}$ + «classical» scalings Kepler observations scaling with coefficient exponent all 860 ± 30 -0.71 ± 0.01 $\nu_{\rm max}$ clump 4700 ± 500 1.10 ± 0.00 Abol (ppm) 100 RGB -0.63 ± 0.02 650 ± 60 agreement between theory and Mosser et al. (2012) observations, at least for RGB stars 10 10 100 $\nu_{\rm max}$ (µHz)

Belkacem et al. (2012)

It has been shown by Samadi et al (2007), for main-sequence stars

$$V_{
m max} \propto \left(rac{L}{M}
ight)^{0.}$$

confirmed by samadi et al. (2012) for red-giants

Now using an adiabatic conversion between mode velocity and luminosity

$$\left(\frac{\delta L}{L}\right)_{\rm max} \propto \left(\frac{L}{M}\right)^{0.7} \times \frac{1}{\sqrt{T_{\rm eff}}}$$

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Adiabatic computation and scaling relation under-estimate mode amplitudes for red giants

Most of the disagreement can be explained by non-adiabatic effects

 $i\sigma T\delta S = -\frac{\partial \delta L}{\partial m} \propto \frac{L}{M}$ • For main-sequence stars L/M ~ 1-5 • For Red-Giants stars L/M ~ 30-100

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 It impacts the shape of the eigenfunction and therefore the driving

 It impacts the conversion between velocity and luminosity perturbations

Conversion between surface velocities and luminosities

Space-borne observations

Theoretical computations







Conversion between surface velocities and luminosities

Space-borne observations

Theoretical computations





 V_{\max}

• adiabatic relation (e.g., Kjeldsen & Bedding 1995) $\left(rac{\delta L}{L}
ight)_{
m max}/V_{
m max} \propto \sqrt{rac{1}{T_{
m eff}}}$

commonly used with a solar calibration

✓ physically not justified✓ it does not work even for the Sun

Conversion between surface velocities and luminosities



10

L/M

100



Improves the agreement with the observations

non-adiabatic effect on mode compressibility



Hence, mode amplitude of red giants is potentially a way to get constraints on non-adiabaticity effects mode displacement rapidly varies at the photosphere due to non-adiabatic effect

> mode compressibility significantly increases, thus mode driving (~ factor 3)

$$\mathcal{P} \propto rac{1}{I} \int dm \left| rac{d\xi_{
m r}}{dr}
ight|^2 E_{
m eddy} \Lambda u_0$$



Concluding remarks

 The wealth of high-quality observational data from CoRoT and Kepler opens up great opportunities for stellar physics, but...

 For period spacings (as well as for large separation!): one has to be careful before any comparison between observations and models

 For mode amplitudes, the overall pictures is being more clear, but...



we crucially need constraints on non-adiabatic processes

ground-based observations, led by E. Poretti, are currently undertaken to this end, for several stars within the CoRoT redgiant working group