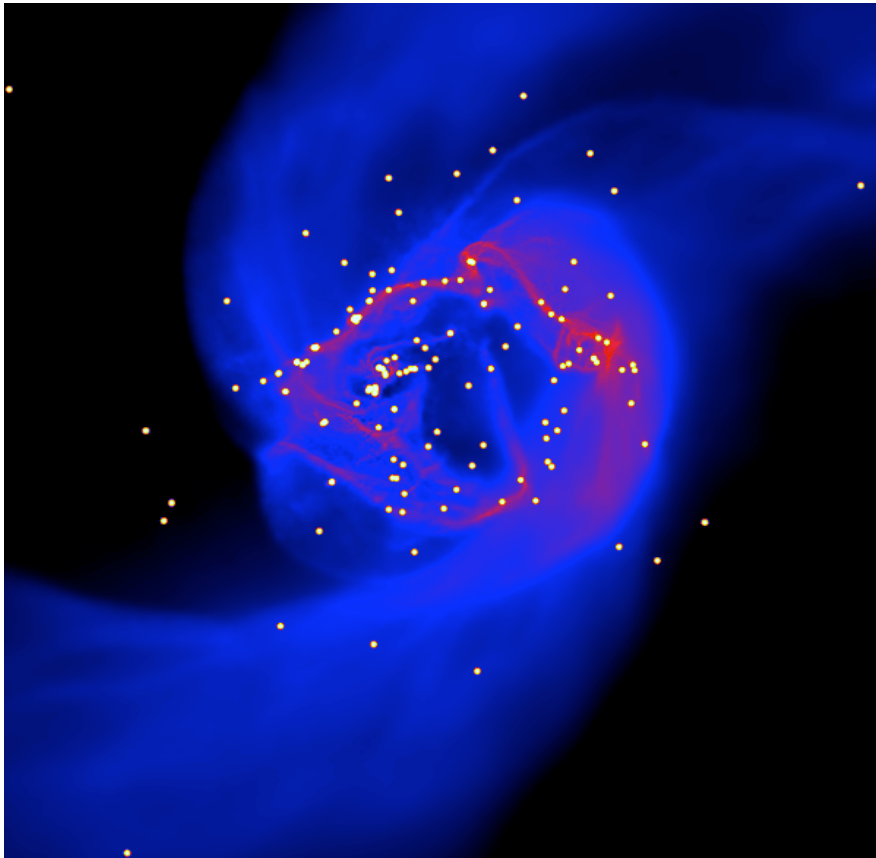


The IMF from primordial to present day star formation



Simon C.O. Glover

Paul C. Clark

Ralf S. Klessen

Rowan Smith

Gustavo Dopcke

Ian A. Bonnell

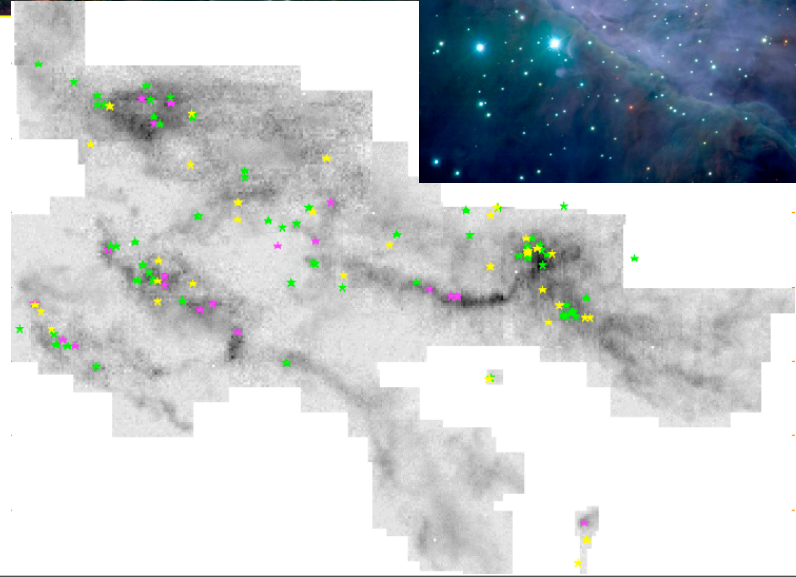
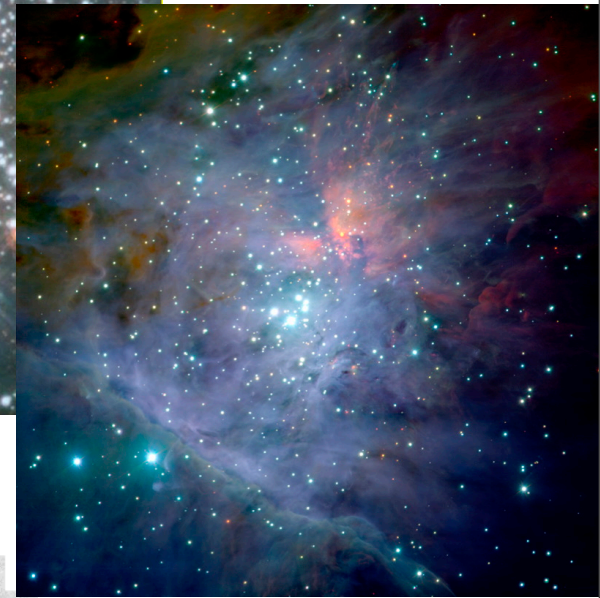
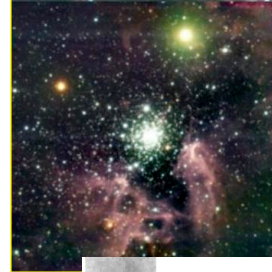
Volker Bromm

Thomas Greif

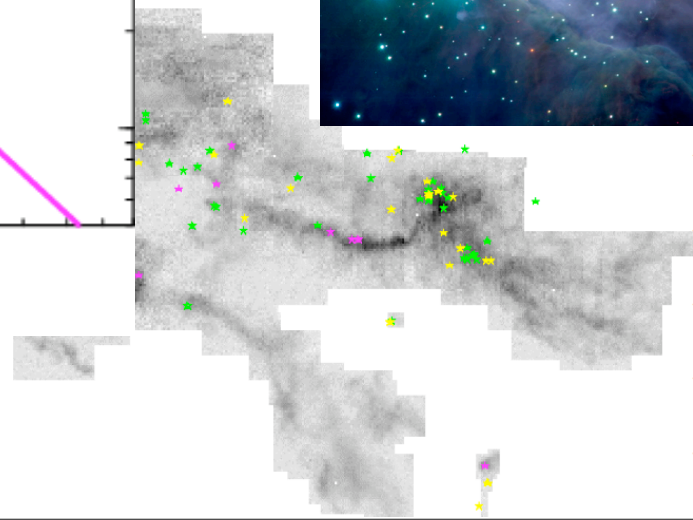
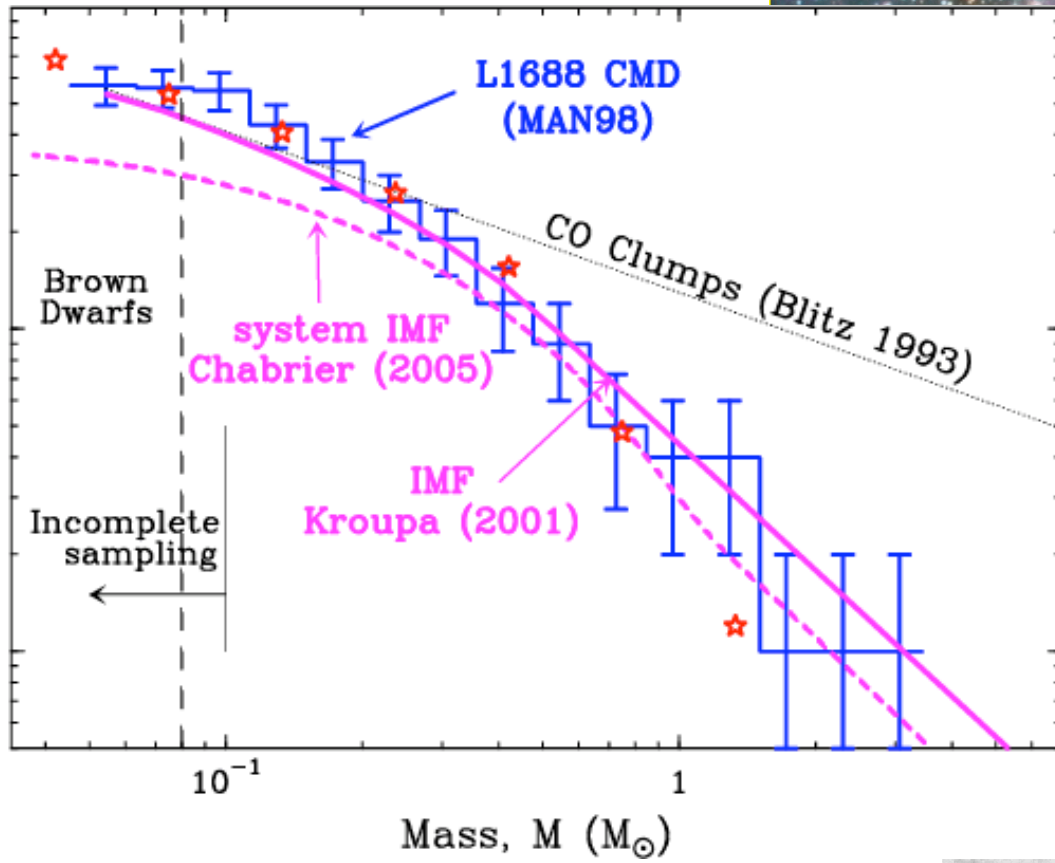
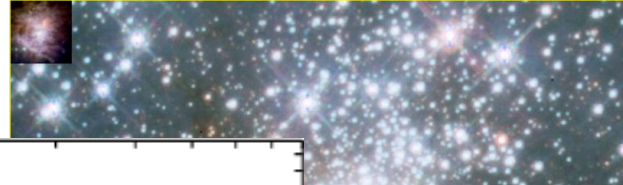
Present-day star formation

3 main features:

- Stars form in clusters/groups/association, not in isolation (e.g. Lada & Lada).
- Stellar mass function seems to be independent of environment and cluster size (Kroupa 2002).
- Most stars have **low-masses**, between 0.1 and $0.5 M_{\odot}$.

