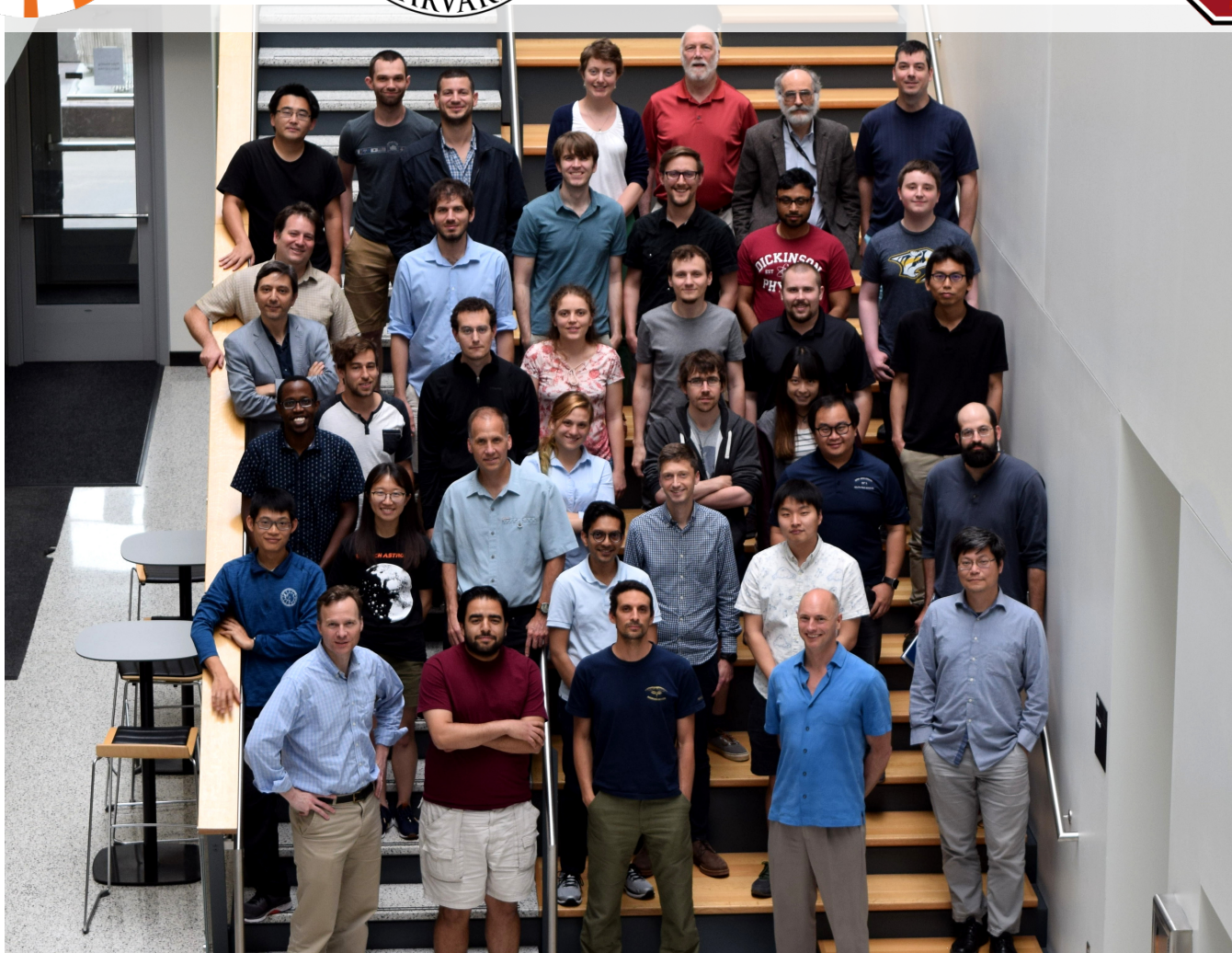


# Searching for Primordial Gravitational Waves with the BICEP/Keck Telescopes

Clem Pryke for the BICEP/Keck Collaboration – Tenerife – Oct 15 2018







# BK15 Results On arxiv Today

BICEP2 / *Keck Array X*: Constraints on Primordial Gravitational Waves using *Planck*, WMAP, and New BICEP2/*Keck* Observations through the 2015 Season

*Keck Array* and BICEP2 Collaborations: P. A. R. Ade,<sup>1</sup> Z. Ahmed,<sup>2</sup> R. W. Aikin,<sup>3</sup> K. D. Alexander,<sup>4</sup> D. Barkats,<sup>4</sup> S. J. Benton,<sup>5</sup> C. A. Bischoff,<sup>6</sup> J. J. Bock,<sup>3,7</sup> R. Bowens-Rubin,<sup>4</sup> J. A. Brevik,<sup>8</sup> I. Buder,<sup>4</sup> E. Bullock,<sup>8</sup> V. Buza,<sup>4,9</sup> J. Connors,<sup>4</sup> J. Cornelison,<sup>4</sup> B. P. Crill,<sup>7</sup> M. Crumrine,<sup>10</sup> M. Dierckx,<sup>4</sup> L. Duband,<sup>11</sup> C. Dvorkin,<sup>9</sup> J. P. Filippini,<sup>12,13</sup> S. Fliescher,<sup>10</sup> J. Grayson,<sup>14</sup> G. Hall,<sup>10</sup> M. Halpern,<sup>15</sup> S. Harrison,<sup>4</sup> S. R. Hildebrandt,<sup>3,7</sup> G. C. Hilton,<sup>16</sup> H. Hui,<sup>3</sup> K. D. Irwin,<sup>14,2,16</sup> J. Kang,<sup>14</sup> K. S. Karkare,<sup>4,17</sup> E. Karpel,<sup>14</sup> J. P. Kaufman,<sup>18</sup> B. G. Keating,<sup>18</sup> S. Kefeli,<sup>3</sup> S. A. Kernasovskiy,<sup>14</sup> J. M. Kovac,<sup>4,9</sup> C. L. Kuo,<sup>14,2</sup> N. A. Larsen,<sup>17</sup> K. Lau,<sup>10</sup> E. M. Leitch,<sup>17</sup> M. Lueker,<sup>3</sup> K. G. Megerian,<sup>7</sup> L. Moncelsi,<sup>3</sup> T. Namikawa,<sup>19</sup> C. B. Netterfield,<sup>20,21</sup> H. T. Nguyen,<sup>7</sup> R. O'Brien,<sup>3,7</sup> R. W. Ogburn IV,<sup>14,2</sup> S. Palladino,<sup>6</sup> C. Pryke,<sup>10,8</sup> B. Racine,<sup>4</sup> S. Richter,<sup>4</sup> A. Schillaci,<sup>3</sup> R. Schwarz,<sup>10</sup> C. D. Sheehy,<sup>22</sup> A. Soliman,<sup>3</sup> T. St. Germaine,<sup>4</sup> Z. K. Staniszewski,<sup>3,7</sup> B. Steinbach,<sup>3</sup> R. V. Sudhwala,<sup>1</sup> G. P. Teply,<sup>3,18</sup> K. L. Thompson,<sup>14,2</sup> J. E. Tolan,<sup>14</sup> C. Tucker,<sup>1</sup> A. D. Turner,<sup>7</sup> C. Umiltà,<sup>6</sup> A. G. Vieregg,<sup>23,17</sup> A. Wandui,<sup>3</sup> A. C. Weber,<sup>7</sup> D. V. Wiebe,<sup>15</sup> J. Willmert,<sup>10</sup> C. L. Wong,<sup>4,9</sup> W. L. K. Wu,<sup>17</sup> H. Yang,<sup>14</sup> K. W. Yoon,<sup>14,2</sup> and C. Zhang<sup>9</sup>

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<sup>12</sup>*Service des Basses Températures, Commissariat à l'Énergie Atomique, 38054 Grenoble, France*

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(Draft As accepted by PRL)

We present results from an analysis of all data taken by the BICEP2/*Keck* CMB polarization experiments up to and including the 2015 observing season. This includes the first *Keck Array* observations at 220 GHz and additional observations at 95 & 150 GHz. The *Q/U* maps reach depths of 5.2, 2.9 and 26  $\mu\text{K}_{\text{rms}}$  arcmin at 95, 150 and 220 GHz respectively over an effective area of  $\approx 400$  square degrees. The 220 GHz maps achieve a signal-to-noise on polarized dust emission approximately equal to that of *Planck* at 353 GHz. We take auto- and cross-spectra between these maps and publicly available WMAP and *Planck* maps at frequencies from 23 to 353 GHz. We evaluate the joint likelihood of the spectra versus a multicomponent model of lensed- $\Lambda\text{CDM}+r+\text{dust}+\text{synchrotron}+\text{noise}$ . The foreground model has seven parameters, and we impose priors on some of these using external information from *Planck* and WMAP derived from larger regions of sky. The model is shown to be an adequate description of the data at the current noise levels. The likelihood analysis yields the constraint  $r_{0.05} < 0.07$  at 95% confidence, which tightens to  $r_{0.05} < 0.06$  in conjunction with *Planck* temperature measurements and other data. The lensing signal is detected at 8.8 $\sigma$  significance. Running maximum likelihood search on simulations we obtain unbiased results and find that  $\sigma(r) = 0.020$ . These are the strongest constraints to date on primordial gravitational waves.

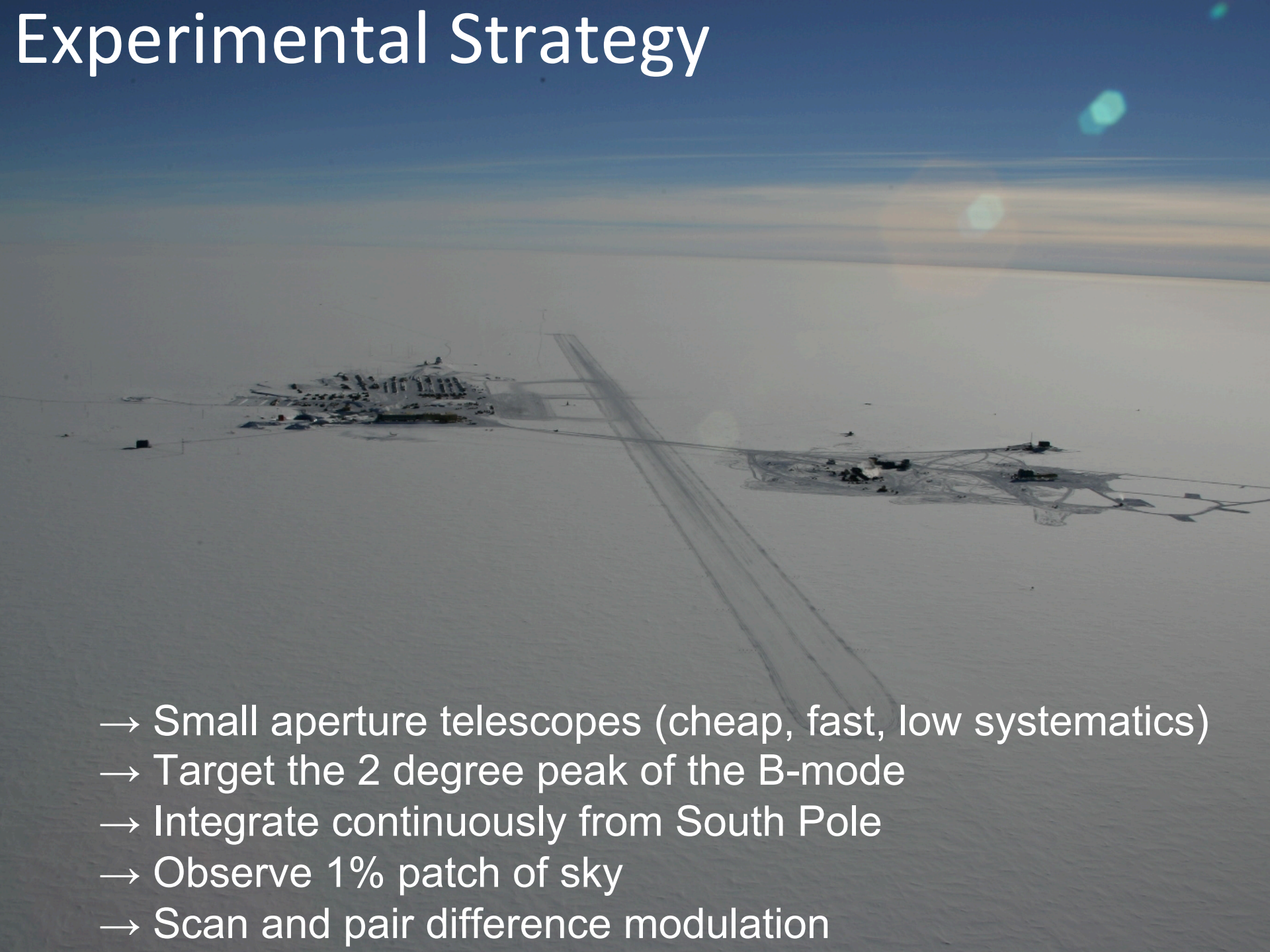
arxiv/1810.05216

BK15 = includes all data taken up to, and including, 2015 season

Three years since BK14 – Sorry for the delay!



# Experimental Strategy



- Small aperture telescopes (cheap, fast, low systematics)
- Target the 2 degree peak of the B-mode
- Integrate continuously from South Pole
- Observe 1% patch of sky
- Scan and pair difference modulation



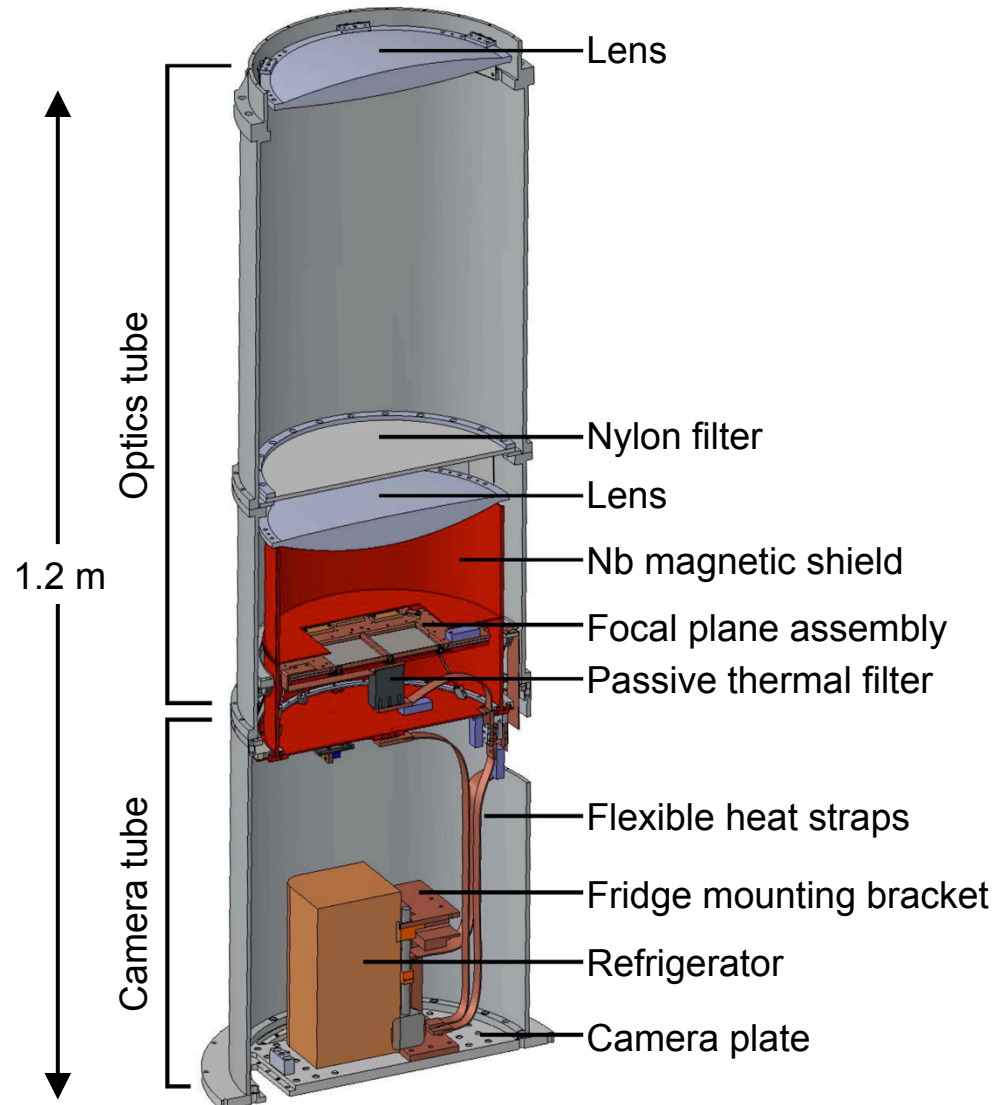
# The BICEP2/Keck Telescopes

Telescope as compact as possible while still having the angular resolution to observe degree-scale features