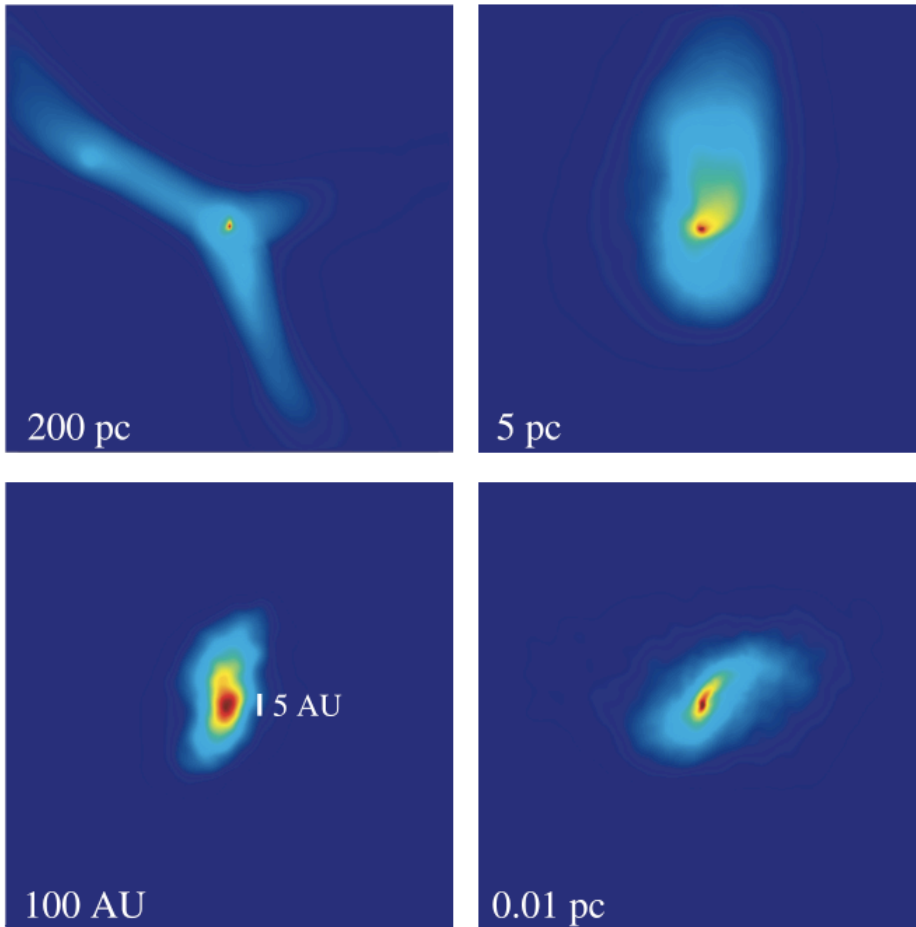


Present-day star formation



Primordial star formation

(Standard picture...)



Stars form in isolation

Stars are typically massive
($> 20M_{\odot}$)

(e.g. Abel et al 2002; Bromm et al 2002; Tan & McKee 2004; Yoshida et al 2006, 2008)

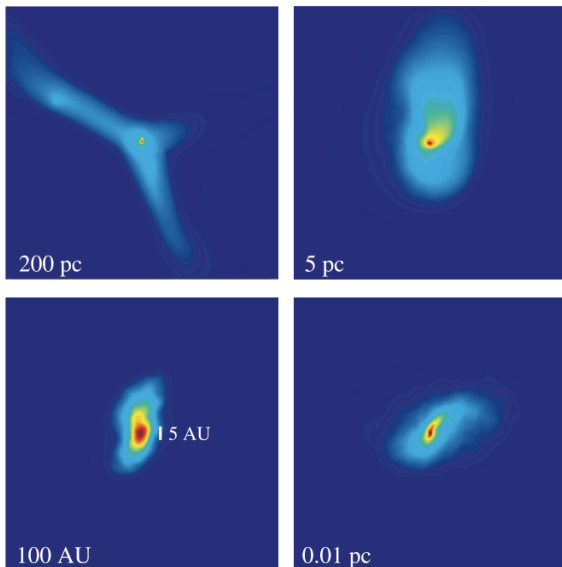
What causes the transition?

Jeans mass and length set the scales:

$$m_J = 1M_{\odot} \left[\frac{\rho}{10^{-19} \text{gcm}^{-3}} \right]^{-1/2} \left[\frac{T}{10\text{K}} \right]^{3/2}$$

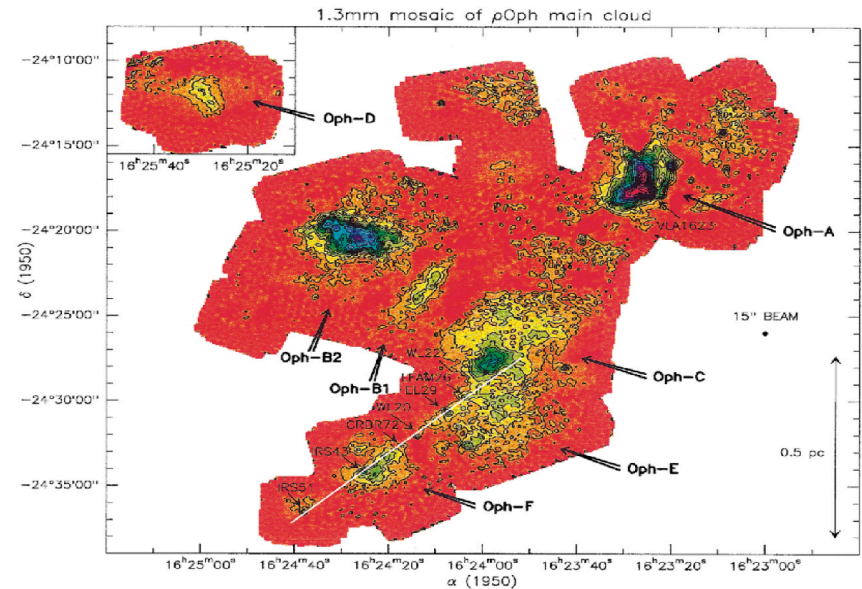
$$\lambda_J = 8500\text{au} \left[\frac{\rho}{10^{-19} \text{gcm}^{-3}} \right]^{-1/2} \left[\frac{T}{10\text{K}} \right]^{1/2}$$

$T \sim 300 \text{ K}$



~ 1 Jeans mass

$T \sim 15 \text{ K}$



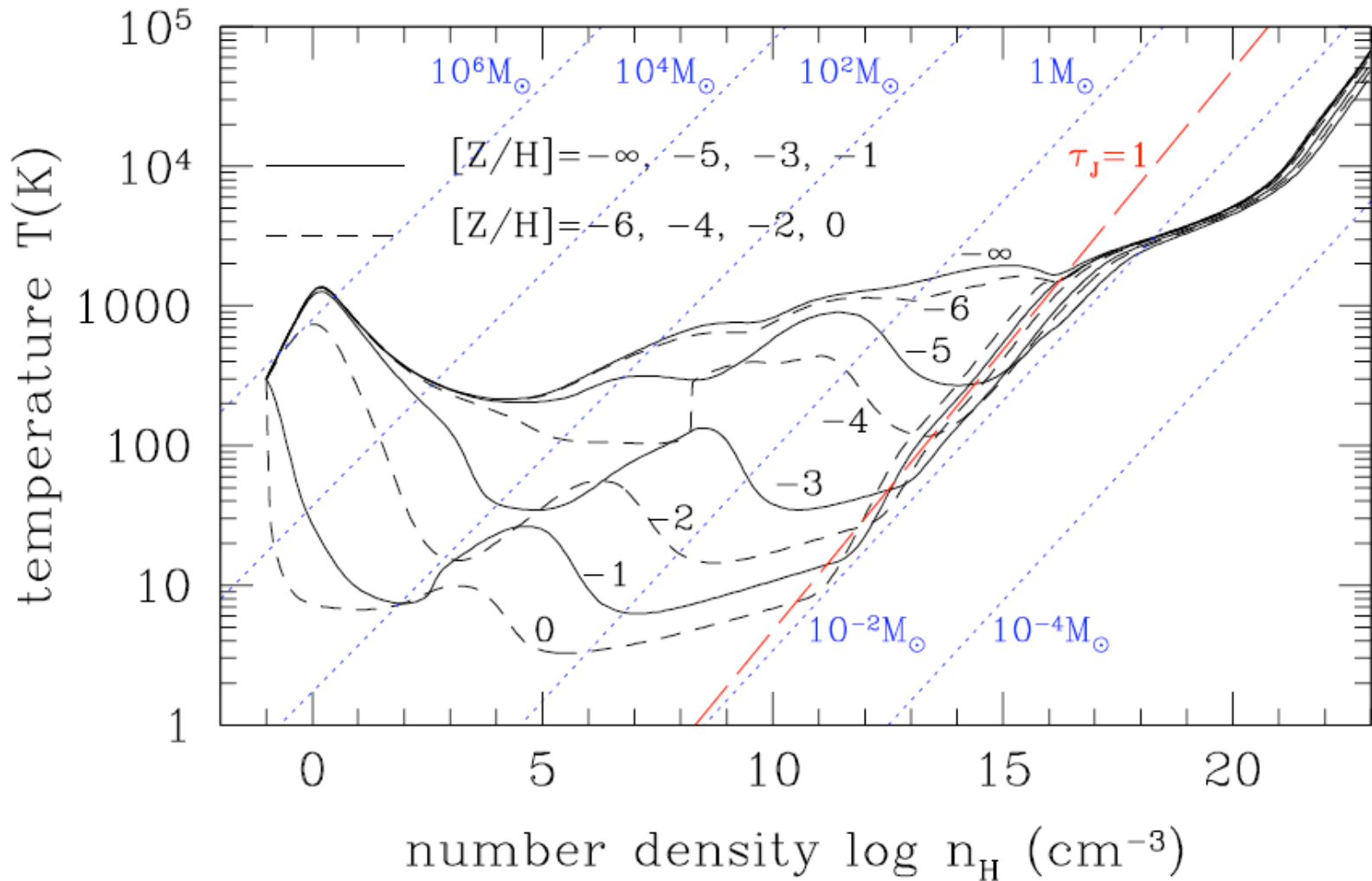
$> 10\text{s}$ of Jeans masses

Cooling causes transition

Two schools of thought:

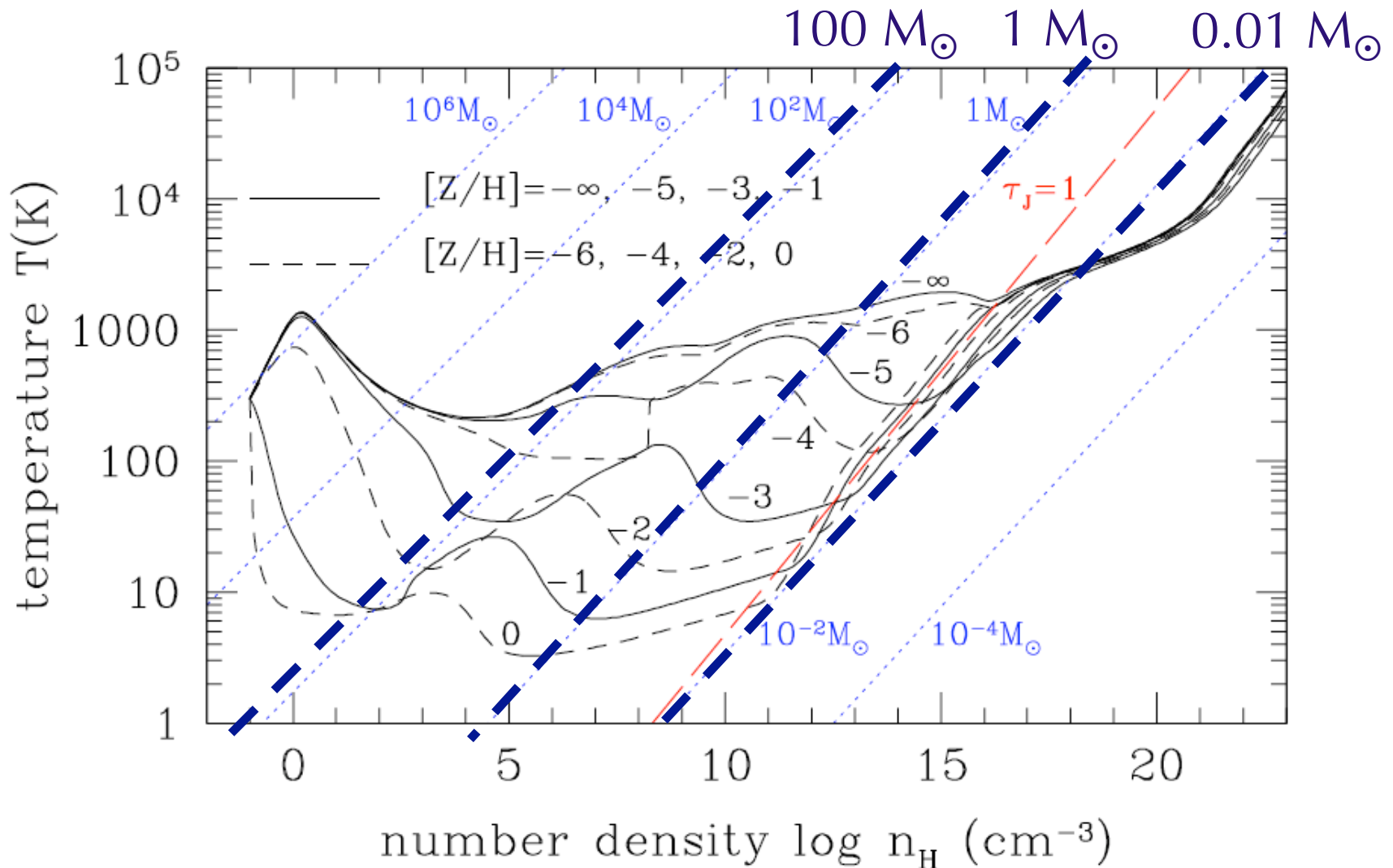
- 1)** C, O fine structure cooling? (e.g. Bromm et al 2001; Bromm & Loeb 2002; Santoro & Shull 2006; Frebel et al 2007; Smith & Sigurdsson 2007)
 - Occurs at low densities ($< 10^4 \text{ cm}^{-3}$)**
 - Sets a critical metallicity of $10^{-3.5} Z_{\odot}$**
- 2)** Dust-cooling induced fragmentation? (e.g. Schneider et al 2002; Omukai et al 2005; Schneider et al 2006)
 - Occurs at very high densities ($> 10^8 \text{ cm}^{-3}$)**
 - Possibly kicks in around 10^{-6} to $10^{-5} Z_{\odot}$**

Dust and line cooling



Omukai, Tsuribe, Schneider & Ferrara (2005)

Dust and line cooling



Omukai, Tsuribe, Schneider & Ferrara (2005)

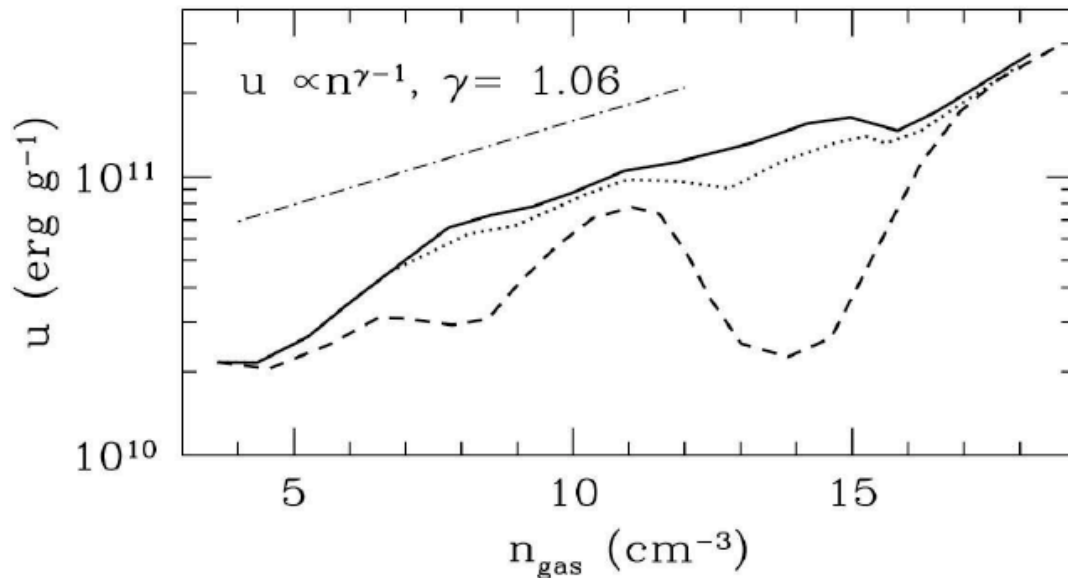
Initial conditions (I)

Gas clump parameters:

500 M_{\odot} with 25 or 2 Million SPH particles.

Mass resolution $\sim 0.002M_{\odot}$ or $\sim 0.02M_{\odot}$.

Piece-wise polytropic fit to Omukai et al (2005) EOS.



Initial conditions (II)

Gas clump parameters:

Initial density = $5 \times 10^5 \text{ cm}^{-3}$

Initial T = 200 ~ 250K (set by EOS)

$E_{\text{therm}} / |E_{\text{grav}}| = 0.39 - 0.32$ (depends on which EOS)

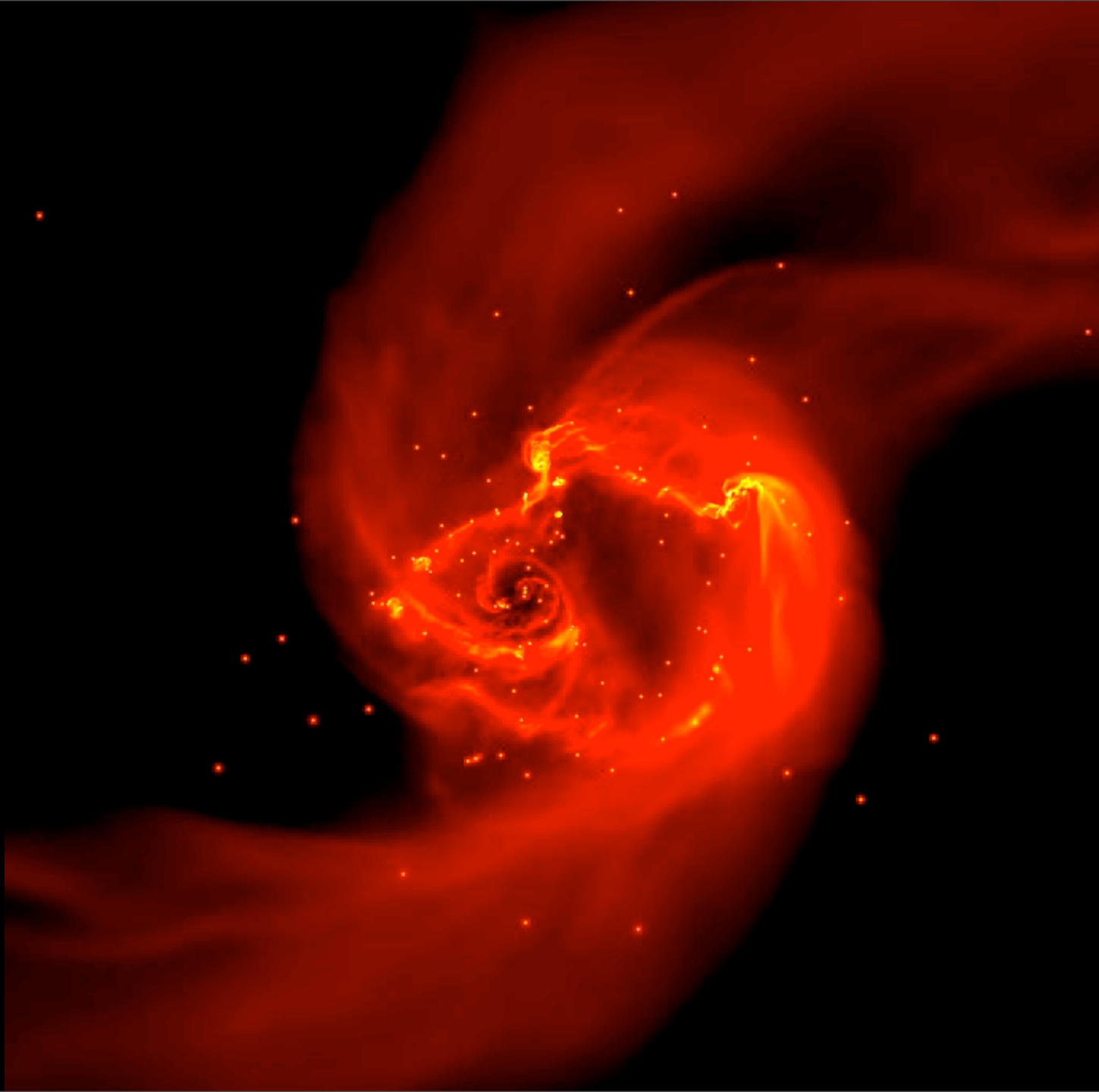
$E_{\text{turb}} / |E_{\text{grav}}| = 0.1$ (subsonic v_{rms})

$E_{\text{rot}} / |E_{\text{grav}}| = 0.02$ (added **on top** of turbulence)

Sinks Particles:

Form at $\sim 10^{17} \text{ cm}^{-3}$

Accretion radii of 0.4 AU

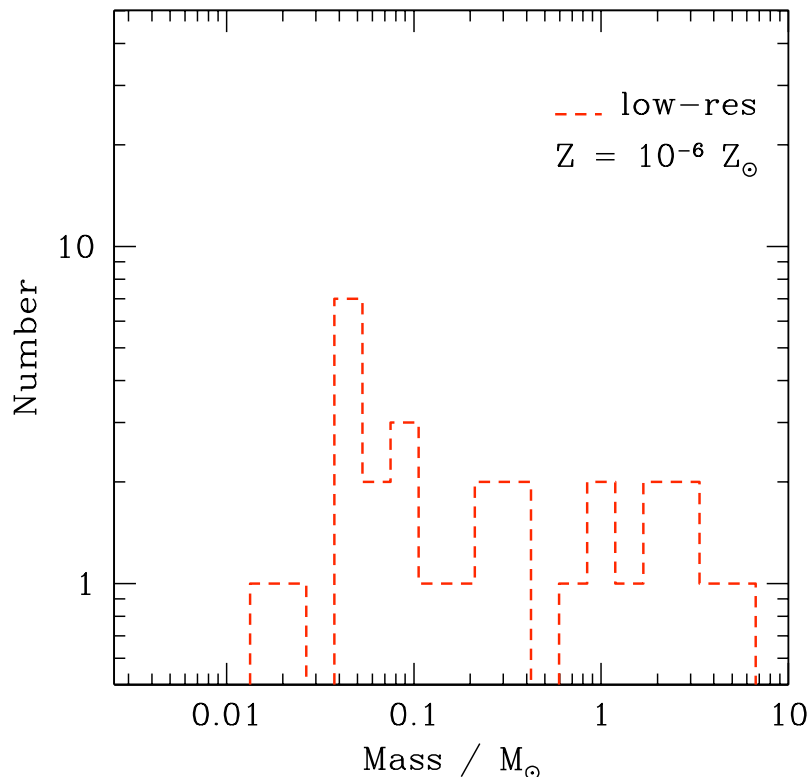


IMF?