On-axis, refractive optics allow the entire telescope to rotate around boresight for polarization modulation

Liquid helium/pulse tube cools the optical elements to 4 K

3-stage helium sorption refrigerator further cools the detectors to 0.27 K





# Planar superconducting detector arrays

...designed to scale  
in frequency

Up to 2013 – all 150GHz

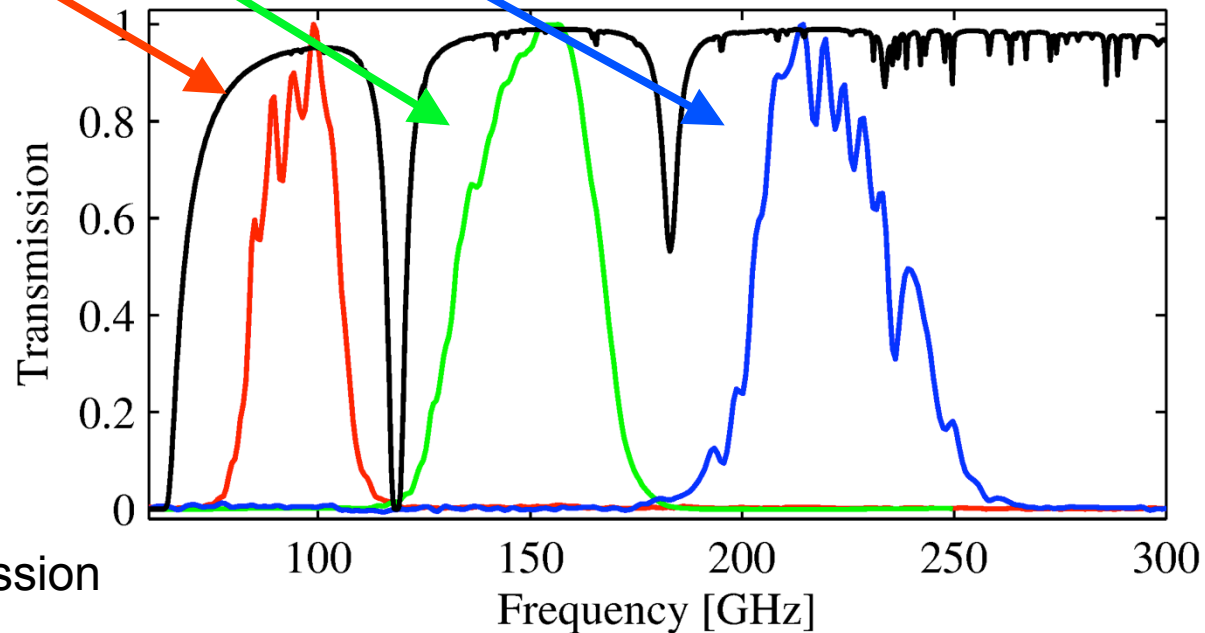
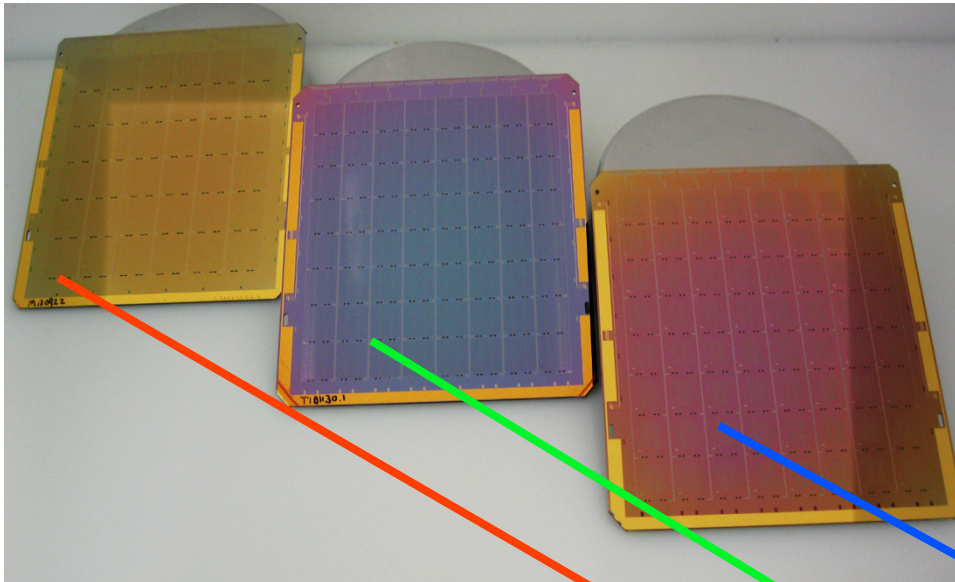
2014 – 2x95 3x150GHz

2015 – 2x95 1x150 2x220GHz

2016 – B3 1x150 4x220GHz

2017 – B3 4x220 1x270GHz

2018 – B3 4x220 1x270GHz

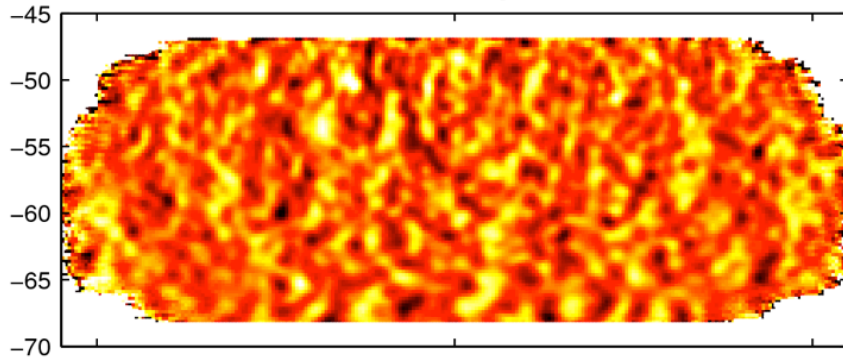


Typical South Pole  
atmospheric transmission

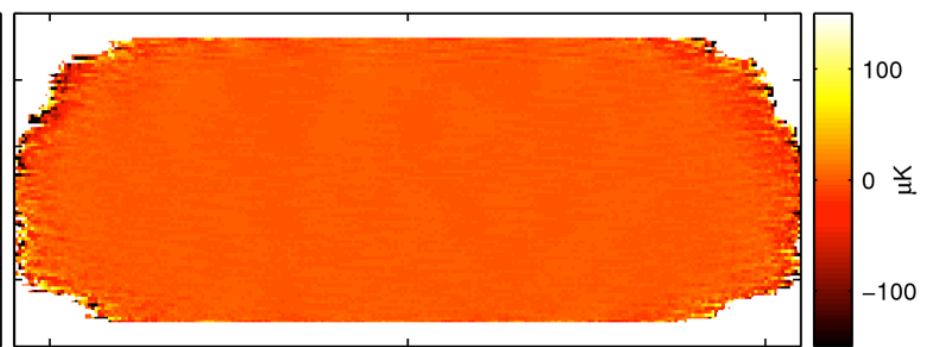


# BK15 95GHz Maps

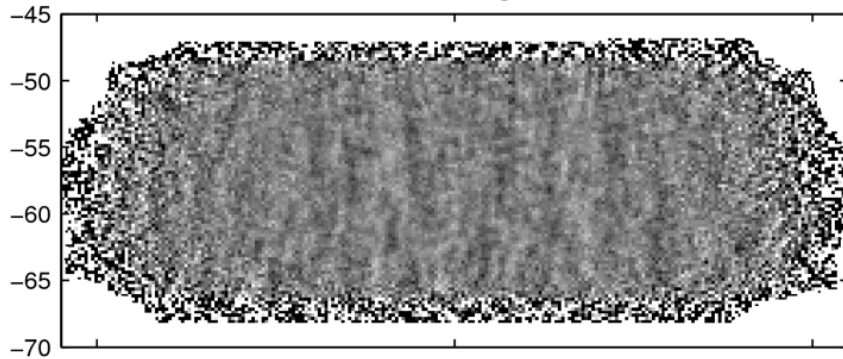
95 GHz T signal



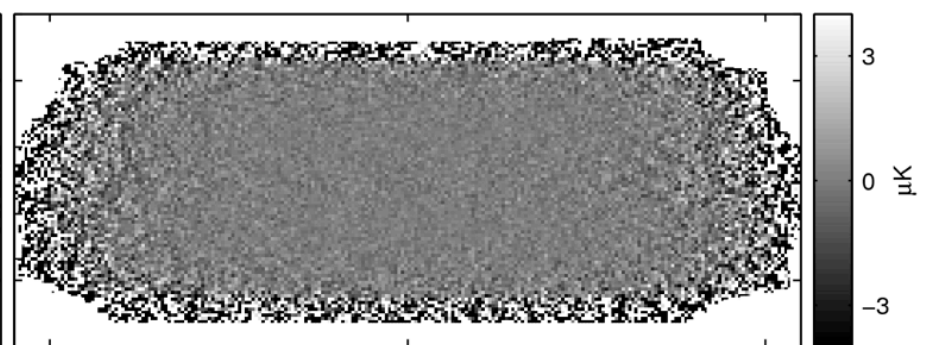
95 GHz T noise



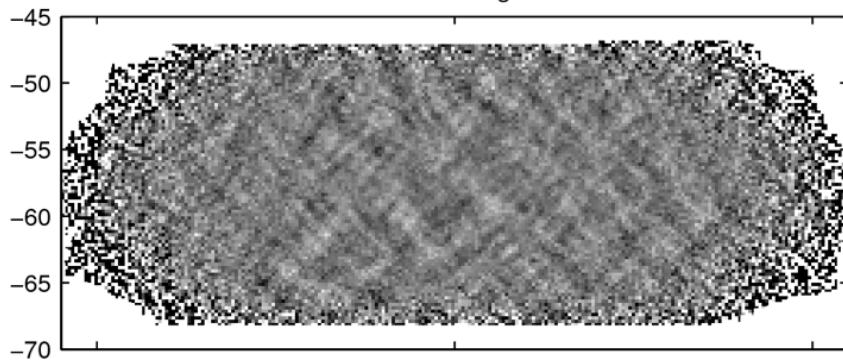
95 GHz Q signal



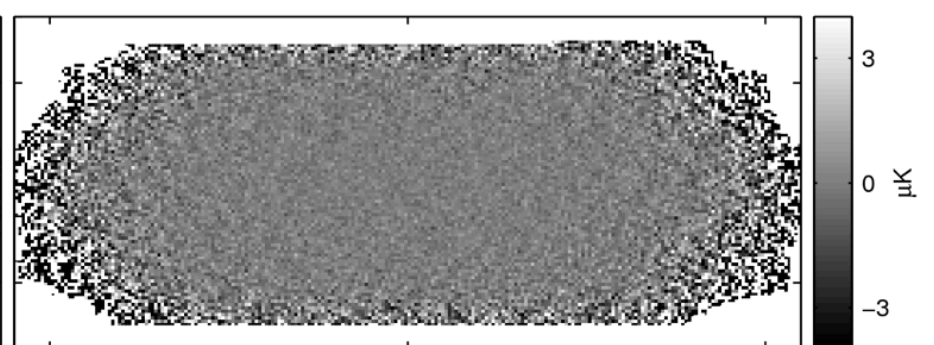
95 GHz Q noise



95 GHz U signal



95 GHz U noise



Declination [deg.]

50

0

-50

Right ascension [deg.]

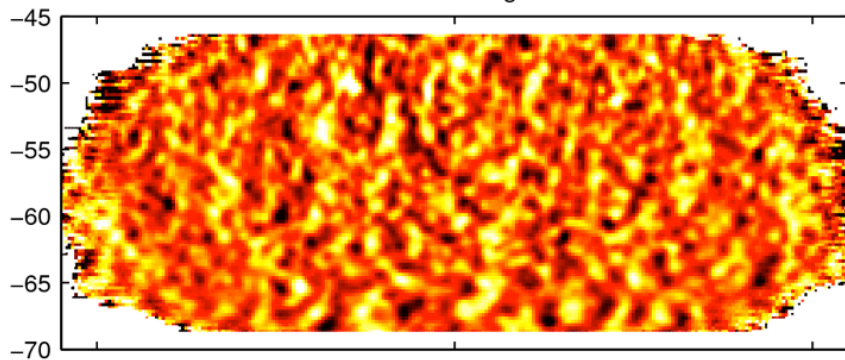
50

0

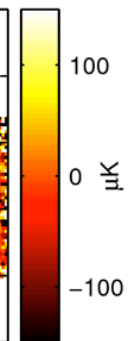
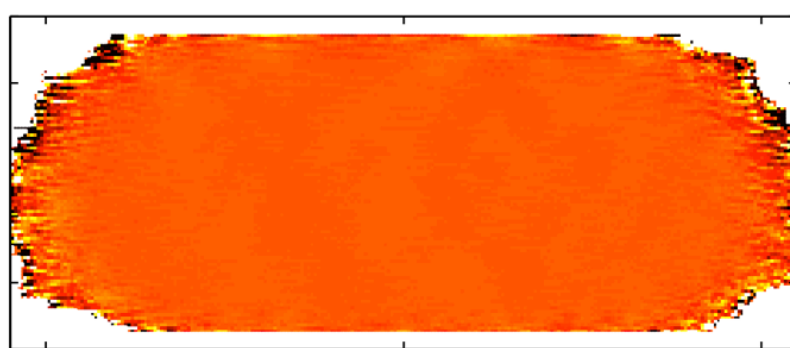
-50

# BK15 150GHz Maps

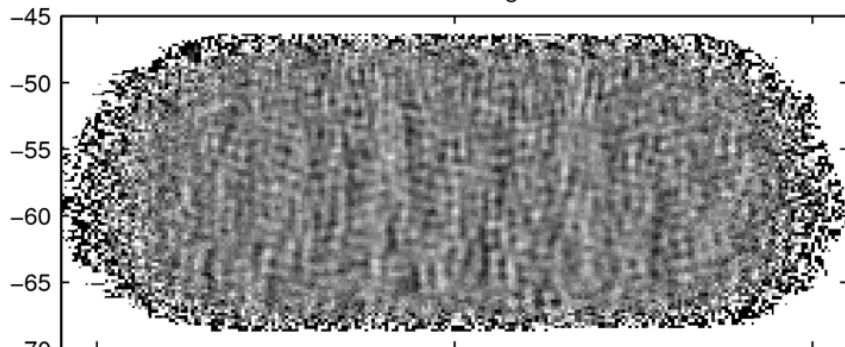
150 GHz T signal



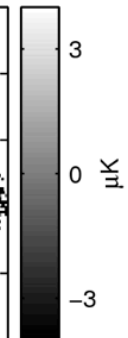
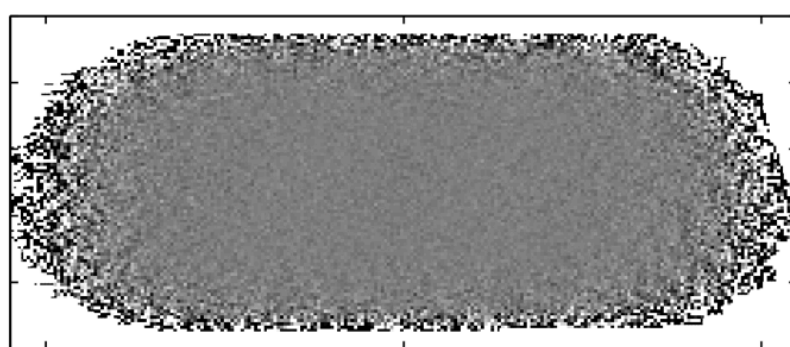
150 GHz T noise



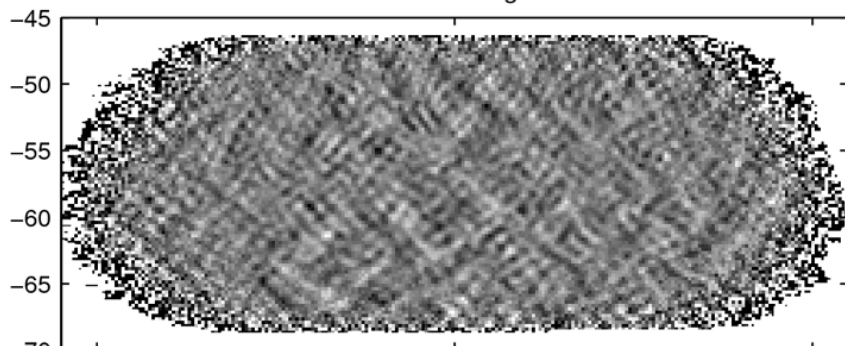
150 GHz Q signal



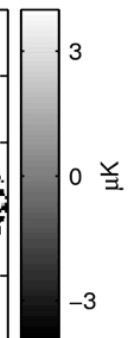
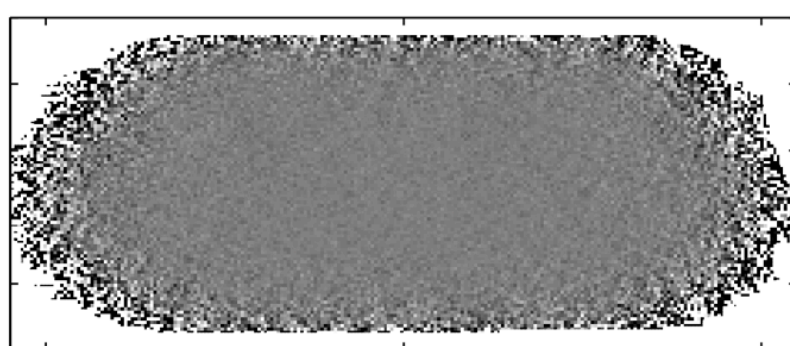
150 GHz Q noise



150 GHz U signal



150 GHz U noise



Declination [deg.]

50

0

-50

Right ascension [deg.]

50

0

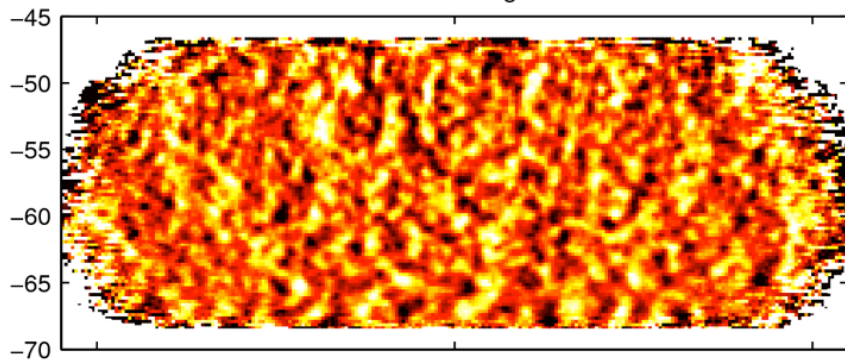
-50

BK15 150GHz – 2.8  $\mu\text{K arcmin}$

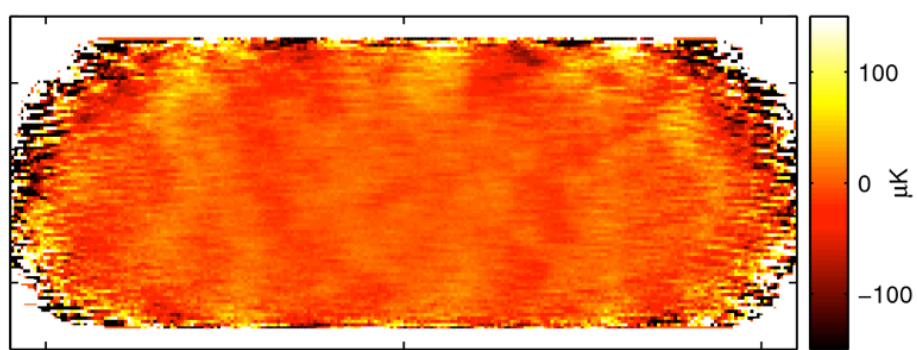


# BK15 220GHz Maps

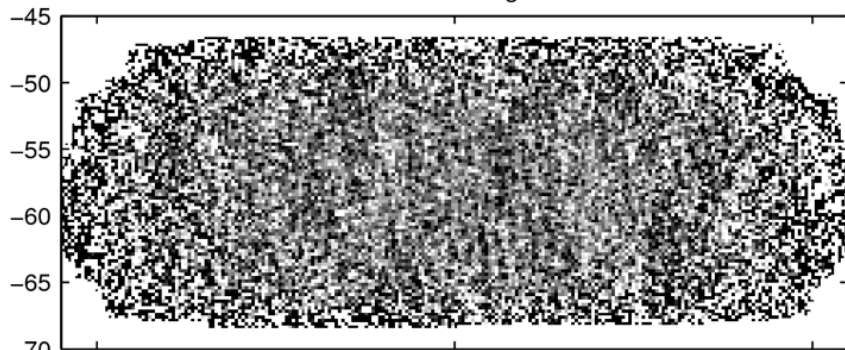
220 GHz T signal



220 GHz T noise



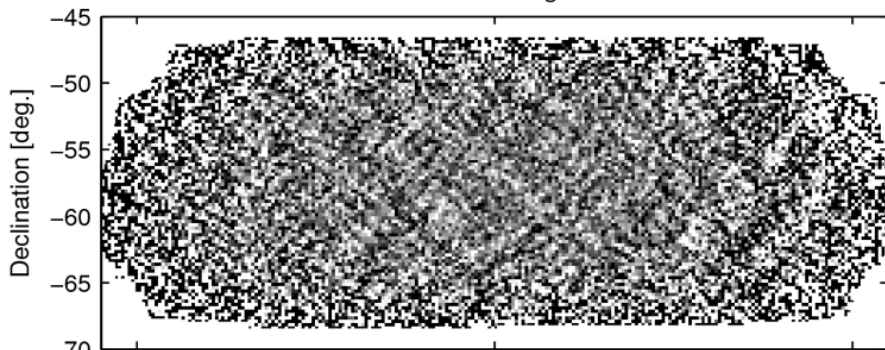
220 GHz Q signal



220 GHz Q noise



220 GHz U signal



220 GHz U noise



Right ascension [deg.]

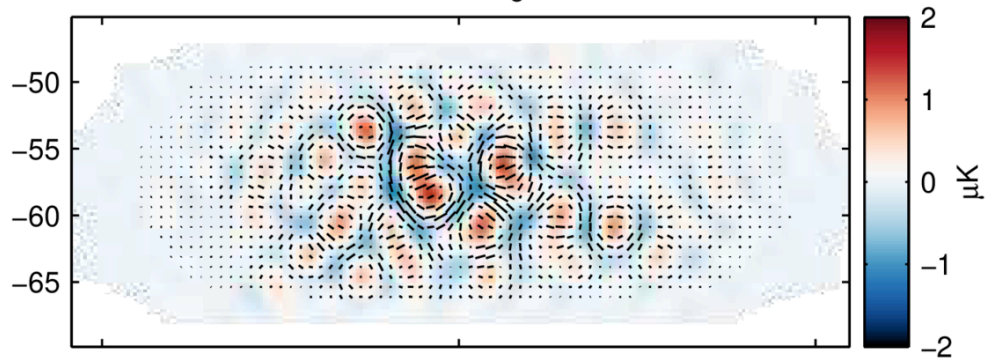
Declination [deg.]



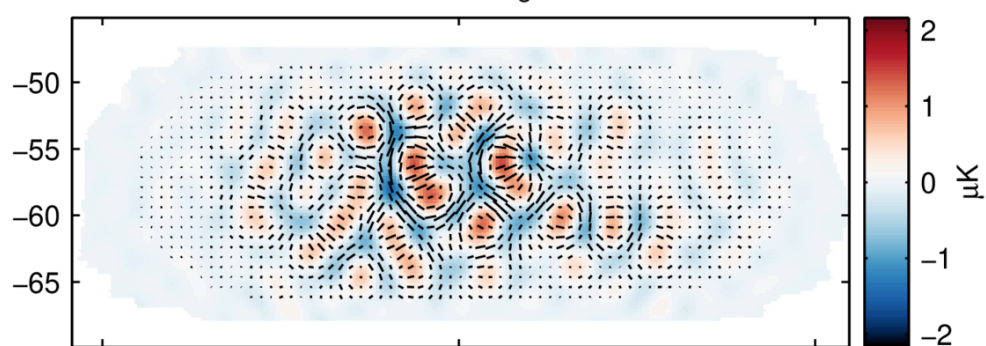


# Just for fun: Keck 2015 single season E-mode maps

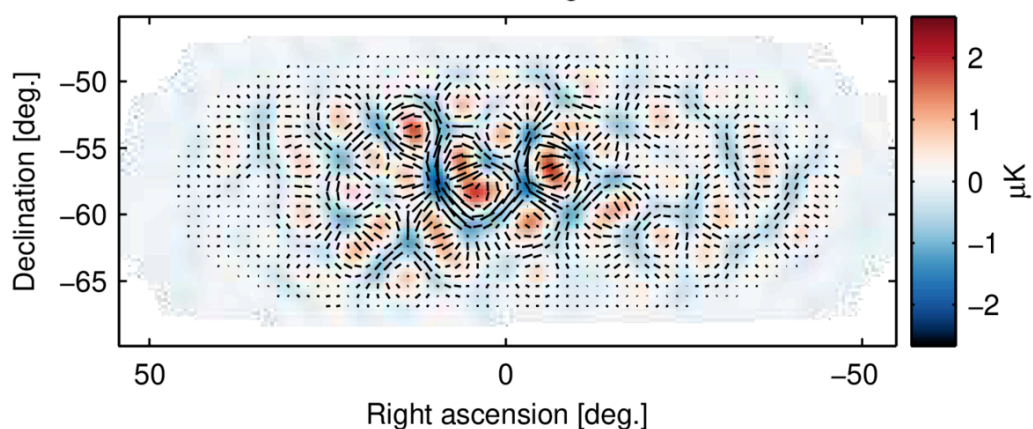
95 GHz E signal



150 GHz E signal



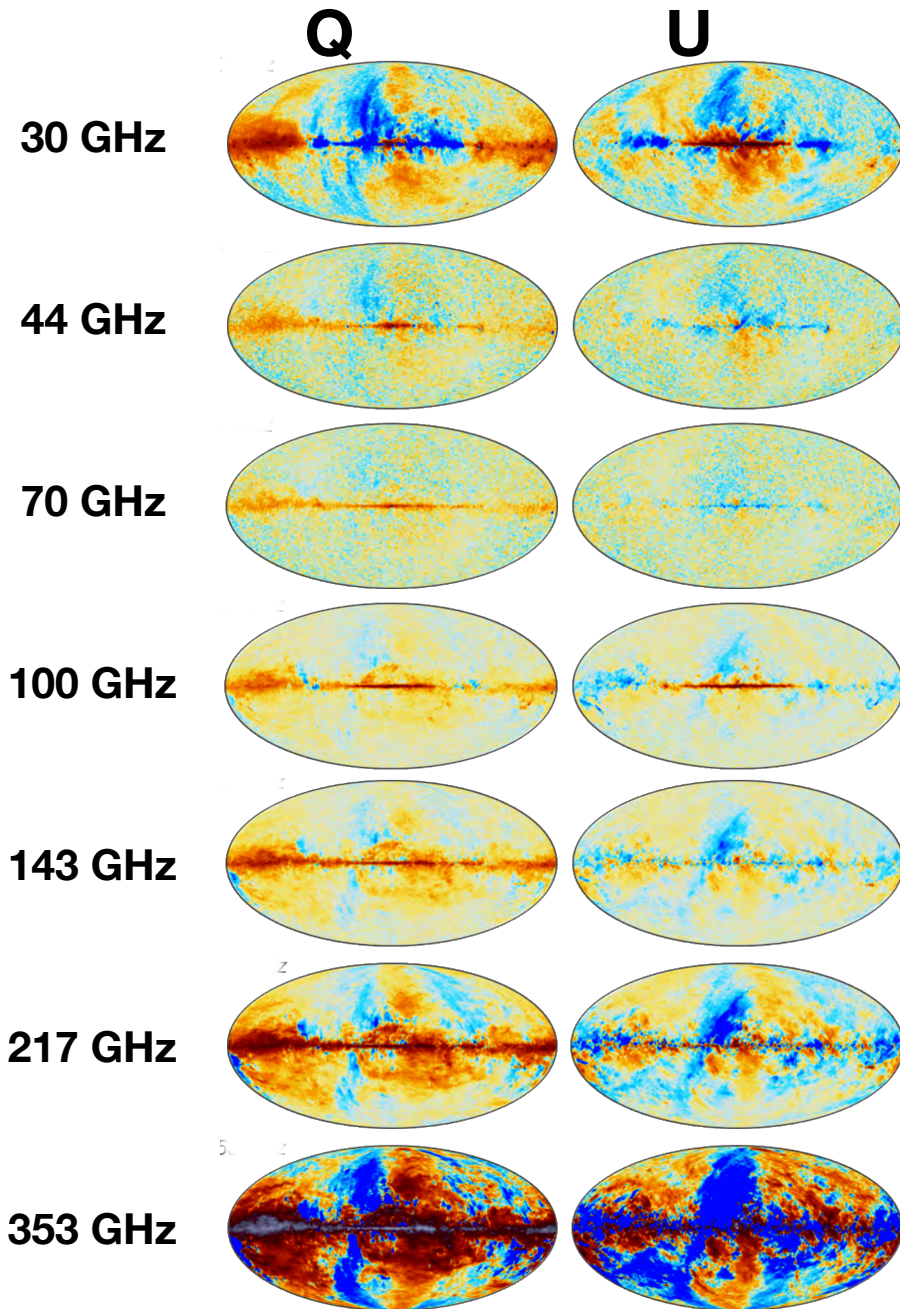
220 GHz E signal



This plot shows LCDM E-modes with high s/n at three frequencies from data taken in a single season!

← Already deeper than Planck 217 GHz

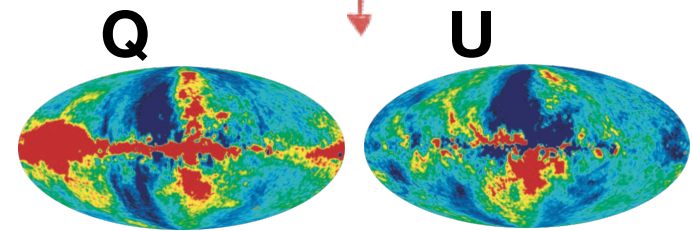
# Add to the mix: Planck at 7 frequencies and WMAP at 2 frequencies



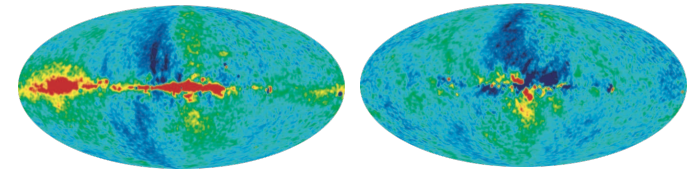
Polarized galactic  
**synchrotron**  
dominates  
at low frequencies



23 GHz



33 GHz



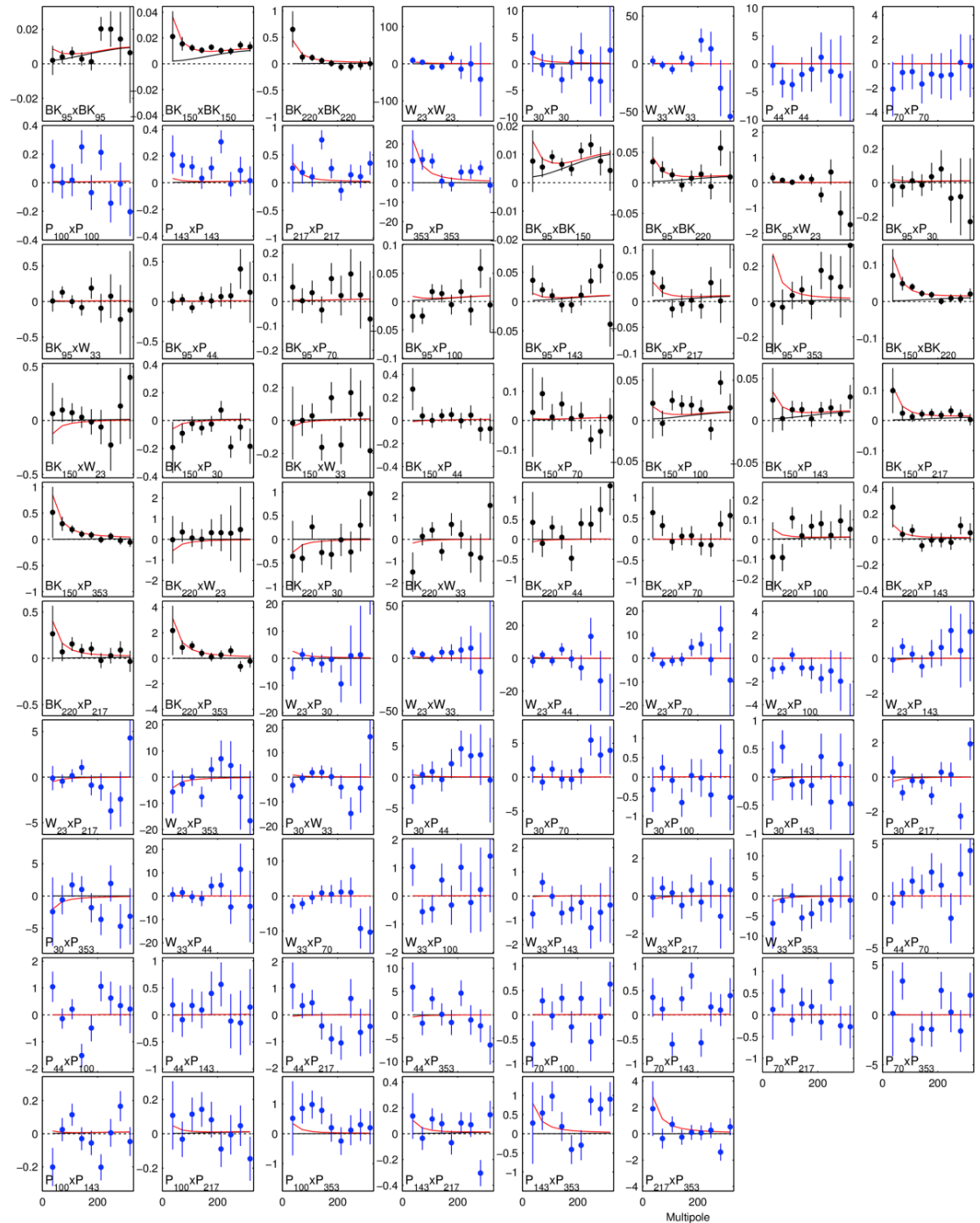
From arxiv 1212.5225

Polarized thermal  
emission ( $\sim 20\text{K}$ ) from  
galactic **dust** aligned in  
magnetic fields  
dominates  
at high frequencies