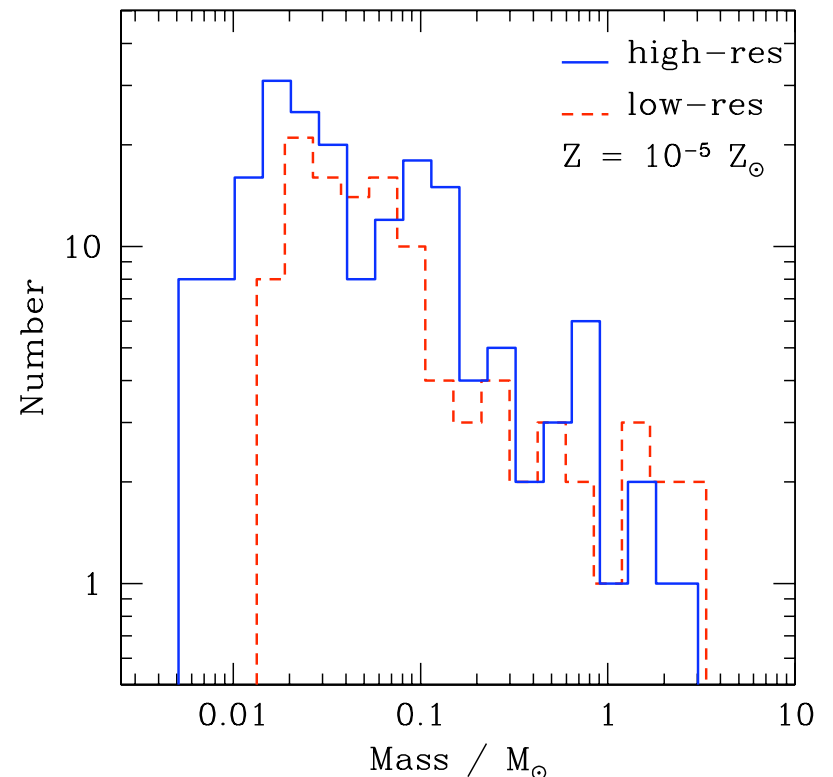
IMFs are plotted when $19 M_{\odot}$ is accreted:

Transition to present-day IMF?

Clark, Glover & Klessen (2008)



$Z = 10^{-6} Z_{\odot}$



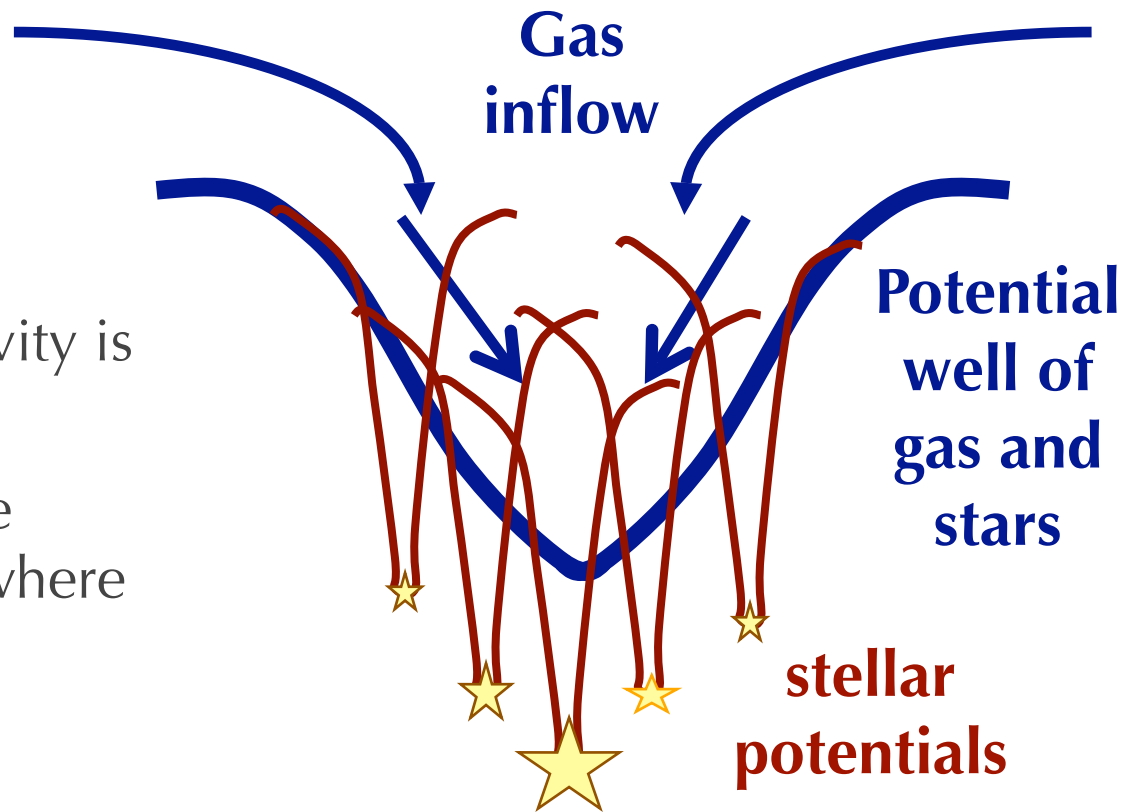
$Z = 10^{-5} Z_{\odot}$

Competitive accretion?

Bonnell & Bate (2006):

Need a situation where gravity is dominating the dynamics.

A collapsing, Jeans unstable region, creates a situation where competitive accretion is **unavoidable**.

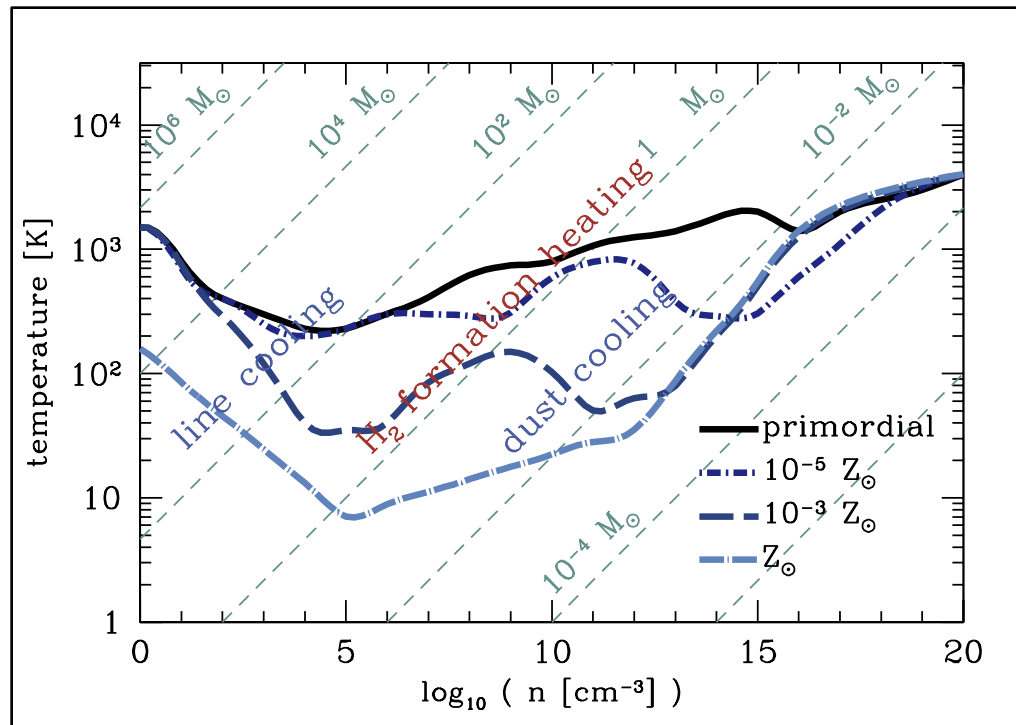


A universal slope for the IMF?

Clark, Glover, Bonnell & Klessen (2009):

Any regime in which $t_{\text{cool}} < t_{\text{ff}}$ during collapse gives you these conditions.

Could be line and/or dust cooling.



Caveats...

We're using an equation of state, so no self-consistent heating-cooling and dynamics.

Assumptions about the dust: are scaled down solar properties valid?