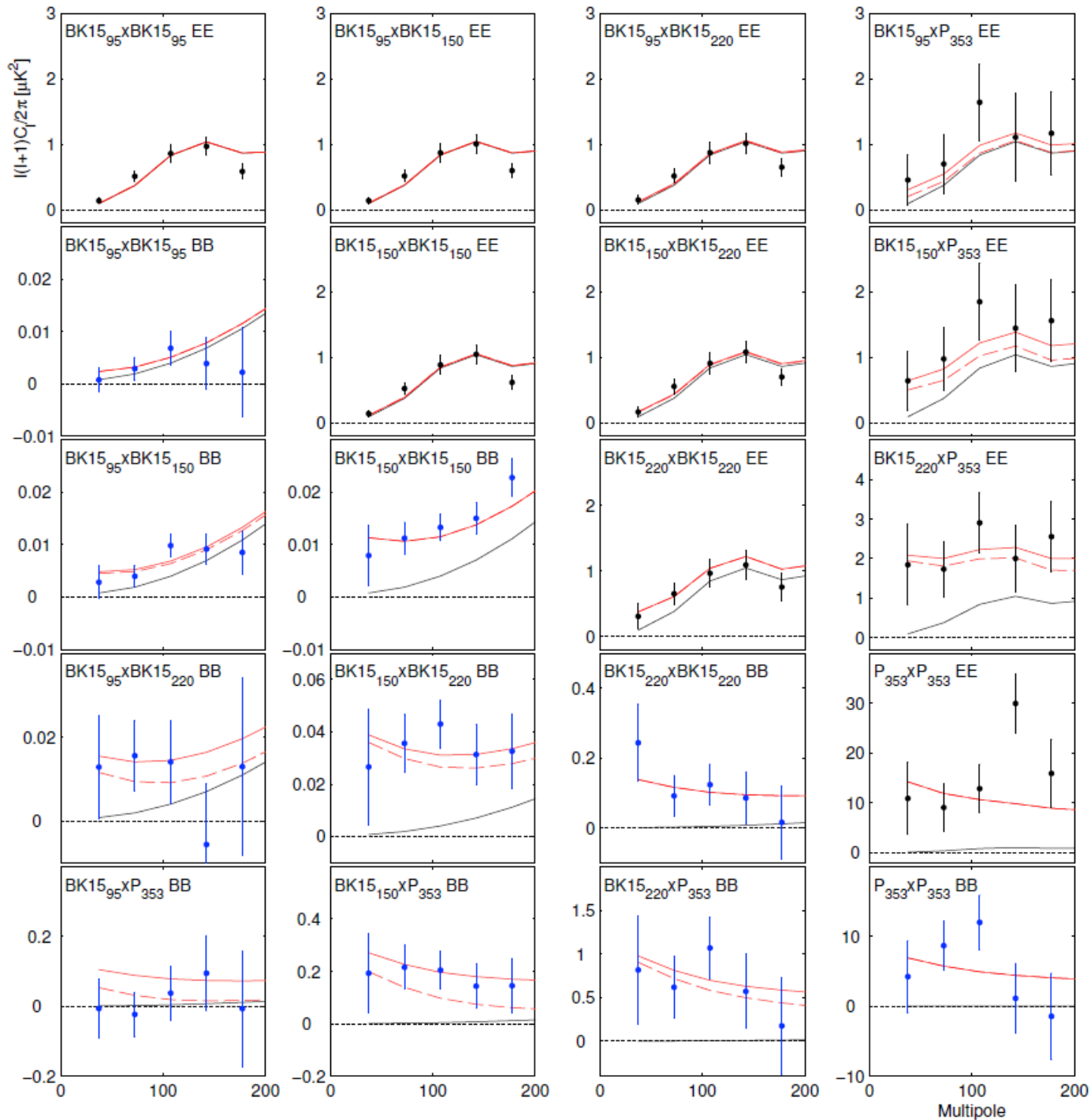
From arxiv 1502.01582

Take all possible auto- and cross spectra between the BICEP/Keck, WMAP, and Planck bands (78 of them)





# BK15+P353 Spectra



Upper/right plots are EE (black points)

Lower/left plots are BB (blue points)

220GHz auto/cross spectra are all new

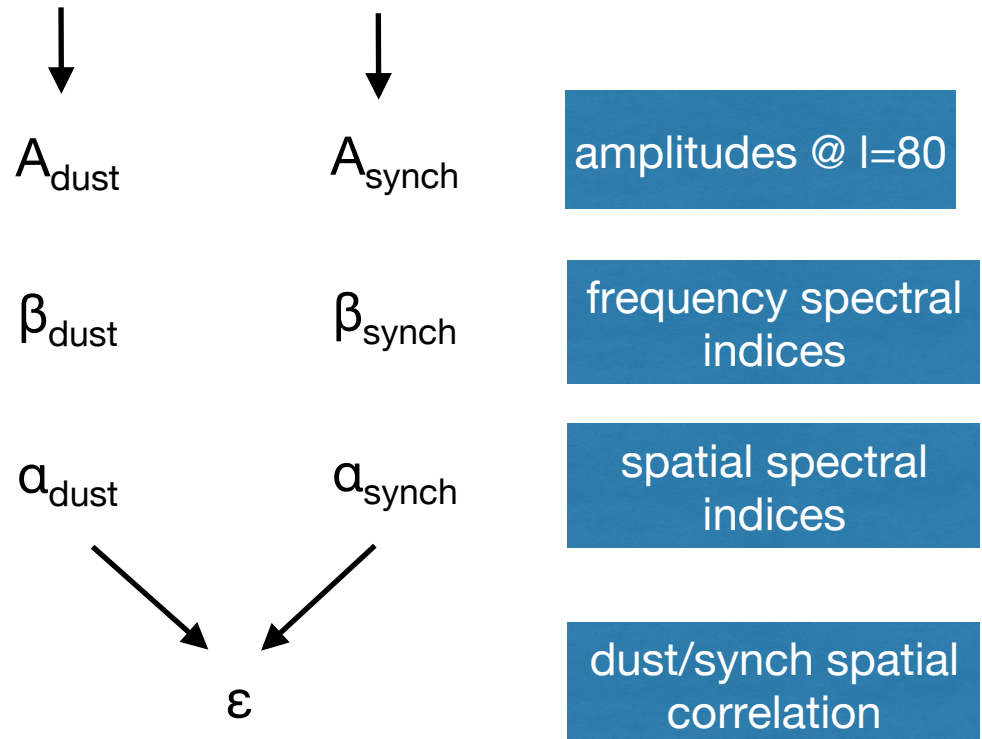
Red solid line is best fit multicomponent model from previous (BK14) analysis - It fits **all** the spectra

Chi-squared is OK - no evidence yet for non-Gaussianity of the dust pattern

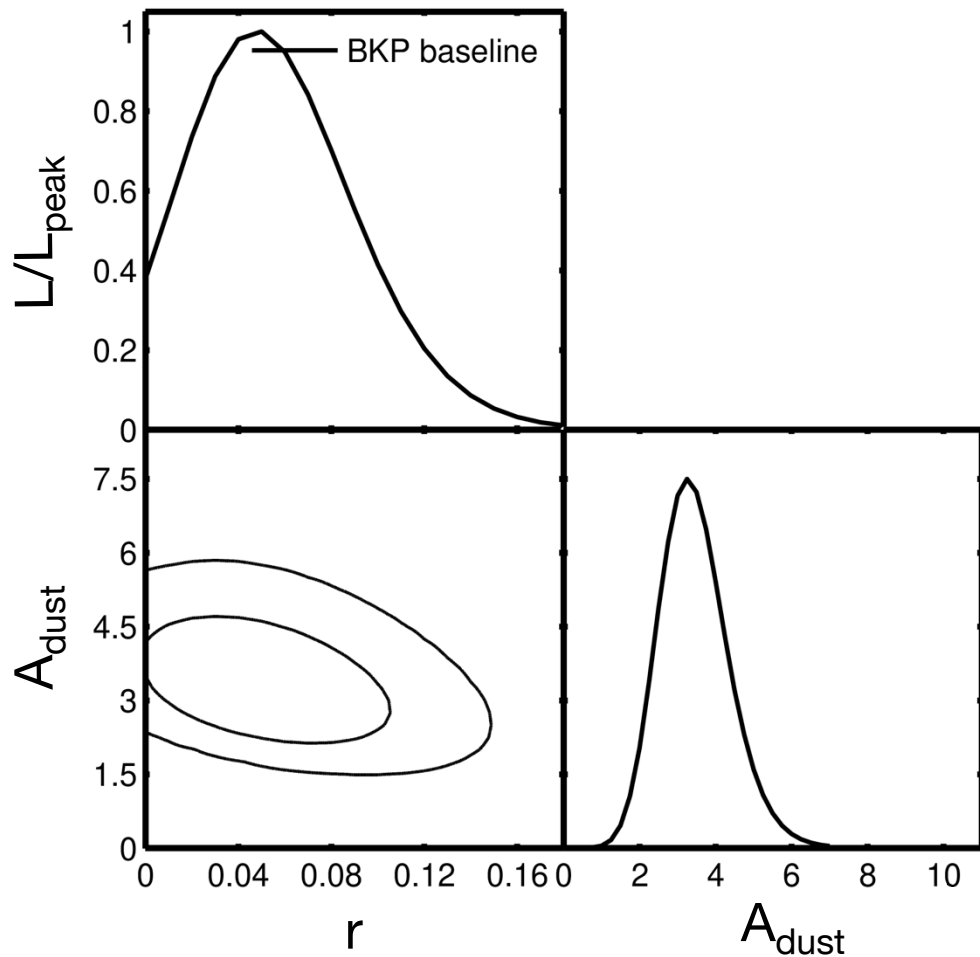
# Multicomponent parametric likelihood analysis

Take the joint likelihood of all the spectra simultaneously vs. model for BB that is the  $\Lambda$ CDM lensing expectation + 7 parameter foreground model + r

foreground model = dust + synchrotron

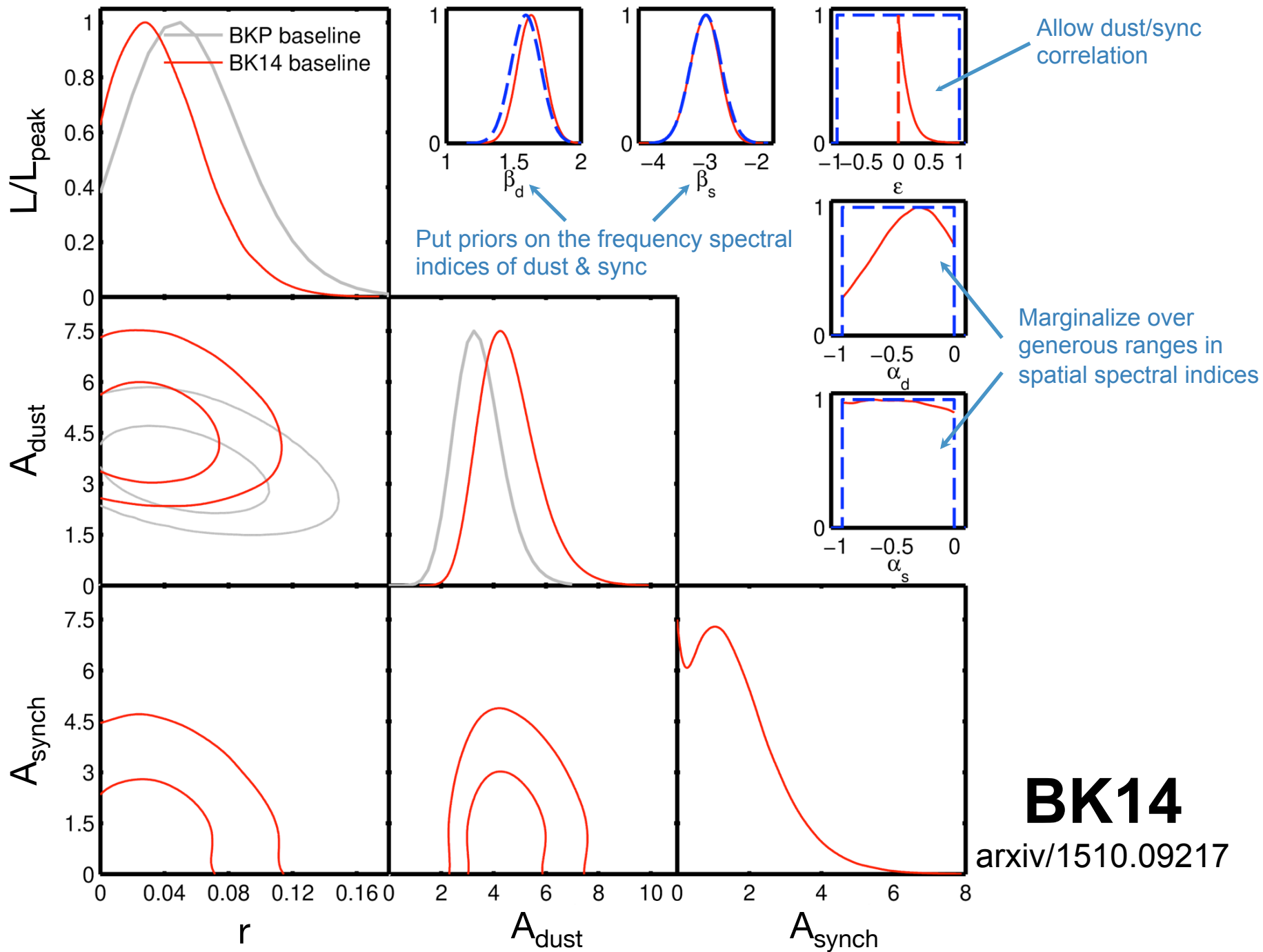


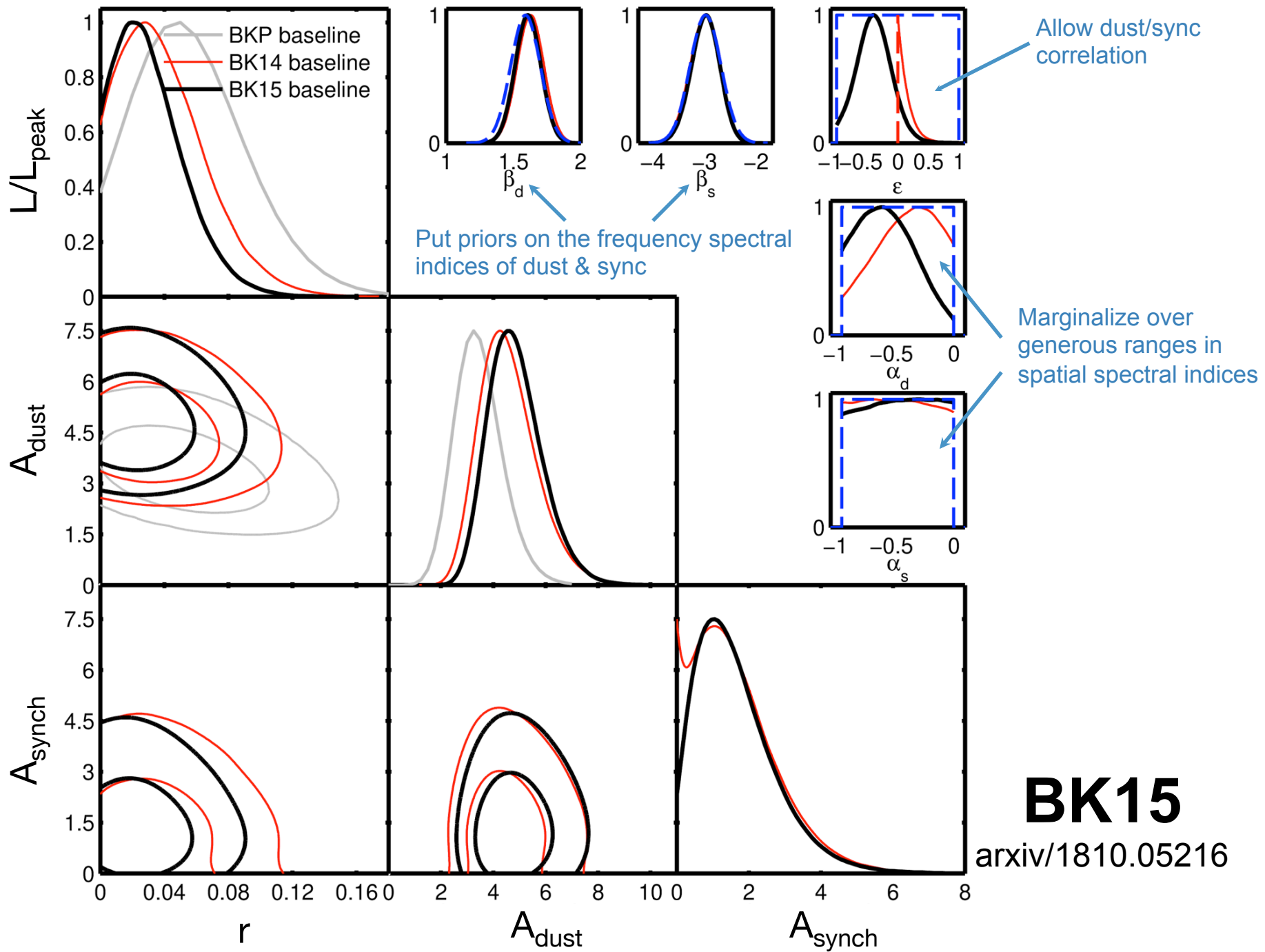




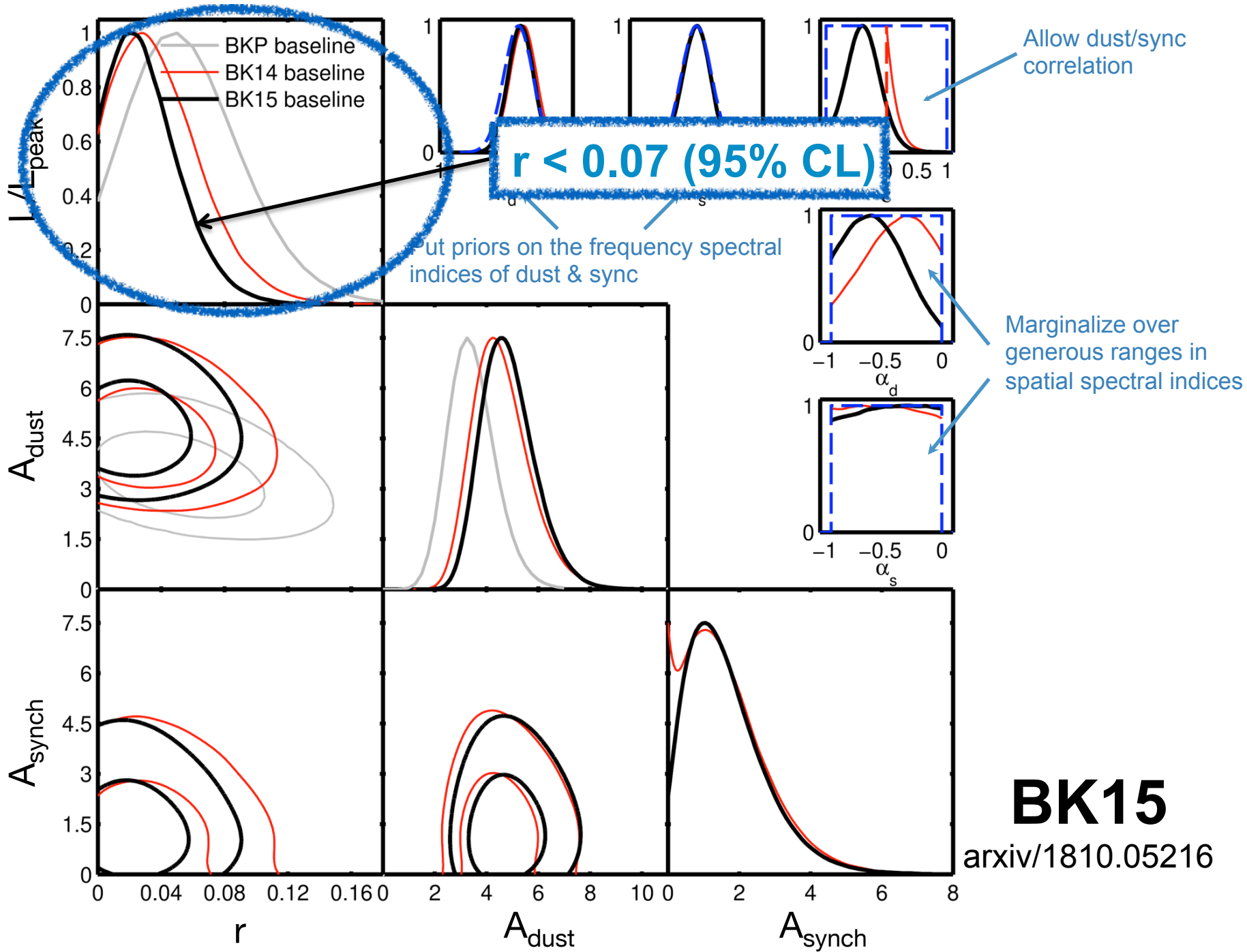
**BKP**

arxiv/1502.00612



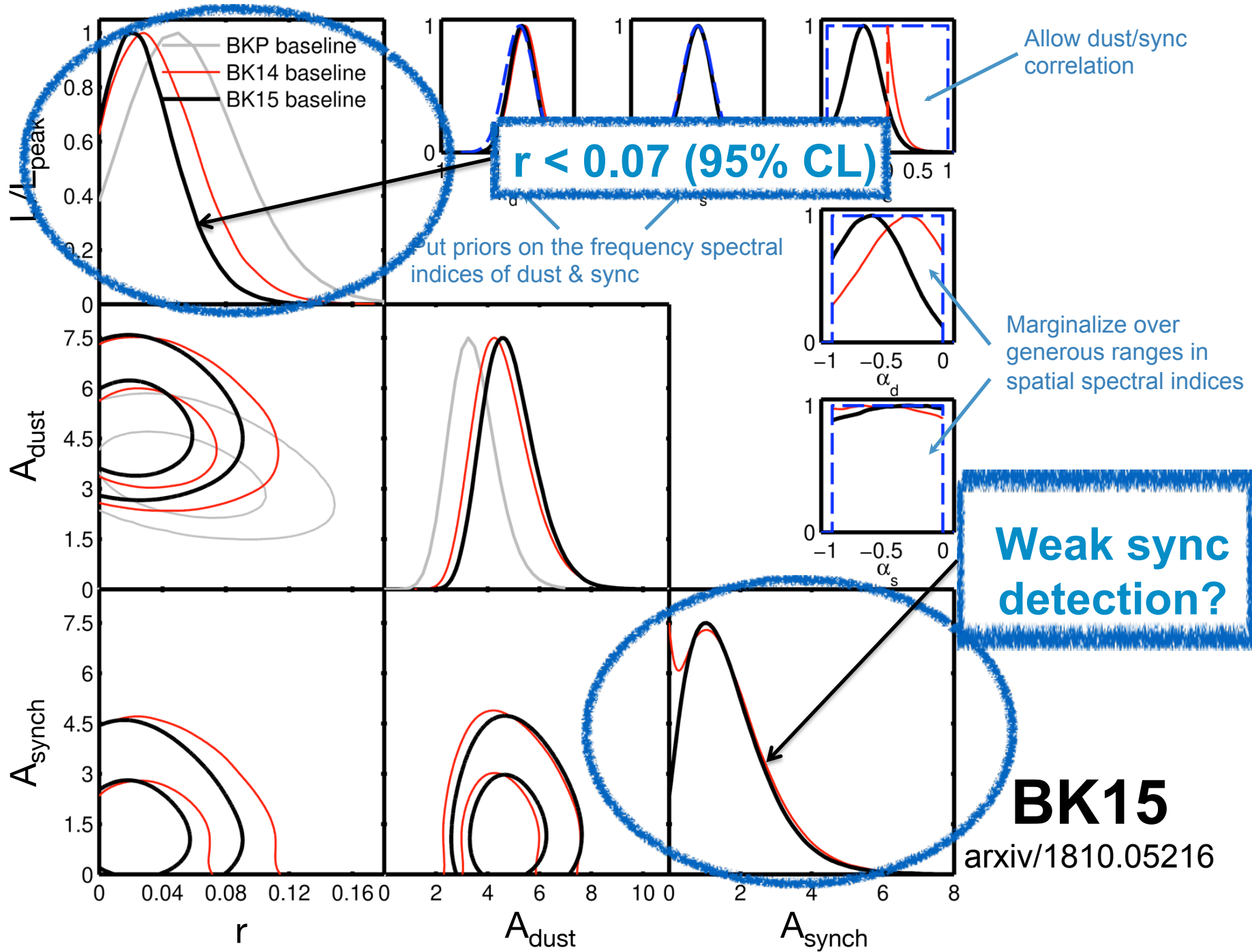




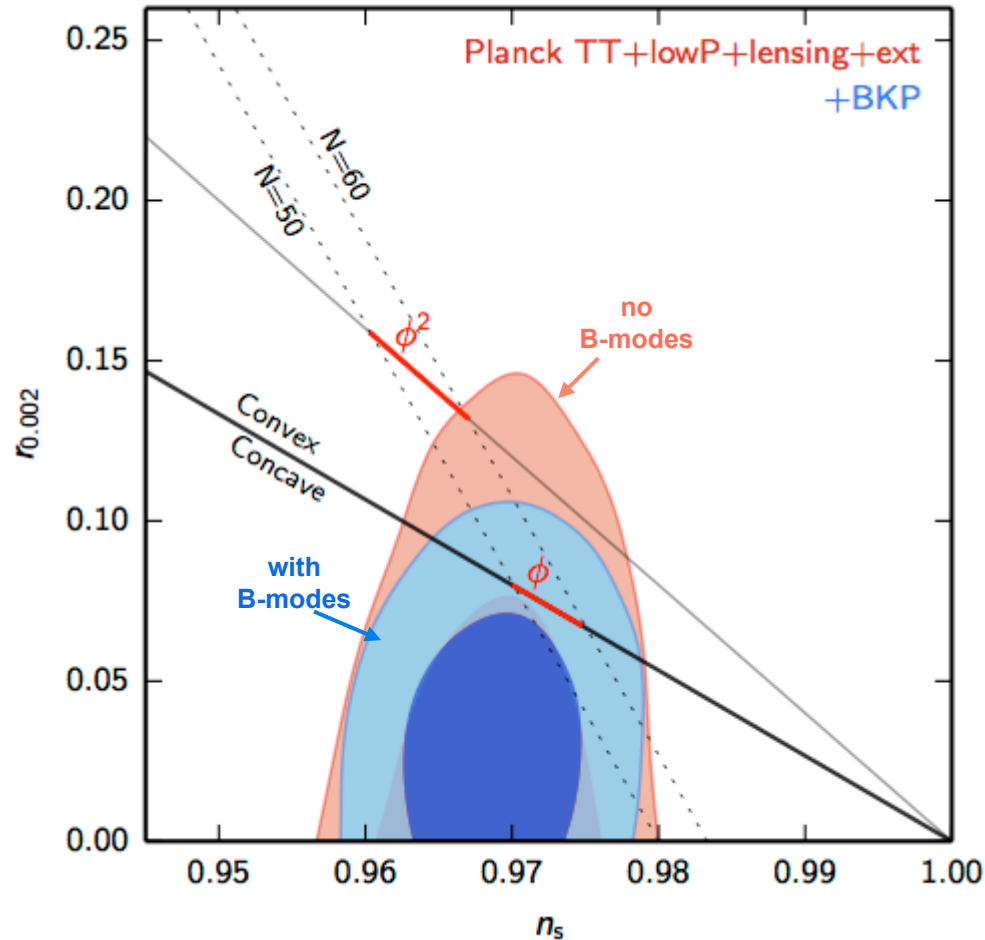


**BK15**

arxiv/1810.05216



# Adding in temperature



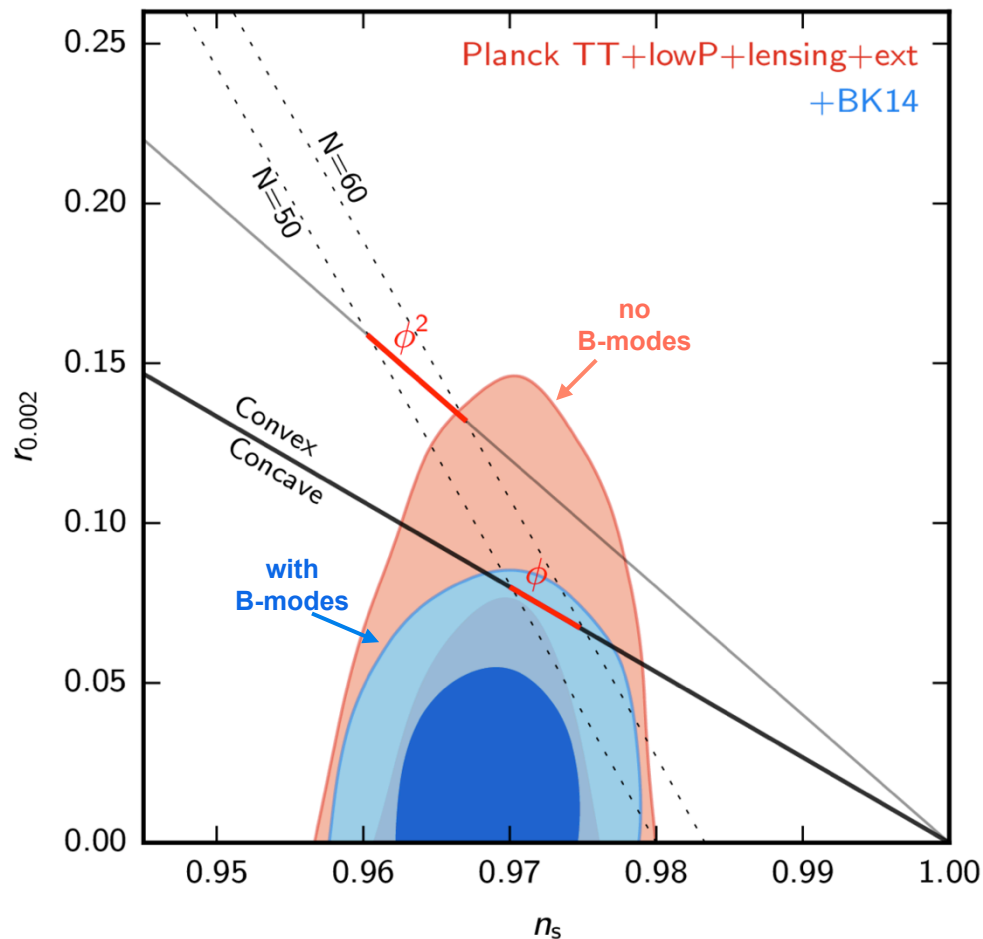
$r < 0.09$

**BKP**

arxiv/1502.00612



# Adding in temperature

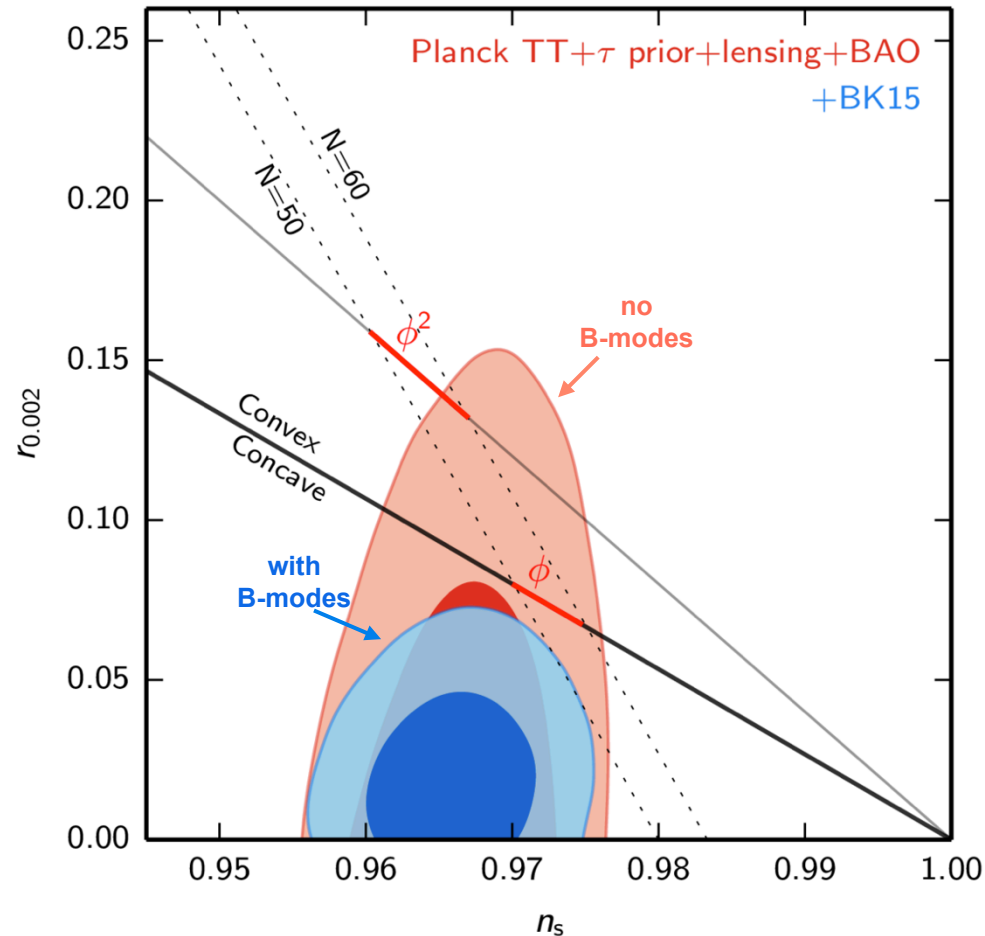


$r_{.05} < 0.07$

**BK14**

arxiv/1510.09217

# Adding in temperature



$r_{.05} < 0.06$

**BK15**

arxiv/1810.05216

# Why BK15 comes 3 years after BK14...

## Planck 2016

### Planck intermediate results. L. Evidence for spatial variation of the polarized thermal dust spectral energy distribution and implications for CMB $B$ -mode analysis