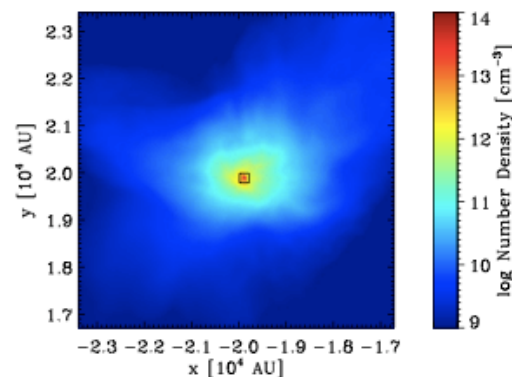
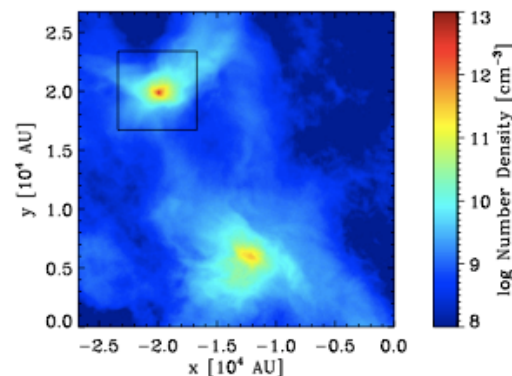
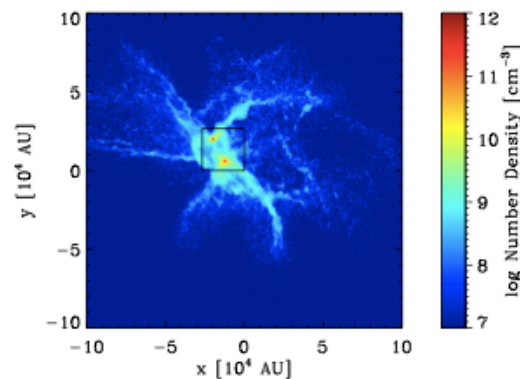
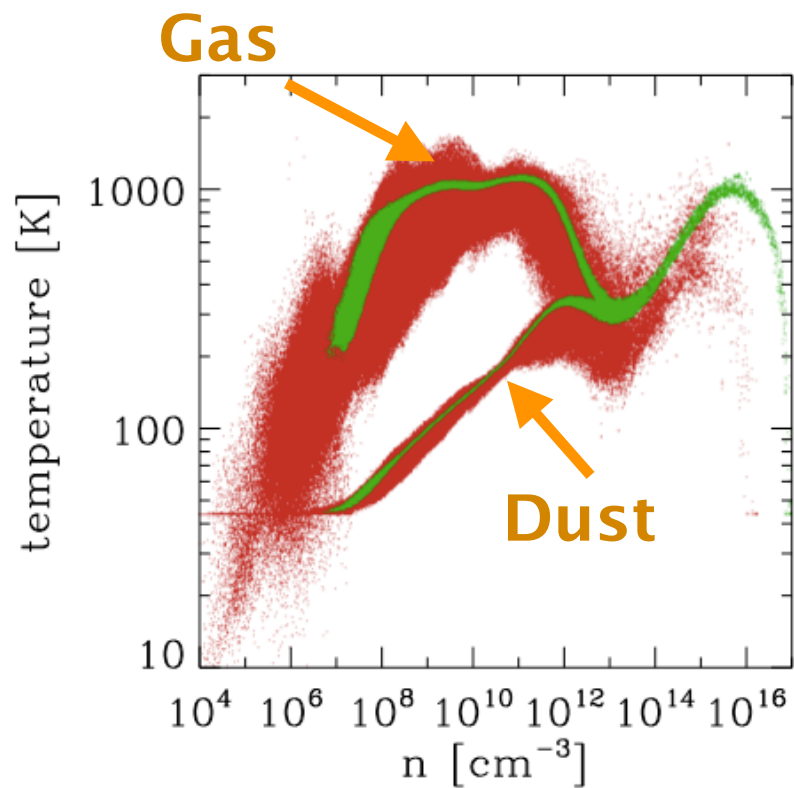
Uncertainties in the chemical reaction rates will affect the temperature - density relationship.

No feed-back from the stars: assuming that dust can cope with the accretion luminosity.

Assumes that chemical enrichment allows $10^{-5} Z_{\odot}$ gas: may always overshoot to higher values.

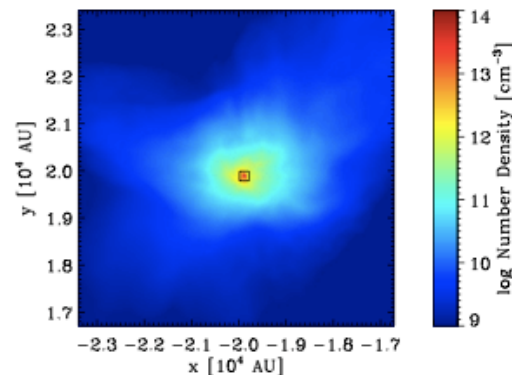
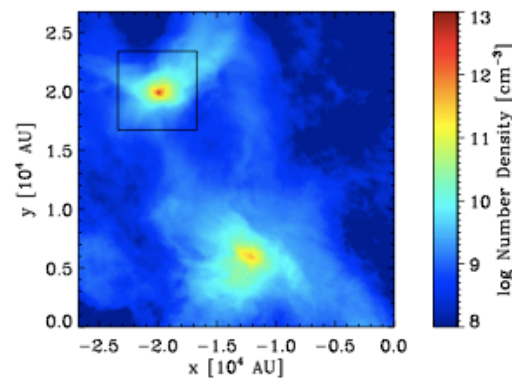
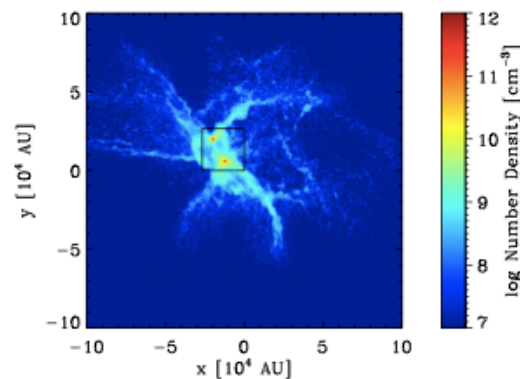
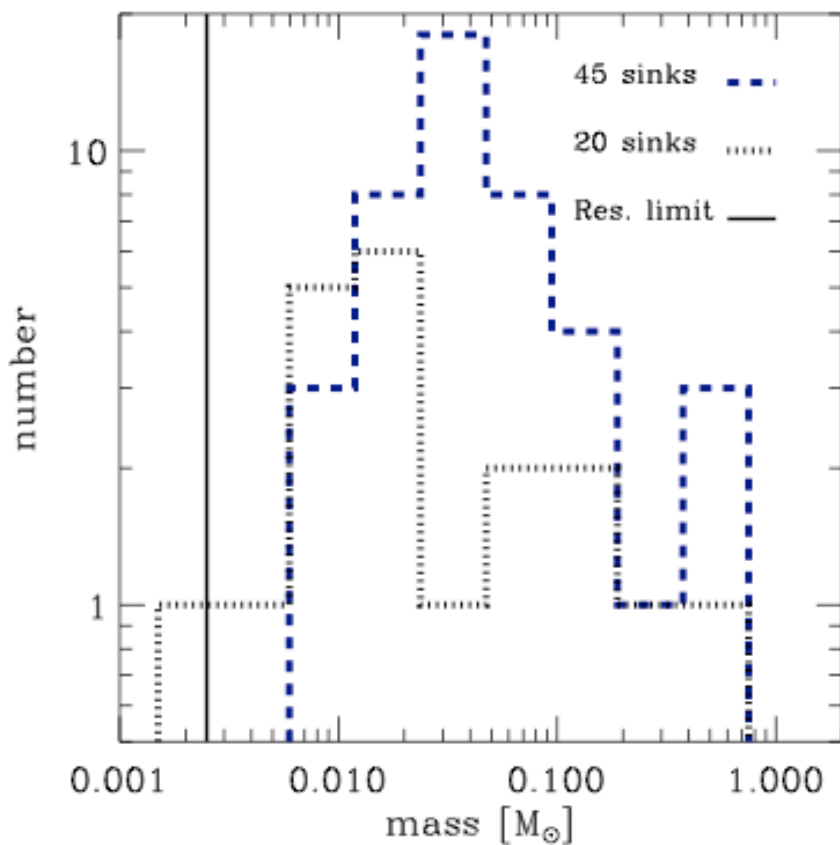
Caveat i) Self-consistent thermodynamics

Gustavo Dopcke's PhD:



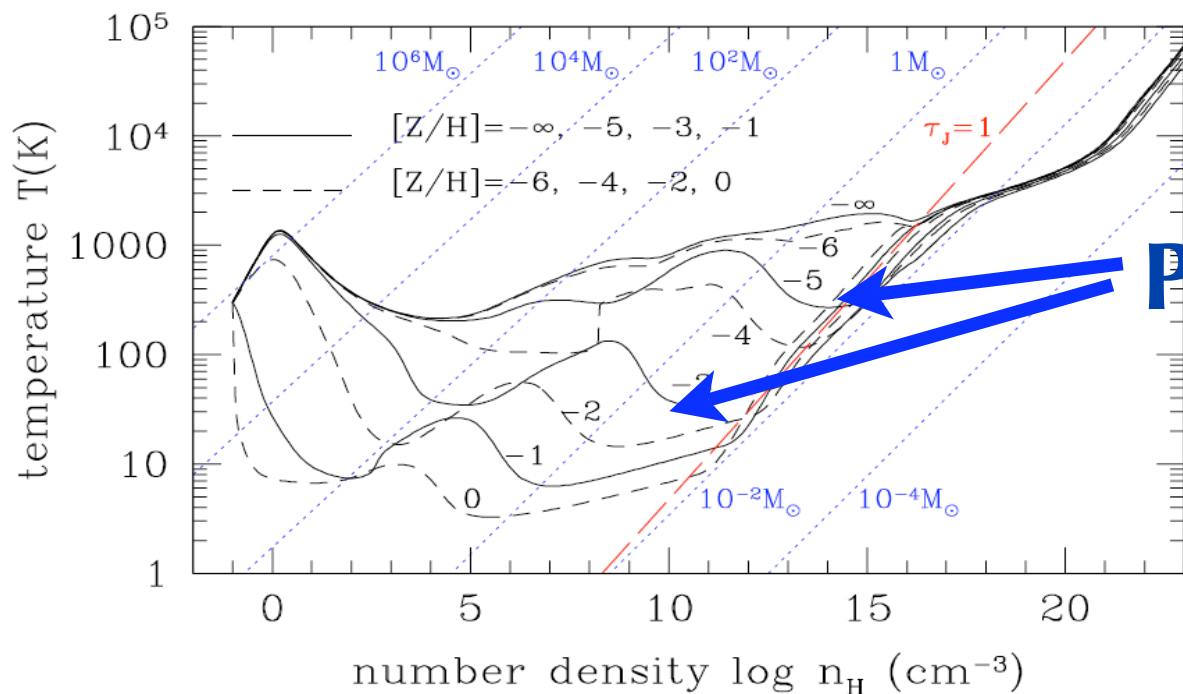
Caveat i) Self-consistent thermodynamics

Gustavo Dopcke's PhD:



Conclusions from these models

- The peak in the mass function is closely related to the Jeans mass at the trough of the cooling curve.
- The upper mass function is roughly Salpeter-like.



Peak + Salpeter

So what about the Pop III IMF?