Planck Collaboration: N. Aghanimi<sup>51</sup>, M. Ashdown<sup>51,6</sup>, J. Aumont<sup>21,4</sup>, C. Baccigalupi<sup>74</sup>, M. Ballardini<sup>25,41,44</sup>, A. J. Banday<sup>82,9</sup>, R. B. Barreiro<sup>56</sup>, N. Bartolo<sup>22,57</sup>, S. Basak<sup>74</sup>, K. Benabed<sup>52,81</sup>, J.-P. Bernard<sup>82,9</sup>, M. Bersanelli<sup>28,42</sup>, P. Bielewicz<sup>71,97,4</sup>, A. Bonaldi<sup>59</sup>, L. Bonavera<sup>15</sup>, J. R. Bond<sup>8</sup>, J. Borrill<sup>11,78</sup>, F. R. Bouchet<sup>52,77</sup>, F. Boulanger<sup>51</sup>, A. Bracco<sup>44</sup>, C. Burigana<sup>41,26,44</sup>, E. Calabrese<sup>79</sup>, J.-F. Cardoso<sup>63,4,2</sup>, H. C. Chiang<sup>21,7</sup>, L. P. L. Colombo<sup>38,25</sup>, C. Combet<sup>66</sup>, B. Comis<sup>66</sup>, B. P. Crill<sup>58,10</sup>, A. Curio<sup>56,6,61</sup>, F. Cuttaia<sup>41</sup>, R. J. Davis<sup>59</sup>, P. de Bernardis<sup>27</sup>, A. de Rosa<sup>41</sup>, G. de Zotti<sup>38,74</sup>, J. Delabrouille<sup>1</sup>, J.-M. Delouis<sup>52,81</sup>, E. Di Valentino<sup>52,77</sup>, C. Dickinson<sup>59</sup>, J. M. Diego<sup>56</sup>, O. Dore<sup>58,10</sup>, M. Douspis<sup>51</sup>, A. Ducout<sup>52,50</sup>, X. Dupac<sup>32</sup>, S. Dusini<sup>27</sup>, G. Efstathiou<sup>51,53</sup>, F. Elsner<sup>19,52,81</sup>, T. A. Enßlin<sup>69</sup>, H. K. Eriksen<sup>54</sup>, E. Falgarone<sup>63</sup>, Y. Fantaye<sup>30,3</sup>, F. Finelli<sup>41,44</sup>, M. Fraix-Burnet<sup>40</sup>, A. A. Fraisse<sup>21</sup>, E. Franceschi<sup>41</sup>, A. Prolov<sup>76</sup>, S. Galotta<sup>40</sup>, S. Galli<sup>40</sup>, K. Ganga<sup>1</sup>, R. T. Génova-Santos<sup>35,14</sup>, M. Gerbino<sup>80,73,27</sup>, T. Ghosh<sup>51</sup>, M. Giard<sup>29</sup>, J. González-Nuevo<sup>15,56</sup>, K. M. Górski<sup>38,84</sup>, A. Gregorio<sup>29,40,48</sup>, A. Gruppioni<sup>51,44</sup>, J. E. Gudmundsson<sup>80,73,21</sup>, F. K. Hansen<sup>54</sup>, G. Helou<sup>10</sup>, D. Herranz<sup>26</sup>, E. Hivon<sup>52,81</sup>, Z. Huang<sup>8</sup>, A. H. Jaffe<sup>50</sup>, W. C. Jones<sup>21</sup>, E. Keihänen<sup>20</sup>, R. Kesitalo<sup>11</sup>, T. S. Kisner<sup>68</sup>, N. Krachmalnicoff<sup>58</sup>, M. Kunz<sup>13,51,3</sup>, H. Kurki-Suonio<sup>20,37</sup>, G. Lagache<sup>55,1</sup>, A. Lähteenmäki<sup>27,7</sup>, J.-M. Lamarca<sup>63</sup>, A. Lasenby<sup>6,61</sup>, M. Lattanzi<sup>26,45</sup>, C. R. Lawrence<sup>58</sup>, M. Le Jeune<sup>1</sup>, F. Levrier<sup>63</sup>, M. Liguori<sup>24,57</sup>, P. B. Lilje<sup>54</sup>, M. López-Cañiego<sup>32</sup>, P. M. Lubin<sup>22</sup>, J. F. Macías-Pérez<sup>66</sup>, G. Maggio<sup>40</sup>, D. Maino<sup>28,42</sup>, N. Mandolesi<sup>11,26</sup>, A. Mangilli<sup>31,62</sup>, M. Maris<sup>40</sup>, P. G. Martin<sup>6</sup>, E. Martínez-González<sup>26</sup>, S. Matarrese<sup>24,57,34</sup>, N. Mauri<sup>44</sup>, J. D. McEwen<sup>10</sup>, A. Melchiorri<sup>27,46</sup>, A. Mennella<sup>28,42</sup>, M. Migliaccio<sup>33,61</sup>, S. Mitra<sup>40,58</sup>, M.-A. Miville-Deschênes<sup>51,8</sup>, D. Molinar<sup>26,41,45</sup>, A. Moneti<sup>32</sup>, L. Montier<sup>82,9</sup>, G. Morgante<sup>41</sup>, A. Moss<sup>75</sup>, P. Naselsky<sup>72,31</sup>, H. U. Nørgaard-Nielsen<sup>12</sup>, C. A. Oxborrow<sup>12</sup>, L. Pagano<sup>27,46</sup>, D. Paoletti<sup>44,44</sup>, B. Partridge<sup>36</sup>, L. Patrizii<sup>44</sup>, O. Perdereau<sup>62</sup>, L. Perotto<sup>66</sup>, V. Pettorino<sup>25,52</sup>, F. Piacentini<sup>27</sup>, S. Płaczczynski<sup>62</sup>, G. Polenta<sup>4,39</sup>, J.-L. Puget<sup>51</sup>, J. P. Rachen<sup>16,69</sup>, M. Reinecke<sup>69</sup>, M. Remazeilles<sup>59,51,1</sup>, A. Renzi<sup>30,47</sup>, G. Rocha<sup>38,10</sup>, M. Rossetti<sup>28,42</sup>, G. Rouderer<sup>1,63,58</sup>, J. A. Rubiño-Martín<sup>55,14</sup>, B. Ruiz-Granados<sup>83</sup>, L. Salvati<sup>43</sup>, M. Sandri<sup>41</sup>, M. Savelainen<sup>19,37</sup>, D. Scott<sup>17</sup>, C. Sirignano<sup>24,57</sup>, G. Sirri<sup>44</sup>, L. Stanco<sup>37</sup>, A.-S. Suur-Uski<sup>20,37</sup>, J. A. Tauber<sup>13</sup>, M. Tenti<sup>61</sup>, L. Toffolatti<sup>15,56,41</sup>, M. Tomasi<sup>26,42</sup>, M. Tristram<sup>62</sup>, J. Valiviita<sup>20,37</sup>, F. Vansyngel<sup>51</sup>, F. Van Tent<sup>61</sup>, P. Vielva<sup>56</sup>, B. D. Wandelt<sup>52,81,22</sup>, I. K. Wehus<sup>58,54</sup>, A. Zacchei<sup>40</sup>, and A. Zonca<sup>22</sup>

(Affiliations can be found after the references)

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#### ABSTRACT

The characterization of the Galactic foregrounds has been shown to be the main obstacle in the challenging quest to detect primordial  $B$ -modes in the polarized microwave sky. We make use of the *Planck*-HFI 2015 data release at high frequencies to place new constraints on the properties of the polarized thermal dust emission at high Galactic latitudes. Here, we specifically study the spatial variability of the dust polarized spectral energy distribution (SED), and its potential impact on the determination of the tensor-to-scalar ratio,  $r$ . We use the correlation ratio of the  $C_{\ell}^{PP}$  angular power spectra between the 217- and 353-GHz channels as a tracer of these potential variations, computed on different high Galactic latitude regions, ranging from 80% to 20% of the sky. The new insight from *Planck* data is a departure of the correlation ratio from unity that cannot be attributed to a spurious decorrelation due to the cosmic microwave background, instrumental noise, or instrumental systematics. The effect is marginally detected on each region, but the statistical combination of all the regions gives more than 99% confidence for this variation in polarized dust properties. In addition, we show that the decorrelation increases when there is a decrease in the mean column density of the region of the sky being considered, and we propose a simple power-law empirical model for this dependence, which matches what is seen in the *Planck* data. We explore the effect that this measured decorrelation has on simulations of the BICEP2-Keck Array/*Planck* analysis and show that the 2015 constraints from those data still allow a decorrelation between the dust at 150 and 353 GHz of the order of the one we measure. Finally, using simplified models, we show that either spatial variation of the dust SED or of the dust polarization angle could produce decorrelations between 217- and 353-GHz data similar to those we observe in the data.

A departure of the correlation ratio from unity that cannot be attributed to a spurious decorrelation due to the cosmic microwave background, instrumental noise, or instrumental systematics... detected at more than 99% confidence

arxiv/1606.07335

## Planck 2018

### Planck 2018 results. XI. Polarized dust foregrounds

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(Affiliations can be found after the references)

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#### ABSTRACT

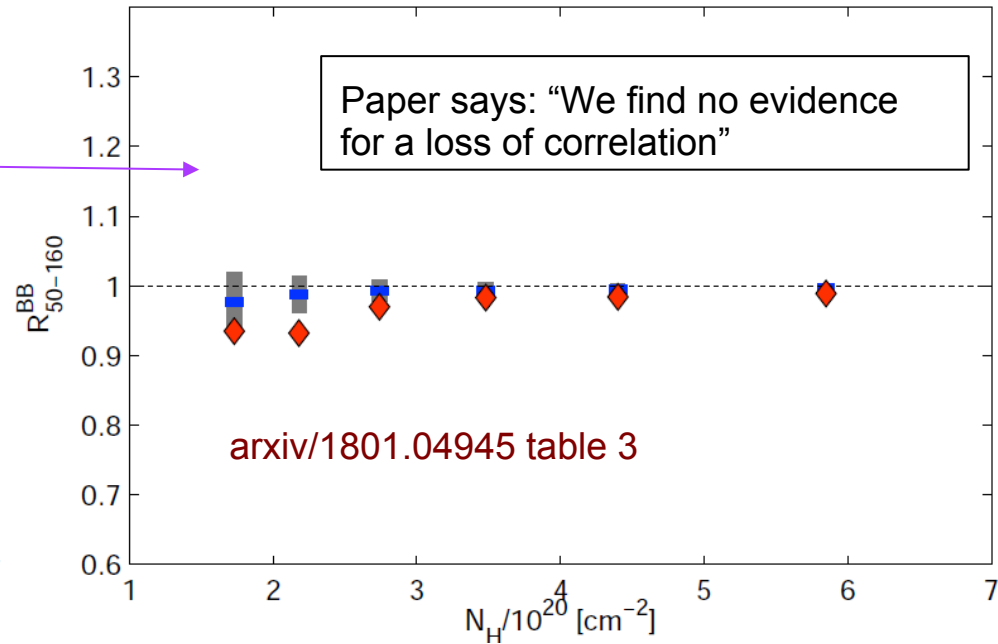
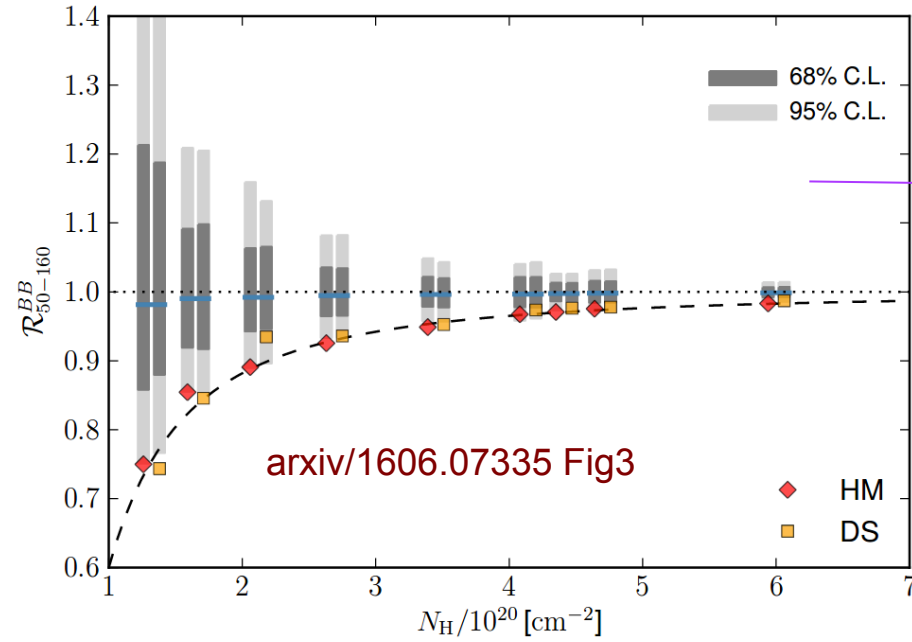
The study of polarized dust emission has become entwined with the analysis of the cosmic microwave background (CMB) polarization in the quest for the curl-like  $B$ -mode polarization from primordial gravitational waves and the low-multipole  $E$ -mode polarization associated with the reionization of the Universe. We use the new *Planck* PR3 maps to characterize Galactic dust emission at high latitudes as a foreground to the CMB polarization and use end-to-end simulations to compute uncertainties and assess the statistical significance of our measurements. We present *Planck*  $EE$ ,  $BB$ , and  $TE$  power spectra of dust polarization at 353 GHz for a set of six nested high-Galactic-latitude sky regions covering from 24 to 71% of the sky. We present power-law fits to the angular power spectra, yielding evidence for statistically significant variations of the exponents over sky regions and a difference between the values for the  $EE$  and  $BB$  spectra, which for the largest sky region are  $\alpha_{EE} = -2.42 \pm 0.02$  and  $\alpha_{BB} = -2.54 \pm 0.02$ , respectively. The spectra show that the  $TE$  correlation and  $E/B$  power asymmetry discovered by *Planck* extend to low multipoles that were not included in earlier *Planck* polarization papers due to residual data systematics. We also report evidence for a positive  $TB$  dust signal. Combining data from *Planck* and WMAP, we determine the amplitudes and spectral energy distributions (SEDs) of polarized foregrounds, including the correlation between dust and synchrotron polarized emission, for the six sky regions as a function of multipole. This quantifies the challenge of the component-separation procedure that is required for measuring the low- $\ell$  reionization CMB  $E$ -mode signal and detecting the reionization and recombination peaks of primordial CMB  $B$  modes. The SED of polarized dust emission is fit well by a single-temperature modified blackbody emission law from 353 GHz to below 70 GHz. For a dust temperature of 19.6 K, the mean dust spectral index for dust polarization is  $\beta_d^p = 1.53 \pm 0.02$ . The difference between indices for polarization and total intensity is  $\beta_d^p - \beta_d^I = 0.05 \pm 0.03$ . By fitting multi-frequency cross-spectra between *Planck* data at 100, 143, 217, and 353 GHz, we examine the correlation of the dust polarization maps across frequency. We find no evidence for a loss of correlation and provide lower limits to the correlation ratio that are tighter than values we derive from the correlation of the 217- and 353-GHz maps alone. If the *Planck* limit on decorrelation for the largest sky region applies to the smaller sky regions observed by sub-orbital experiments, then frequency decorrelation of dust polarization might not be a problem for CMB experiments aiming at a primordial  $B$ -mode detection limit on the tensor-to-scalar ratio  $r \approx 0.01$  at the recombination peak. However, the *Planck* sensitivity precludes identifying how difficult the component-separation problem will be for more ambitious experiments targeting lower limits on  $r$ .

We find no evidence for a loss of correlation... might not be a problem for CMB experiments aiming at a primordial B-mode detection limit on the tensor-to-scalar ratio  $r \sim 0.01$ ...

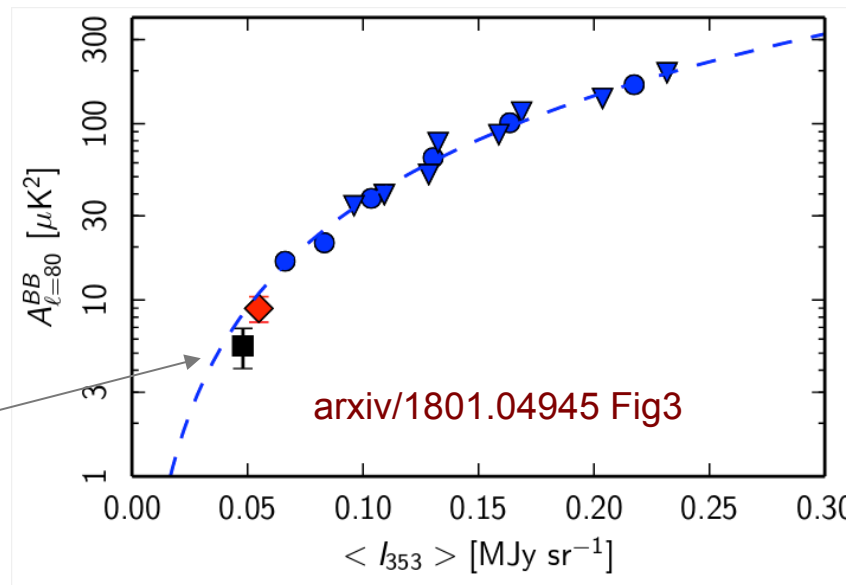
arxiv/1801.04945



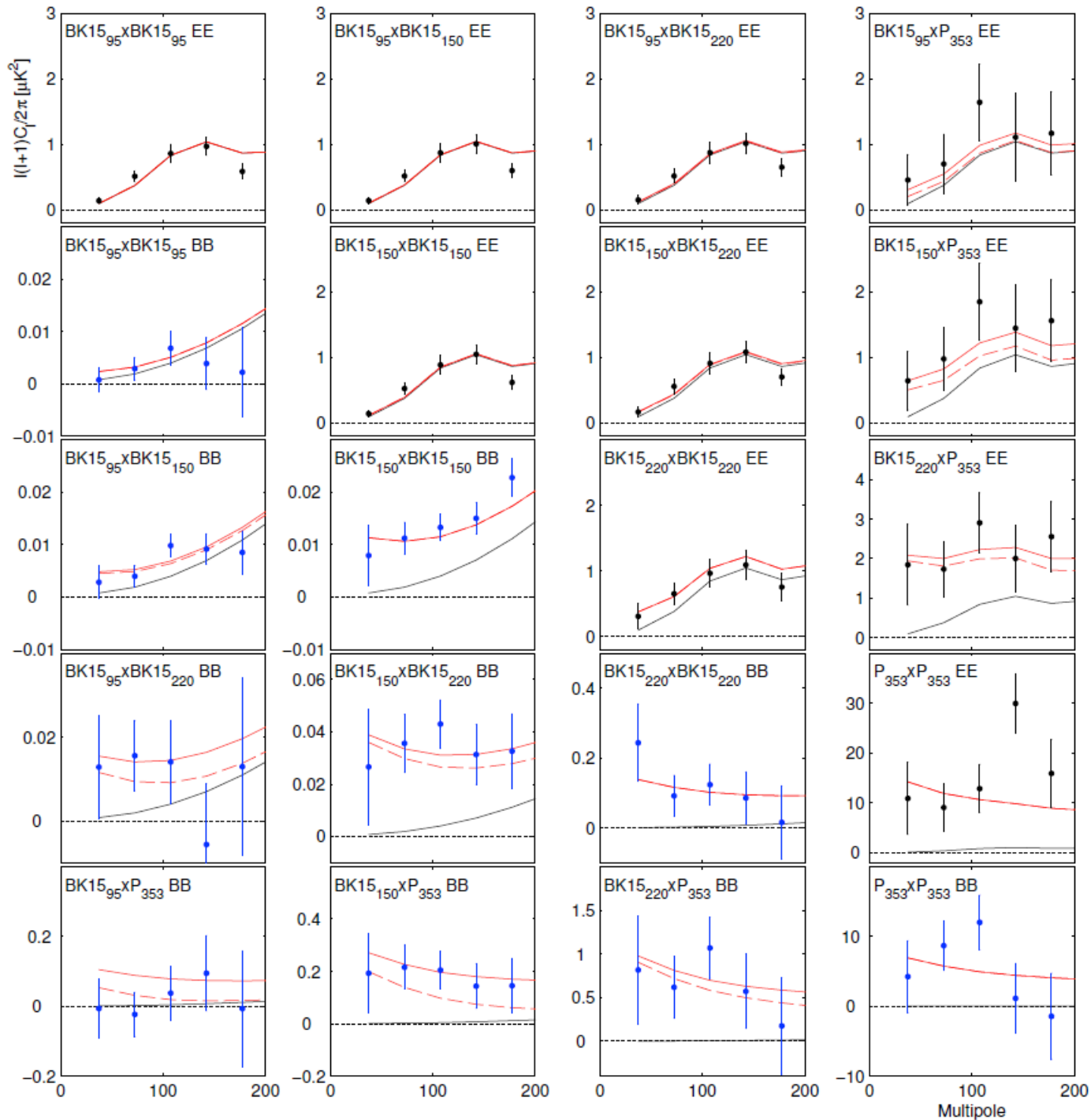
# Latest Planck Dust Analysis



BK patch a few times cleaner than extrapolation



# BK15+P353 Spectra



Upper/right plots are EE (black points)

Lower/left plots are BB (blue points)

220GHz auto/cross spectra are all new

Red solid line is best fit multicomponent model from previous (BK14) analysis - It fits **all** the spectra

Red dashed line is same model but with strong decorrelation - better for 95x353, worse for 150x353

Need better data to say for sure

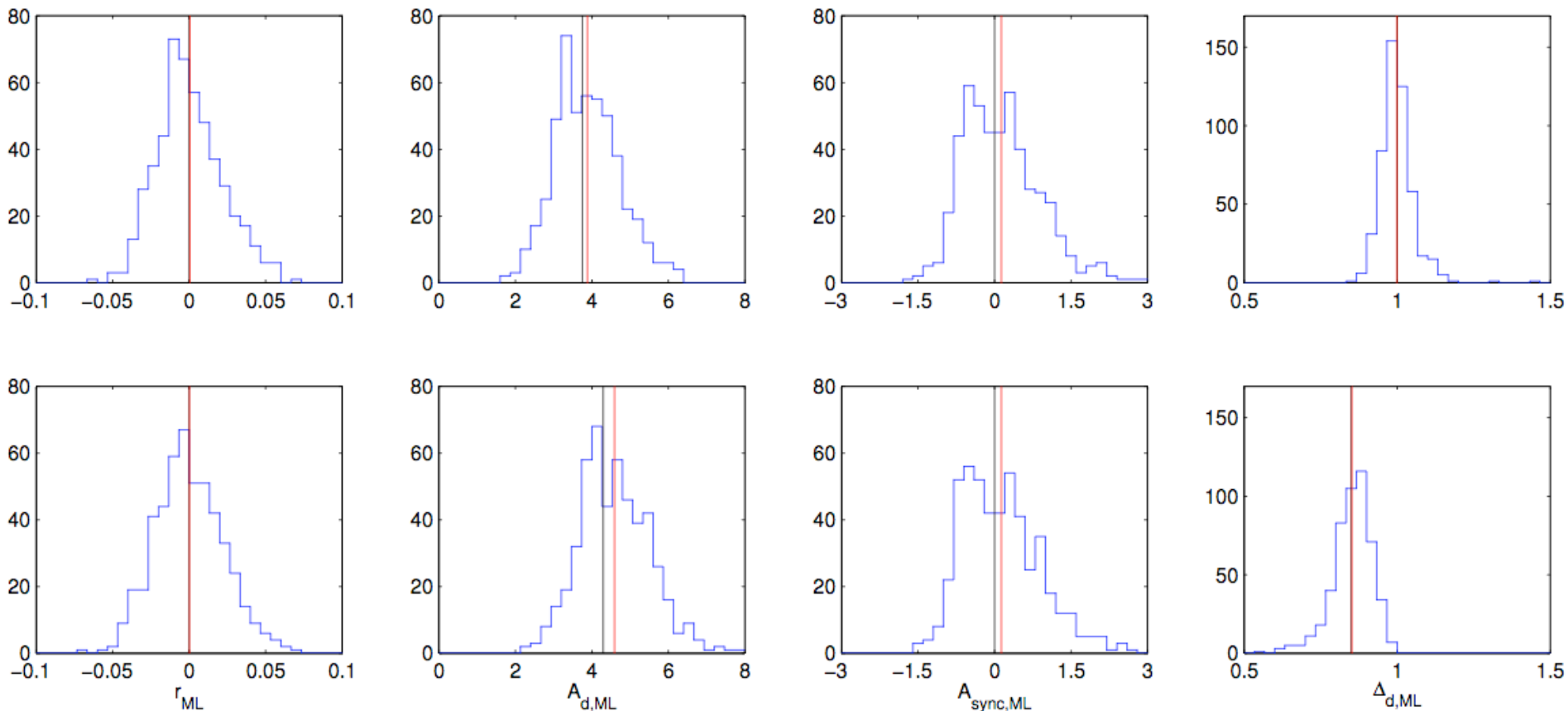
# Include dust correlation parameter?

- The standard BK15 marginalized likelihood analysis (COSMOMC style) is unbiased for  $r$  – i.e. when run on lensed-LCDM+dust+noise sims 50% of the  $r$  curves peak at zero.
- However, if add a dust correlation parameter and restrict to physical range ( $<1$ ) then  $r$  becomes biased – i.e. 72% of curves peak at zero
- This is because  $r$  and dust correlation parameter are partially degenerate
- Can run alternate maximum likelihood searches where parameters are allowed to take unphysical values ( $r < 0$  and dust correlation  $> 1$ )
- These really should be unbiased...



# BK15 ML Search Results for Sims

Upper = standard lensed-LCDM+dust+noise sims



Lower = lensed-LCDM+dust+noise sims with dust decorrelation

$$\Delta_d = \frac{\mathcal{D}_{80}(217 \times 353)}{\sqrt{\mathcal{D}_{80}(217 \times 217)\mathcal{D}_{80}(353 \times 353)}}$$

# BK15 ML Search Results for Sims

Also run ML search on sims containing 3<sup>rd</sup> party foreground models which do not necessarily conform to the foreground parameterization we are using in the re-fit. For the models considered so far bias is small compared to  $\sigma(r)$  – additional models are welcome.

Model	$\overline{A_d}$ ( $\mu\text{K}^2$ )	$\overline{A_s}$ ( $\mu\text{K}^2$ )	$\beta_d$ prior	$\sigma(r), \bar{r}/\sigma(r)$ $\beta_d$ free	with decorr.
Gaussian	3.8	0.1	0.020, $+0.1\sigma$	0.023, $0.0\sigma$	0.021, $+0.0\sigma$
PySM 1	10.9	1.1	0.026, $+0.2\sigma$	0.028, $+0.2\sigma$	0.028, $+0.1\sigma$
PySM 2	24.2	0.9	0.028, $+0.1\sigma$	0.029, $+0.1\sigma$	0.032, $+0.1\sigma$
PySM 3	12.1	1.1	(0.030, $+0.4\sigma$ )	0.031, $+0.1\sigma$	(0.032, $+0.2\sigma$ )
MHDv2	2.9	5.6	0.020, $+0.2\sigma$	0.027, $-0.2\sigma$	0.021, $-0.1\sigma$
G. Decorr.	4.6	0.1	(0.023, $+1.5\sigma$ )	(0.026, $+1.3\sigma$ )	0.022, $+0.0\sigma$

# 2016 onwards: BICEP3 “Super receiver”

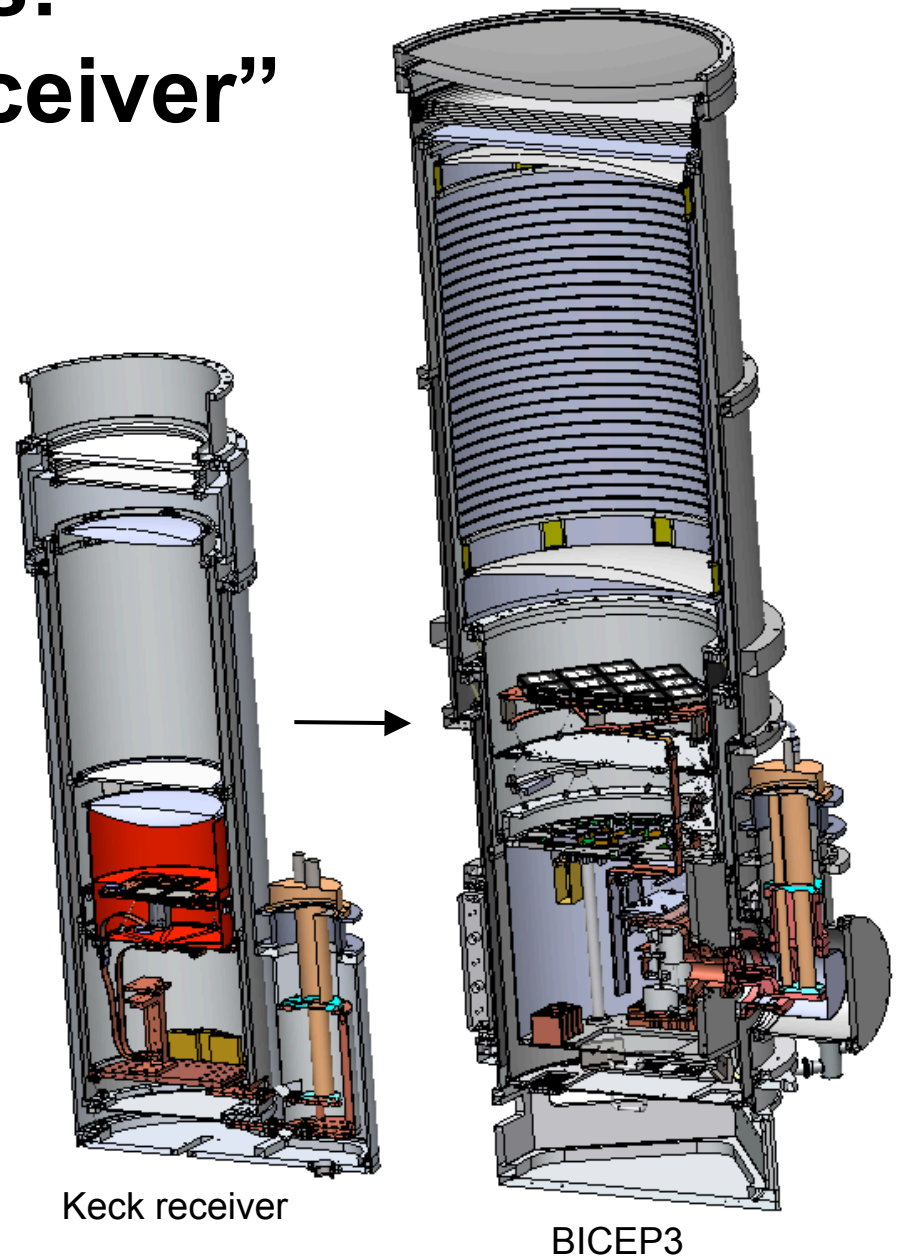
All 95 GHz

2560 detectors in modular  
focal plane

Larger-aperture optics

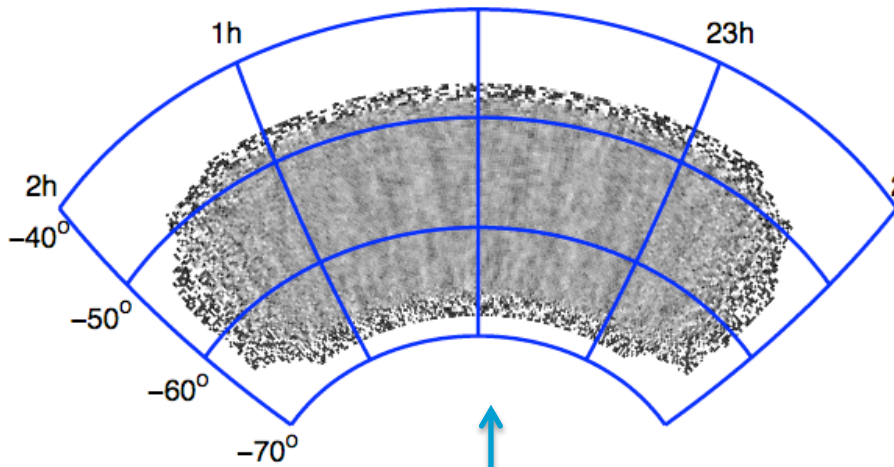
**> 10x optical throughput  
of single BICEP2/Keck  
receiver**