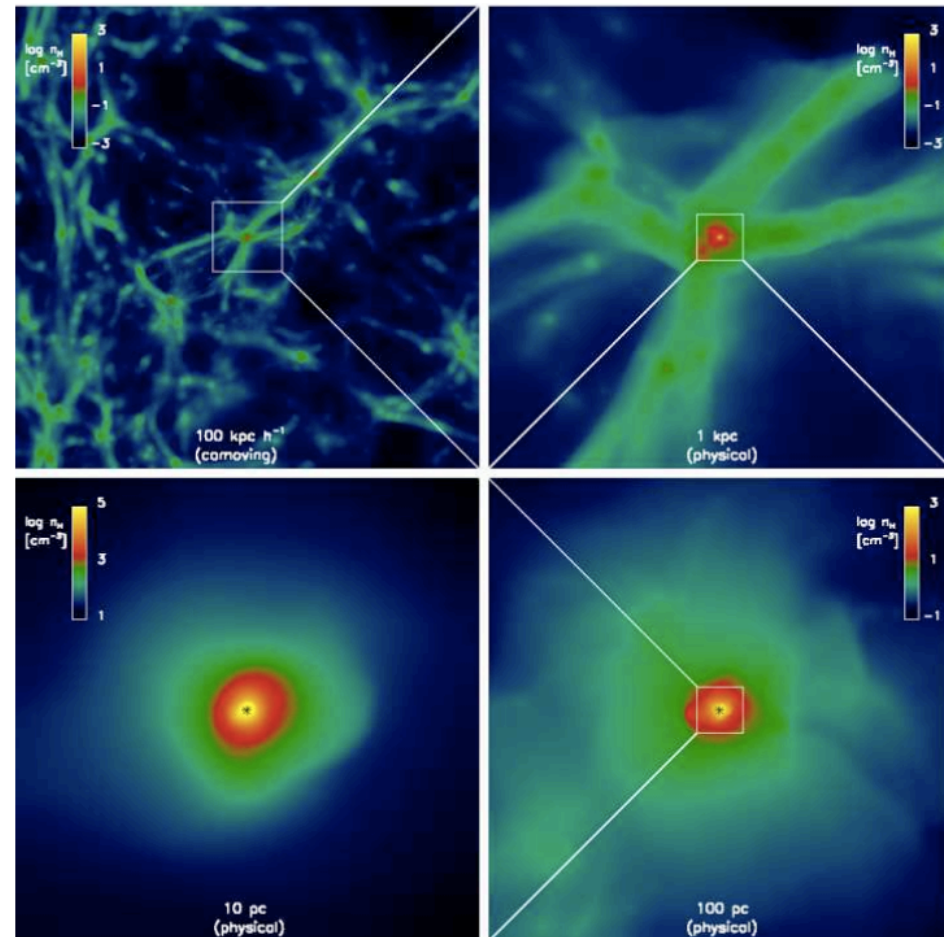
Cosmological initial conditions

Initial conditions from a cosmological GADGET2 simulation (performed by T. H. Greif, similar to that in Stacy et al 2010):

- 200 kpc (co-moving) cosmological box
- Λ -CDM: $\Omega_m = 1 - \Omega_\Lambda = 0.3$; $\Omega_b = 0.04$;
 $h = H_0/100 \text{ km s}^{-1} \text{ Mpc}^{-1} = 0.7$; $\sigma_8 = 0.9$
- Evolved from $z = 99$ to 22, when baryons first become self-gravitating

Then re-zoom the simulation:

- Go from Thomas' $5 M_{\text{sun}}$ SPH resolution to $0.05 M_{\text{sun}}$, and run collapse to $n = 10^{13} \text{ cm}^{-3}$.
- Refine final stages of collapse to $0.001 M_{\text{sun}}$ resolution --> 200,000 SPH particles in disc



Stacy, Greif & Bromm (2010)

Additional components

Chemical processes:

- 3-body H₂ formation heating.
- Rotational + vibrational line-emission from H₂ (Glover & Abel 2008).
- At high densities, H₂ energy levels are computed accounting for the escape-probability for the photons (Yoshida et al 2006).
- Collision induced emission (CIE; Ripamonti & Abel 2004) + reduction by continuum absorption (see Matt Turk).

Luminosity feedback:

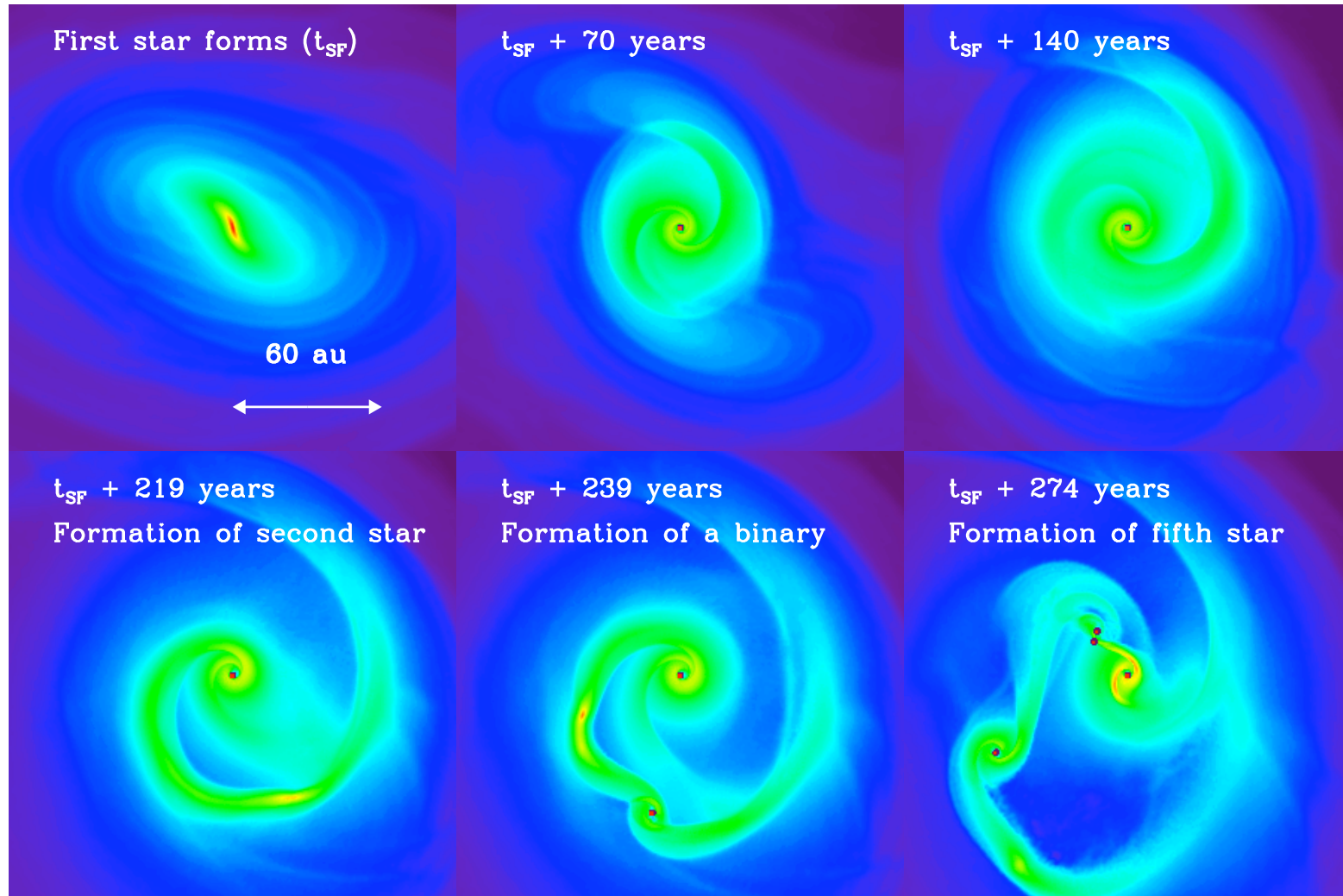
$$R_* = 26 M_*^{0.27} (\dot{M}/10^{-3})^{0.41}$$

$$L_{acc} = GM_* \dot{M} / R_*$$

$$\Gamma_{acc} = \rho_g K_P \left(\frac{L_{acc}}{4\pi r^2} \right)$$

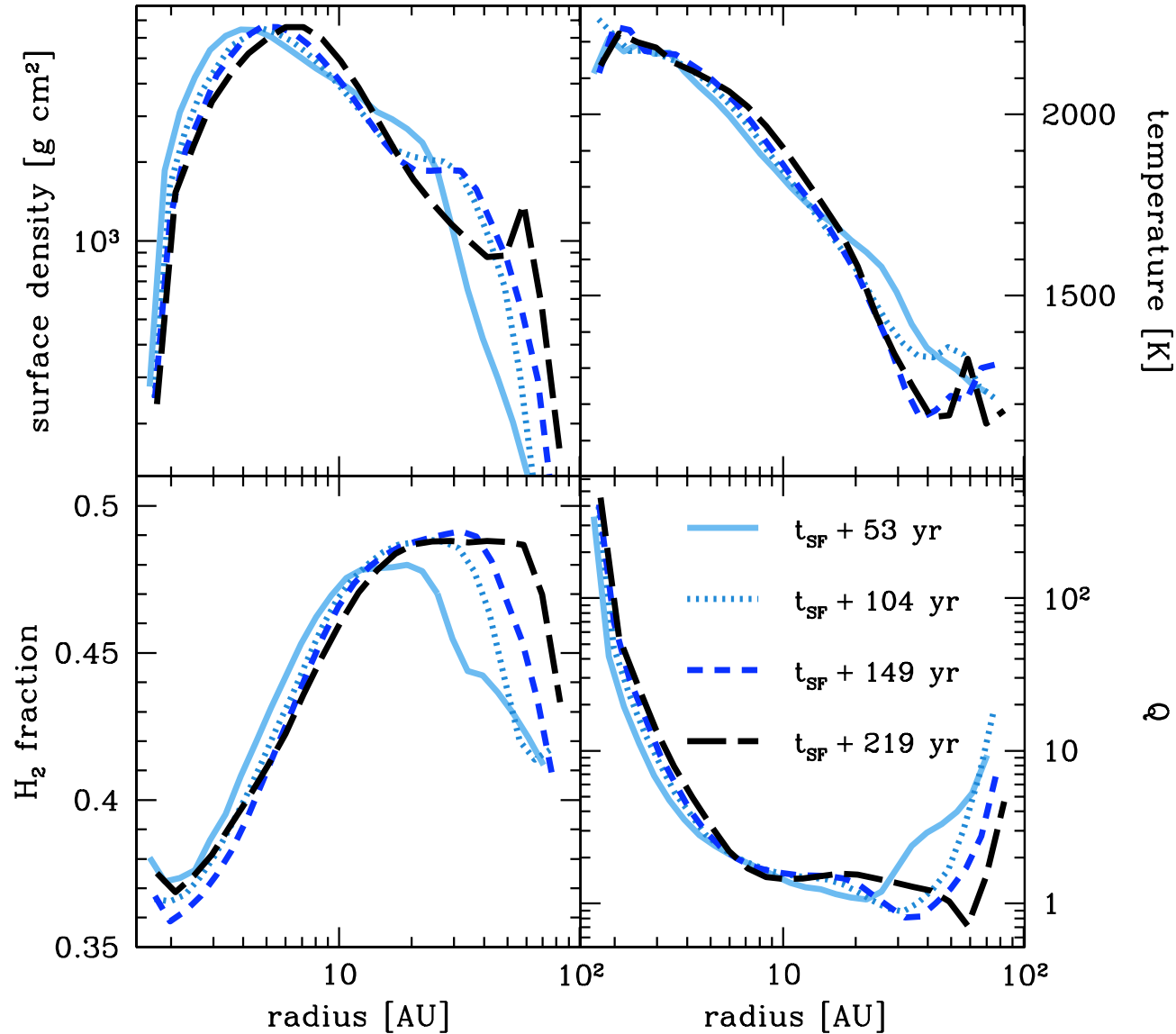
- Mass-radius relationship from Stahler, Palla & Salpeter (1986).
- Plank mean opacities from Mayer & Duschl (2005).
- We fix L_{acc} at $10^{-2} M_{sun} \text{ yr}^{-1}$ in our current simulations.

Fragmentation around Pop III protostellar discs?

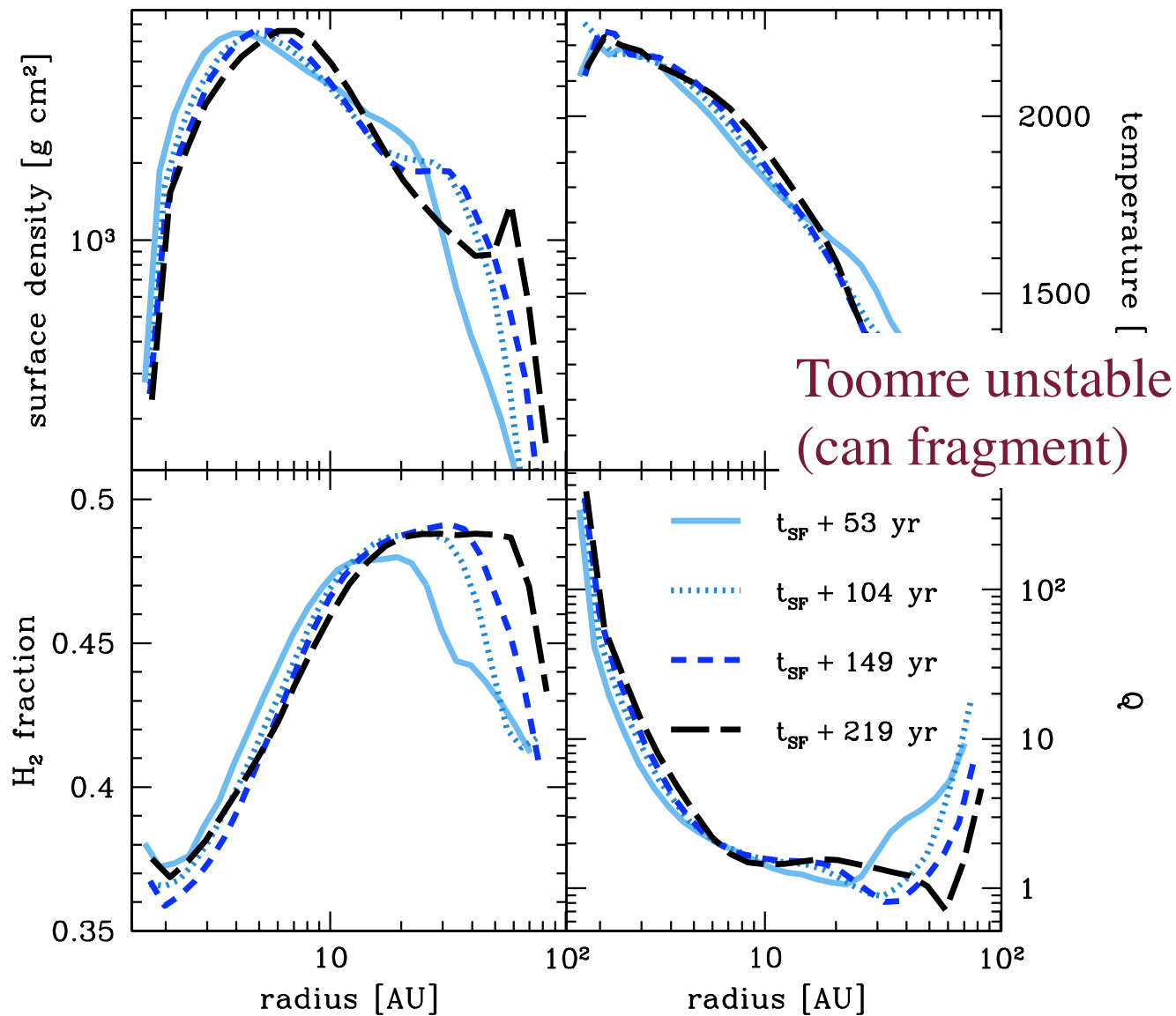


Clark et al (2010, *submitted*)

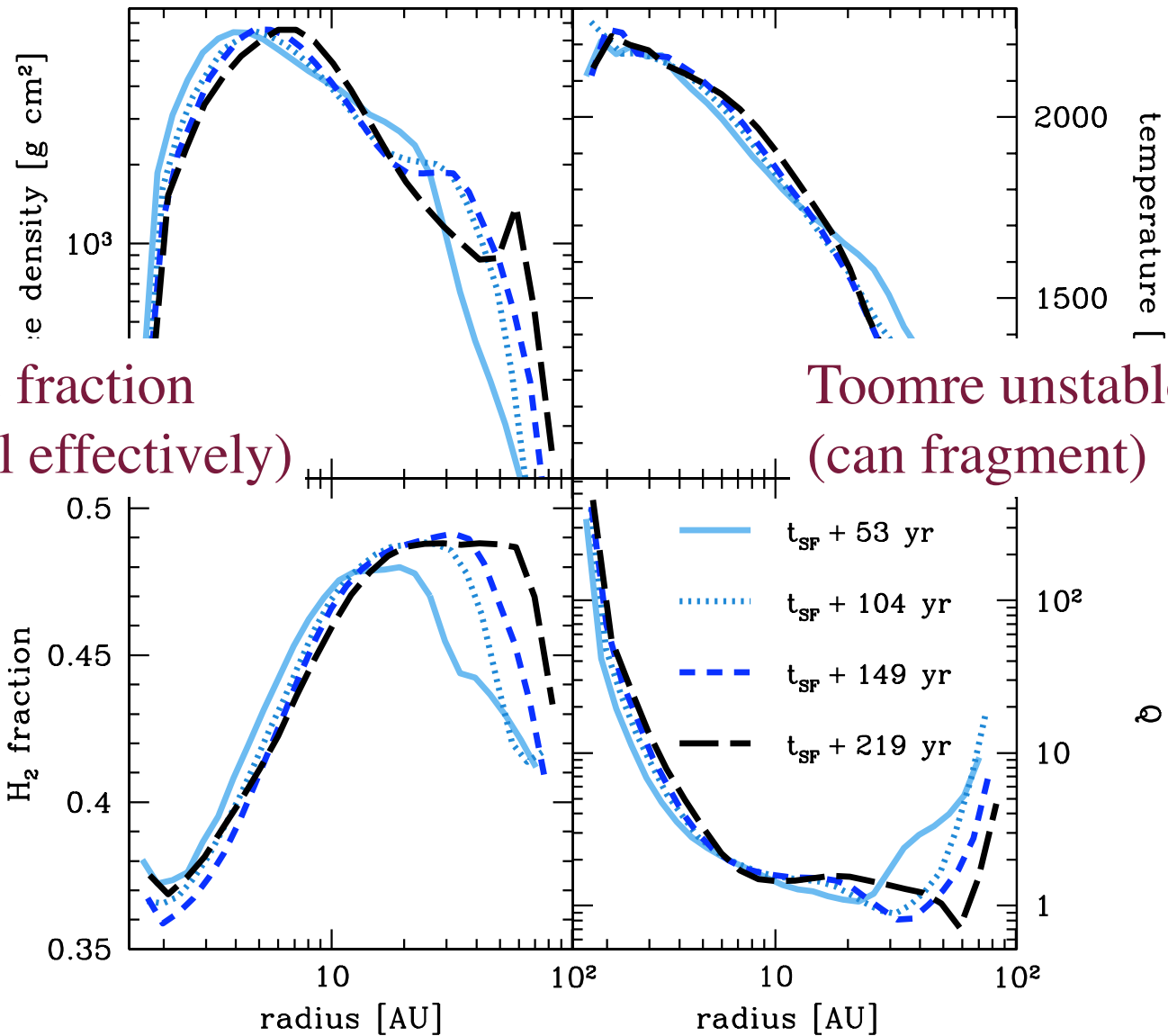
So why does the disc form stars?



So why does the disc form stars?



So why does the disc form stars?



High H_2 fraction
(can cool effectively)

Toomre unstable
(can fragment)

IMF of Pop III stars?

Too early to tell at these scales.

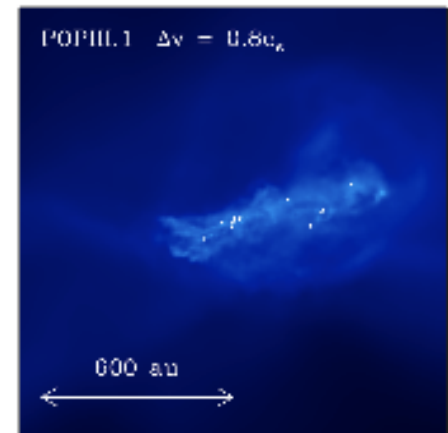
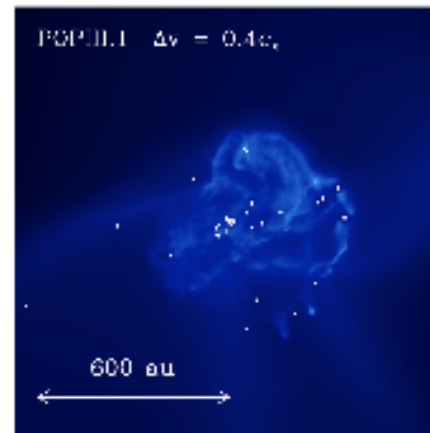
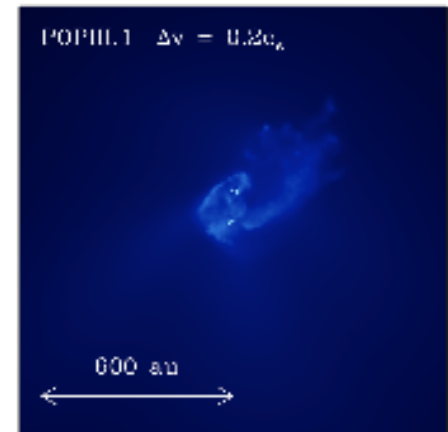
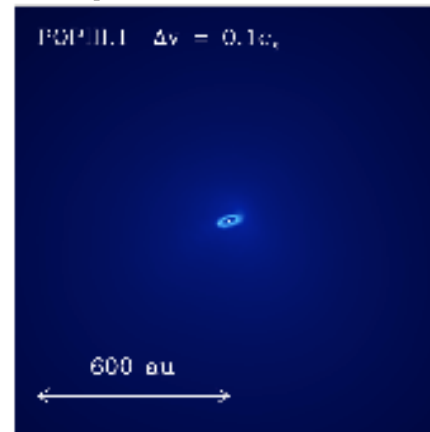
**But binary (and small N)
systems seem unavoidable**

What about larger scales?