Means larger field of view and  
lower noise faster

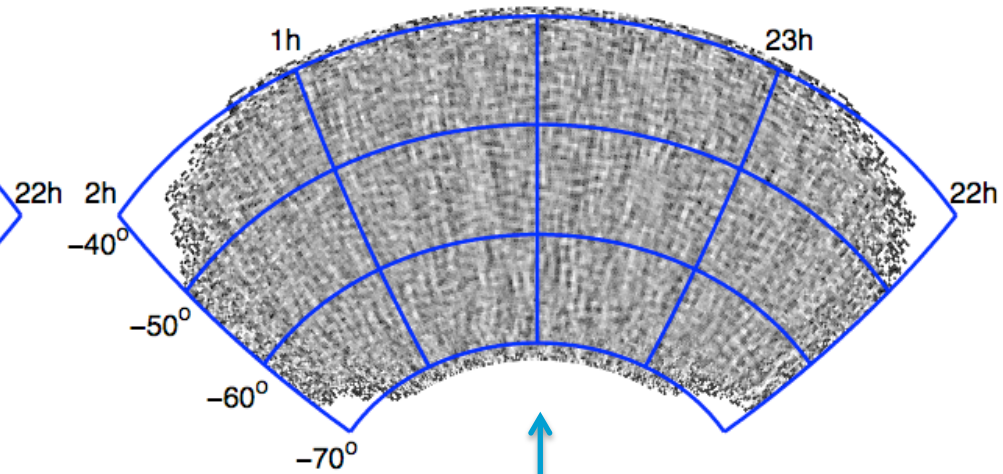




# Larger receiver = more sky area



Keck 95 GHz Q map after 4  
receiver years



BICEP3 95 GHz Q map after  
1 receiver year (2017)  
(Increased area, angular-  
resolution and sensitivity)

## Stage 2

## Stage 3

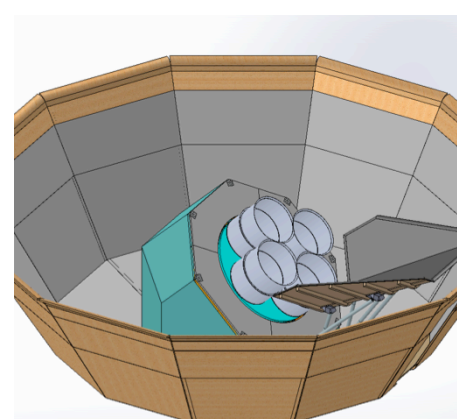
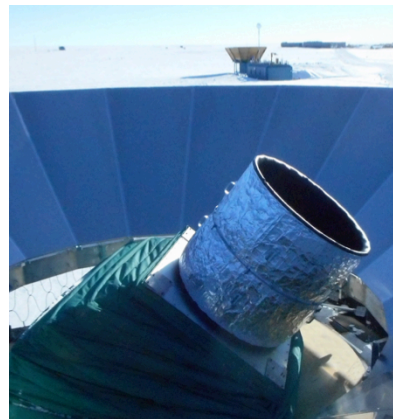
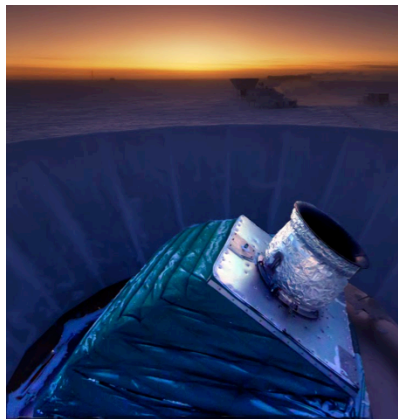
**BICEP2**  
(2010-2012)

**Keck Array**  
(2012-2019)

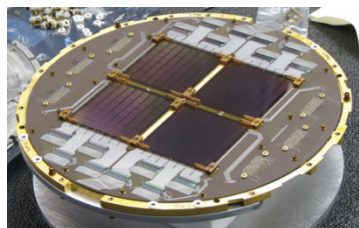
**BICEP3**  
(2016-)

**BICEP Array**  
(2020-)

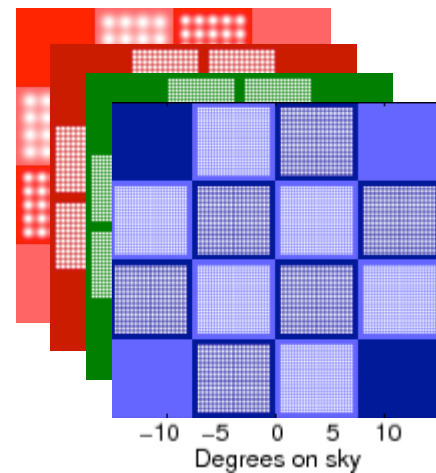
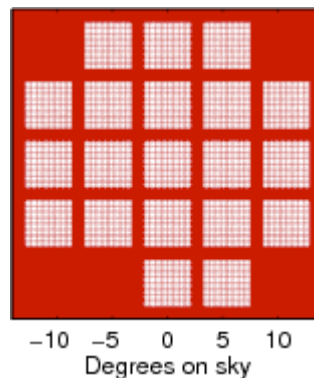
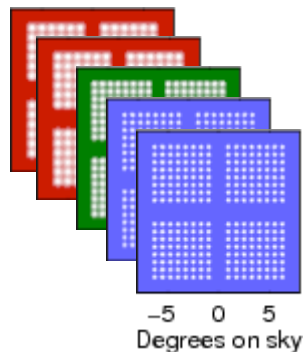
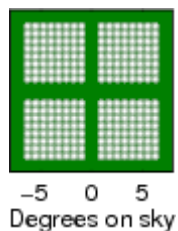
Telescope and Mount



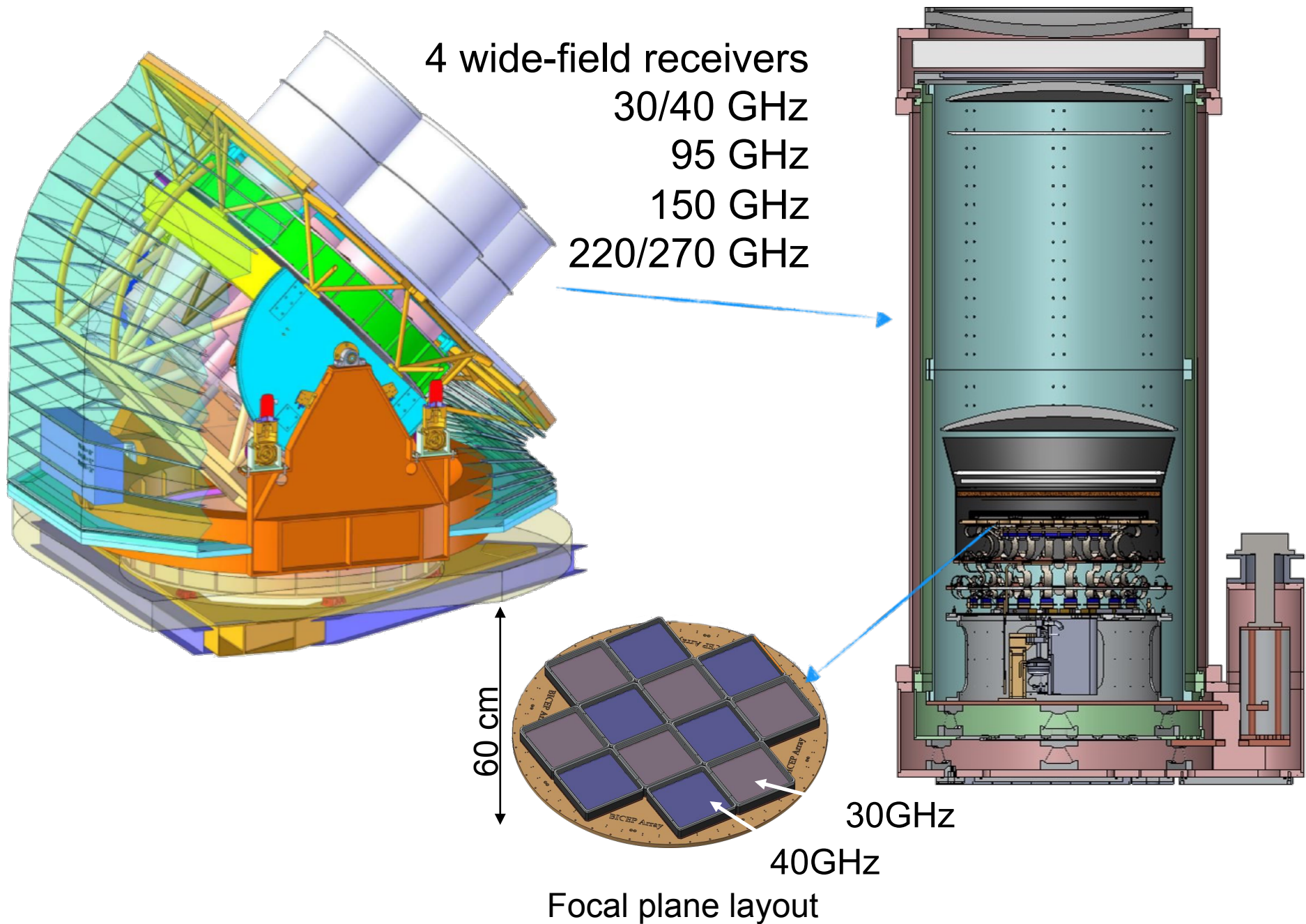
Focal Plane



Beams on Sky

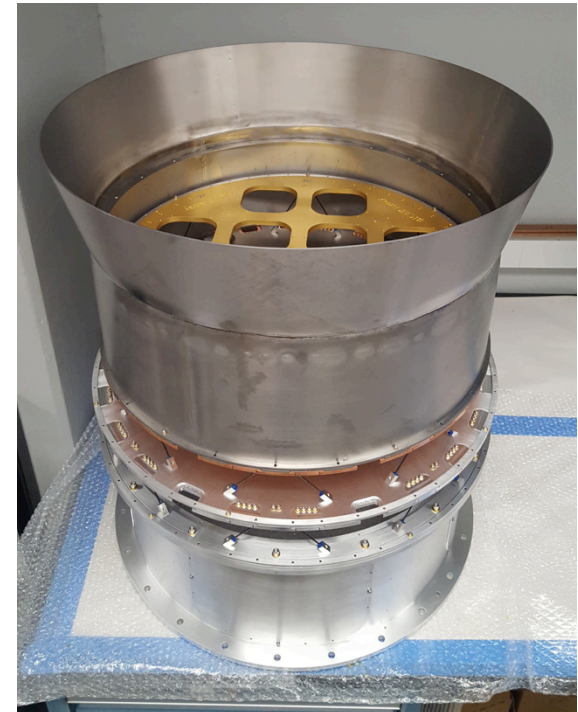


# Next Gen Experiment BICEP Array Under Construction

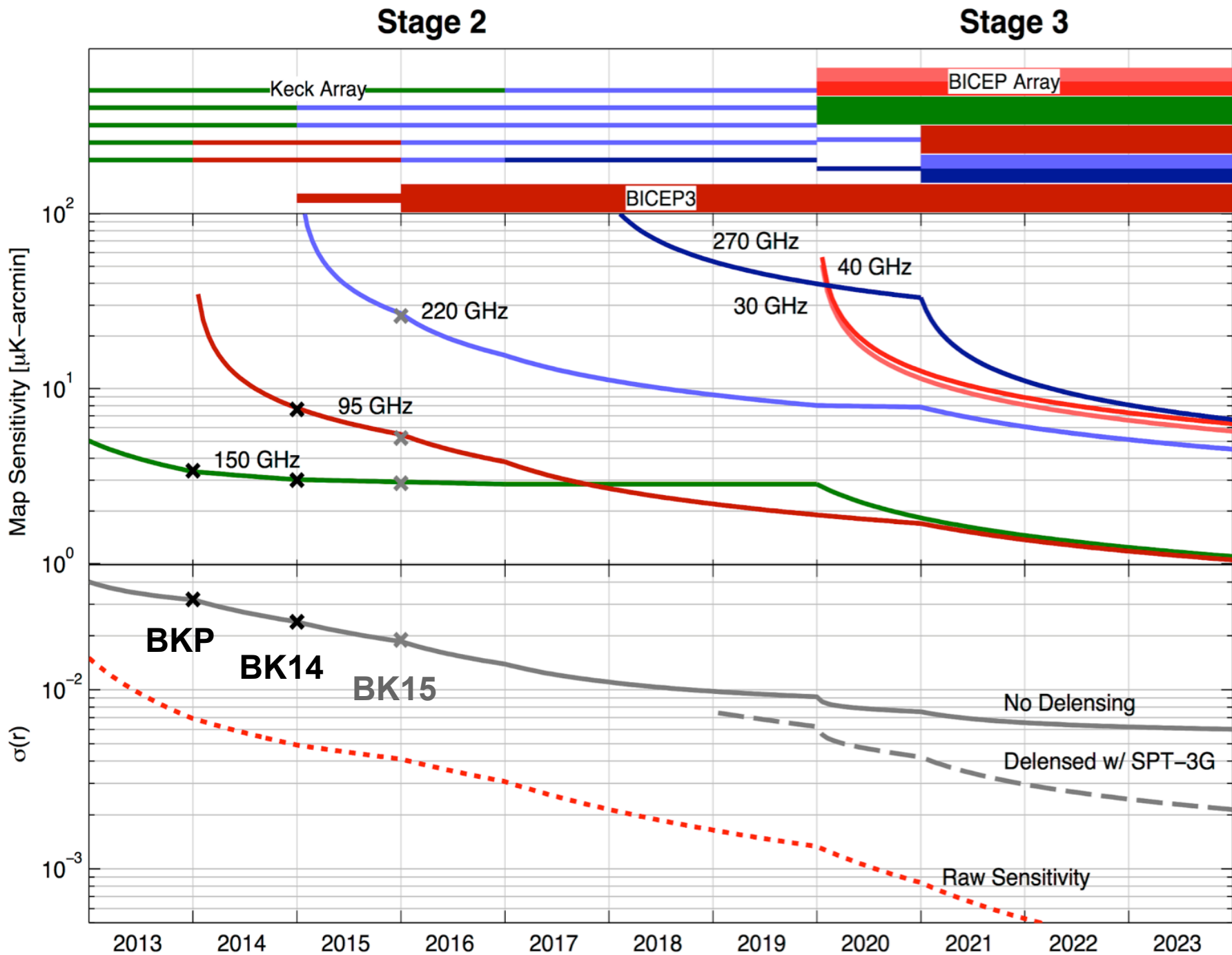




# BICEP Array Is Under Construction







# Conclusions

- BICEP/Keck lead the field in the quest to detect or set limits on inflationary gravitational waves:
- New BK15 result sets  $r_{0.05} < 0.07$  and  $\sigma(r) = 0.020$
- BICEP3 is running since 2016 with high sensitivity at 95GHz, and Keck Array continues to run at 220GHz, plus new 270GHz band
- We intend to go straight to BK17 (or BK18) analysis which will approach  $\sigma(r) = 0.010$
- BICEP Array is under construction and will go much further
- Next gen. receivers in five bands
- Delensing in conjunction with SPT3G is under development
- Project BK23  $\sigma(r) < 0.003$
- And beyond that is mega experiment CMB-S4...
- Foreground complexity will remain a serious issue – the hope is that we can measure it *and* constrain  $r$  simultaneously without a large loss of sensitivity. Additional ground/balloon measurements at low/high frequencies may be able to help.