BE spheres injected with subsonic turbulence:

- Primordial gas Pop III.1 and Pop III.2 channels.
- Find significant fragmentation in the Pop III.1 but **VERY** little in the Pop III.2 case.
- If baryonic component of mini-halos has significant turbulent component (e.g. Turk, Abel & O'Shea 2009), fragmentation may be common.

Pop III.1

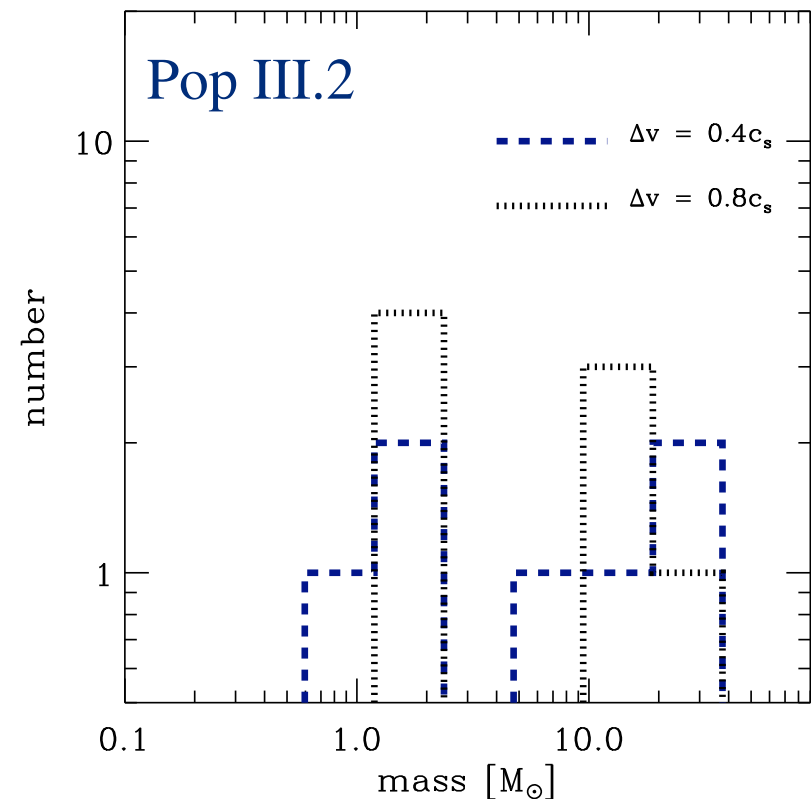
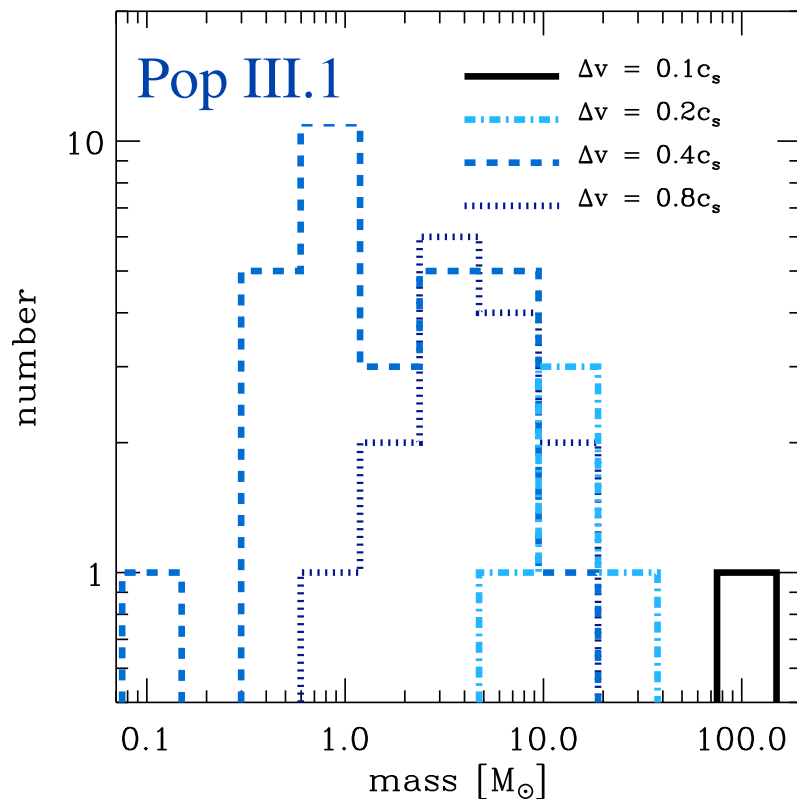


Clark, Glover, Klessen & Bromm (2010)

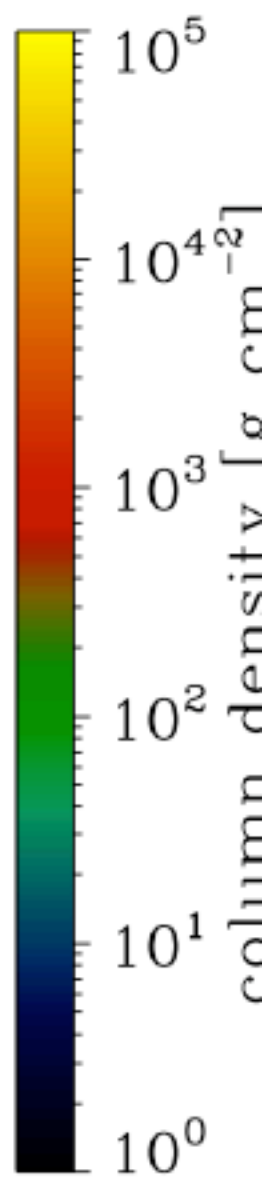
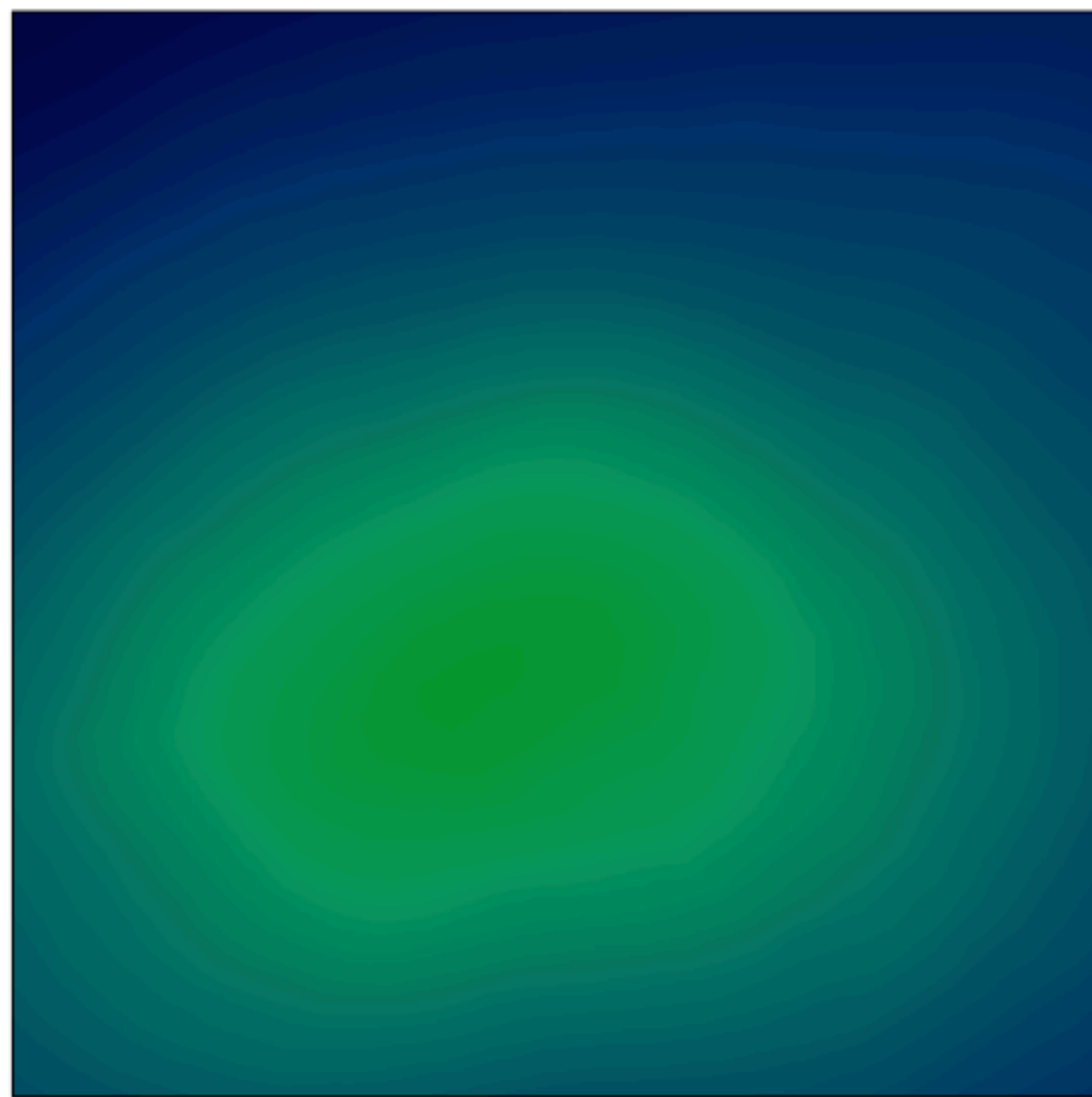
Pop III.1/III.2 IMF?

So if our little experiment is related to reality (!):

Clark, Glover, Klessen & Bromm (2010)



Clearly full cosmological simulations are necessary to test how likely these initial conditions are.



Summary

- Dust cooling can provide an efficient transition to the type of low-mass star formation that we see today.
- The upper mass function is roughly Salpeter-like due to competitive accretion.
- The turnover in the IMF (or characteristic mass) appears to be connected to the T-rho space at which dust cooling is most effective.
- The Pop III IMF is more uncertain, however fragmentation is inevitable, and a broad IMF seems likely.
- If Pop III stars do form in clusters, then Salpeter slope is possible with a turnover $> 1M_{\odot}$